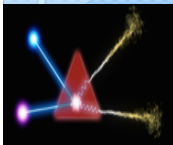


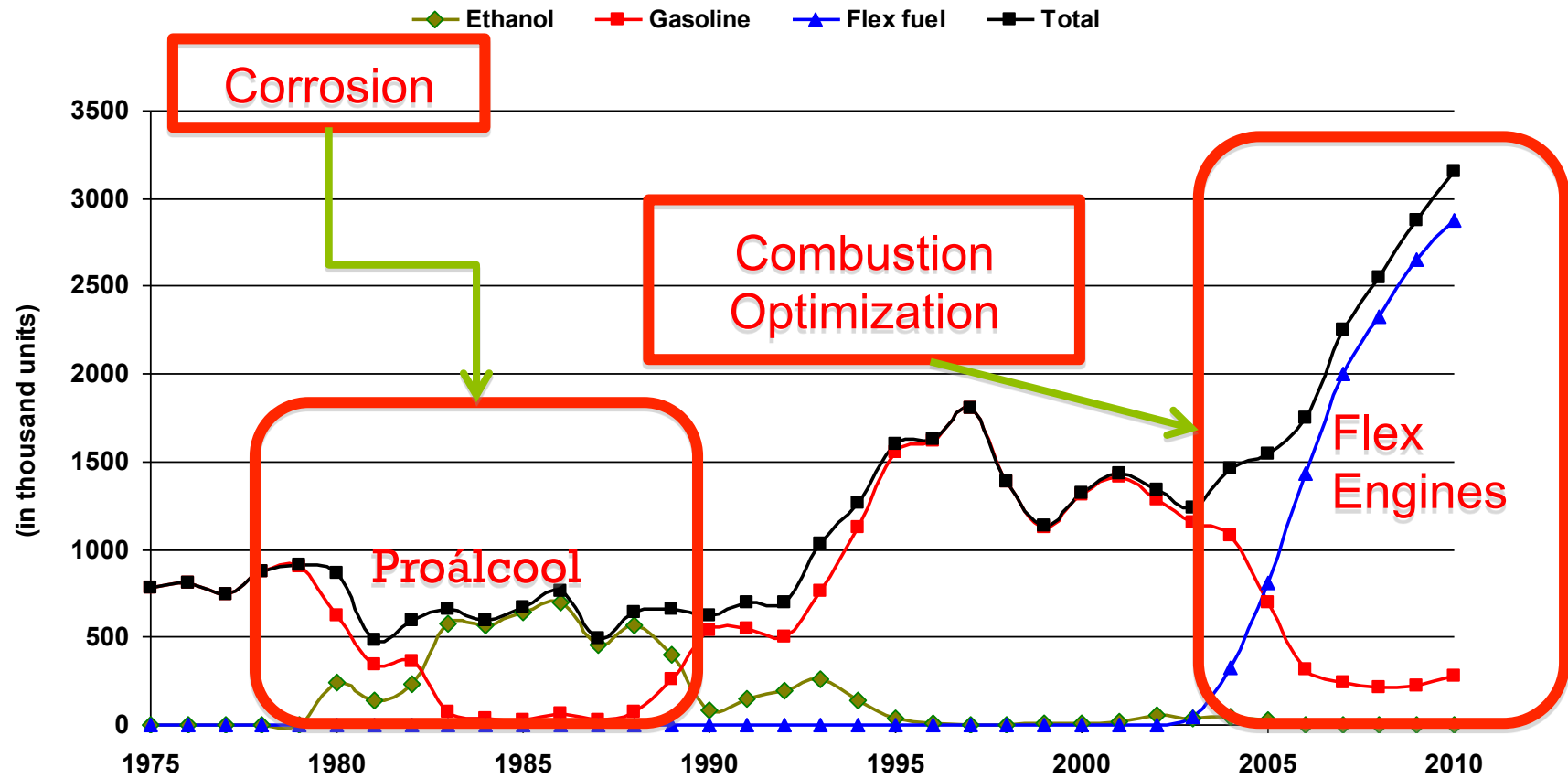
Espalhamento de elétrons por fragmentos moleculares da biomassa

de

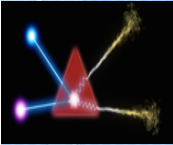
Marco A. P. Lima
Unicamp



Motivation I: large scale use of ethanol in engines



Brazilian Sales of light fleet Vehicles (1975-2010)

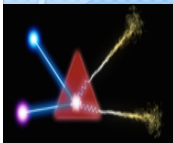


Ethanol as Fuel: Plasma Ignition for Vehicle Engines

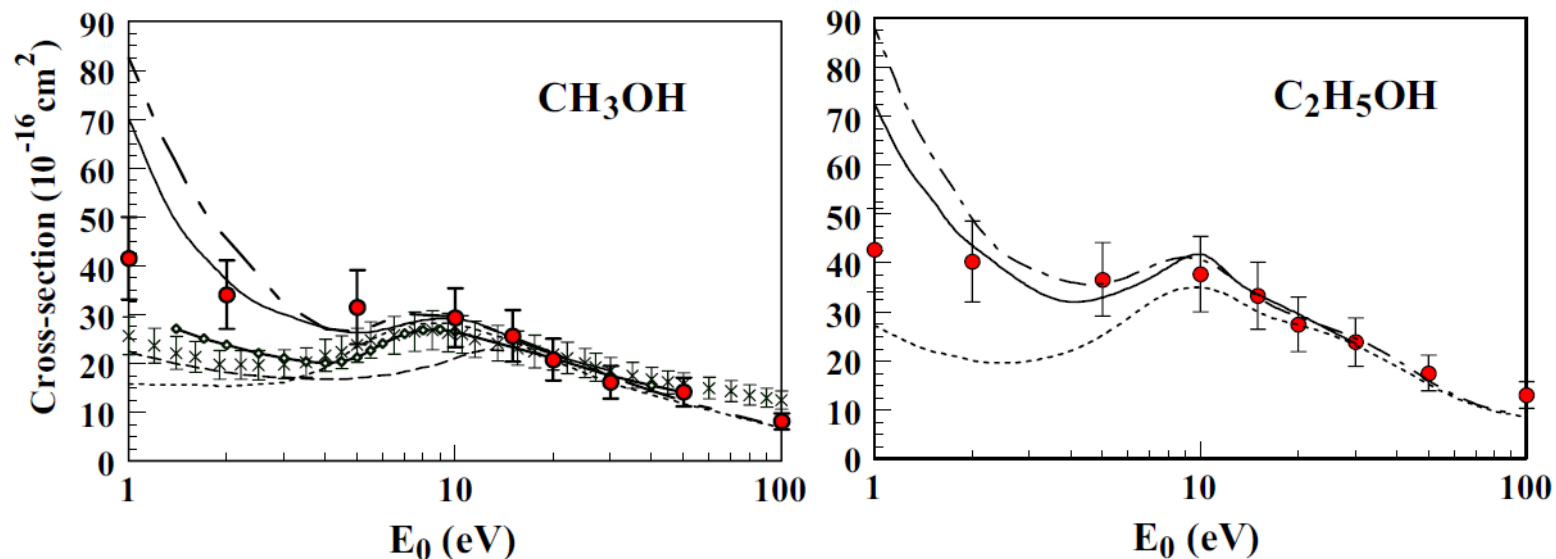


Theoretical support for an application project working on:

- Investigation of processes occurring during the ignition of plasma and its consequences in post-discharge for an internal combustion engine;
- The proper parameters to be applied in cars that operate on "poor mixtures" reducing pollutants released into the atmosphere, especially considering the spark plug discharge.



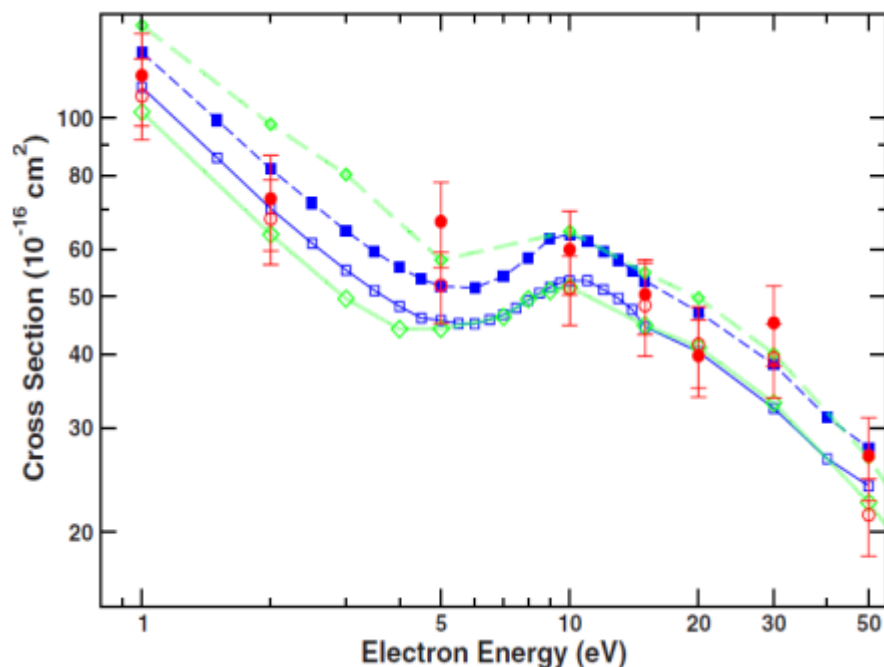
Low-energy elastic scattering from methanol and ethanol, M.A. Khakoo, J. Blumer, K. Keane, C. Campbell, H. Silva, M. C.A. Lopes, C. Winstead, V. Mckoy, R. F. da Costa, L. G. Ferreira, M.A. P. Lima, and M. H. F. Bettega, *Phys. Rev. A* **77**, 042705 (2008).



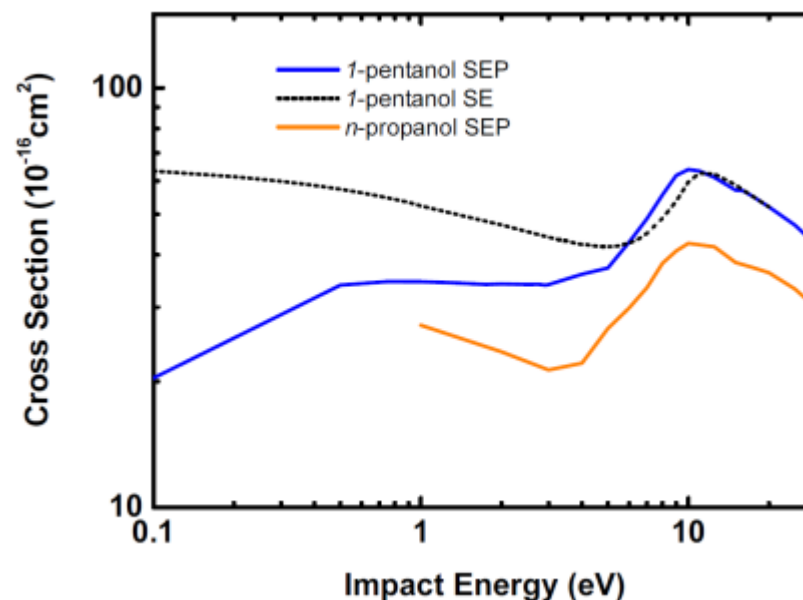
Integral elastic scattering cross sections for CH_3OH . Legend: ●: present experiment; —: SMCPP SEP; —: SMC SEP; ---- (short dashes): SMC SE which is similar to SMCPP SE; and ---- (long dashes): *R*-matrix ICSs of Bouchiha *et al.* (without Born correction) [10]. × Total cross section measurements of Szmytkowski and Krzysztofowicz [24] and —◆— of Schmieder [22]. ---- (short dashes) are from the SMCPP SE which is similar to SMC SE.



Low-energy elastic scattering from propanol and butanol, M.A. Khakoo, J. Blumer, K. Keane, C. Campbell, H. Silva, M. C.A. Lopes, C. Winstead, V. Mckoy, M. H. F. Bettega, E. M. de Oliveira, M.T. do N. Varella, R. F. da Costa, and M.A. P. Lima, and, *Phys. Rev. A* **78**, 0622714 (2008).

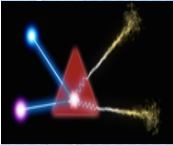


Electron scattering of slow electrons by 1-pentanol (a drop in fuel)



Integral elastic cross sections for electron collisions with *n*-propanol and *n*-butanol. Red circles are experimental values, thin blue lines with squares results from all-electron calculations, and thick green lines with diamonds results from pseudopotential calculations. Open symbols and solid lines are propanol data; solid symbols and dashed lines are butanol data.

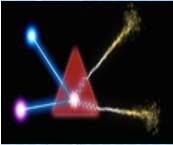
NSF (Caltech/Fullerton)/CNPq (Unicamp/UFJF/UFPR/USP/UFABC) project



Motivation II: large scale production of ethanol



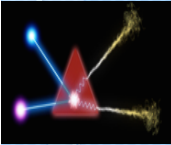
A sugarcane industry of Sugar/Ethanol/Bioelectricity



Motivation II: large scale production of ethanol



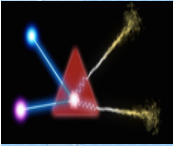
Biomass: a source of energy and carbon



Motivation II: large scale production of ethanol



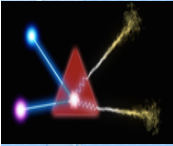
Biomass: a source of energy and carbon



Motivation II: large scale production of ethanol



First generation ethanol: crushing the cane for the juice



Motivation II: large scale production of ethanol



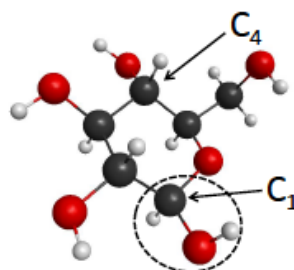
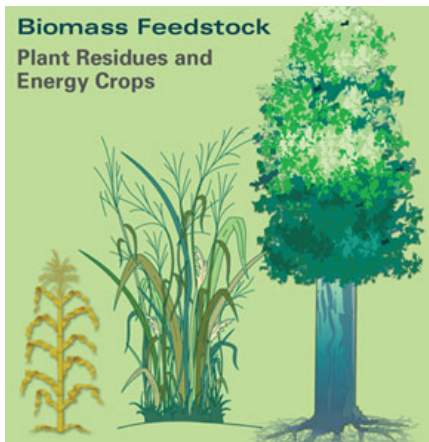
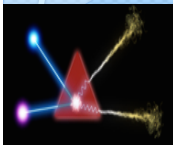
**Bagasse piles
at the mill.**

**2nd generation
ethanol?
Other high value
bioproducts?**

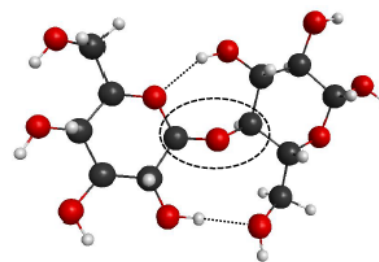


Biomass: a source of energy and carbon

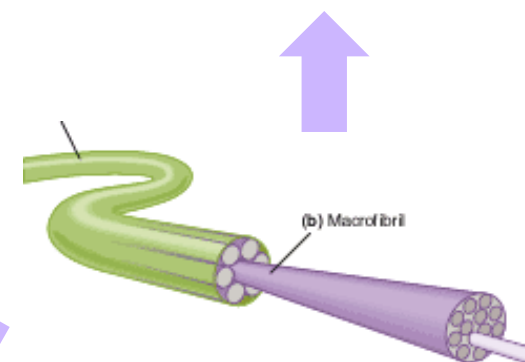
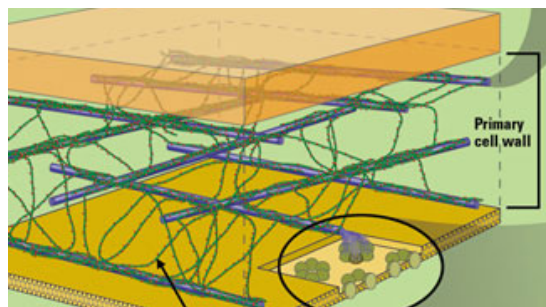
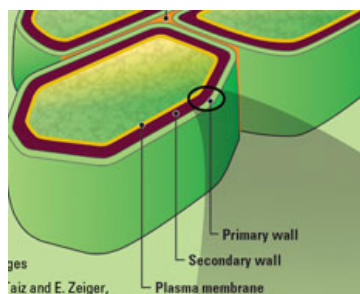
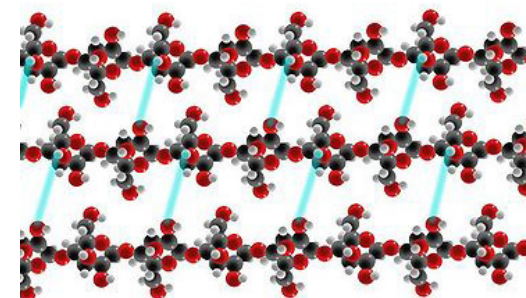
Biomass is Made Up with Fermentable Sugars



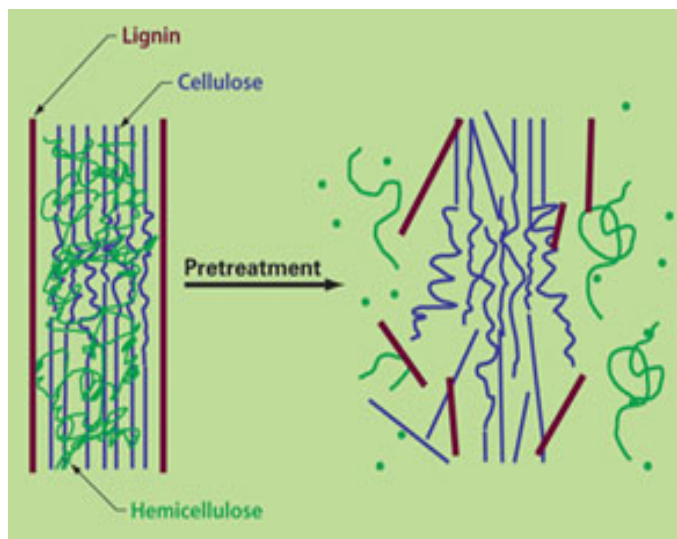
β -D-glucose



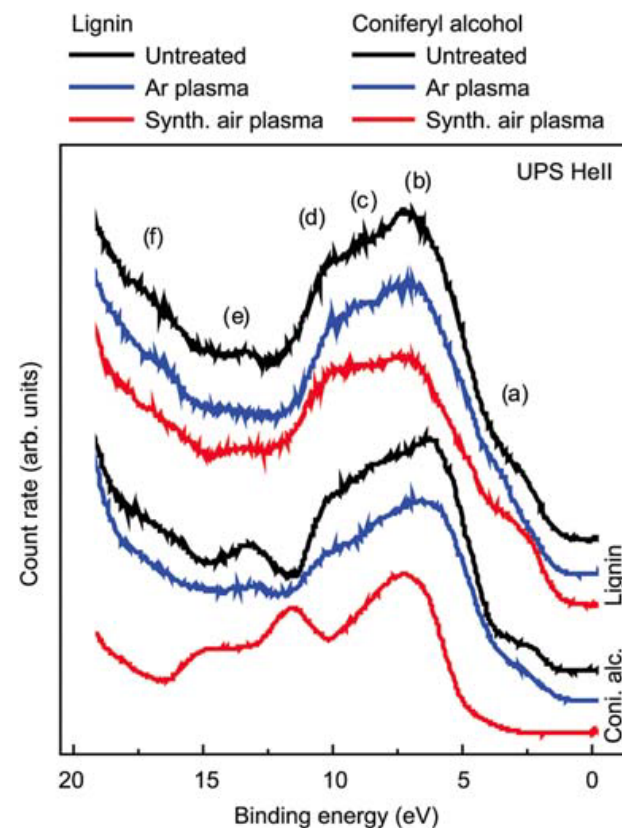
cellobiose



Lignocellulose is Resistant to Hydrolysis



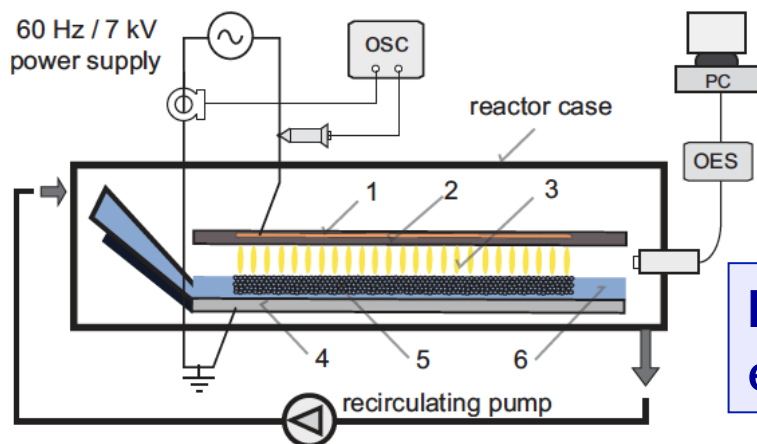
Pretreatment: bio- and physical-chemical processes to expose the cellulose fibers

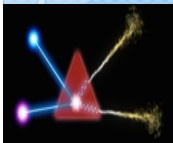


Lothar Klarhöfer¹, Wolfgang Viöl^{2,3,*} and Wolfgang Maus-Friedrichs¹

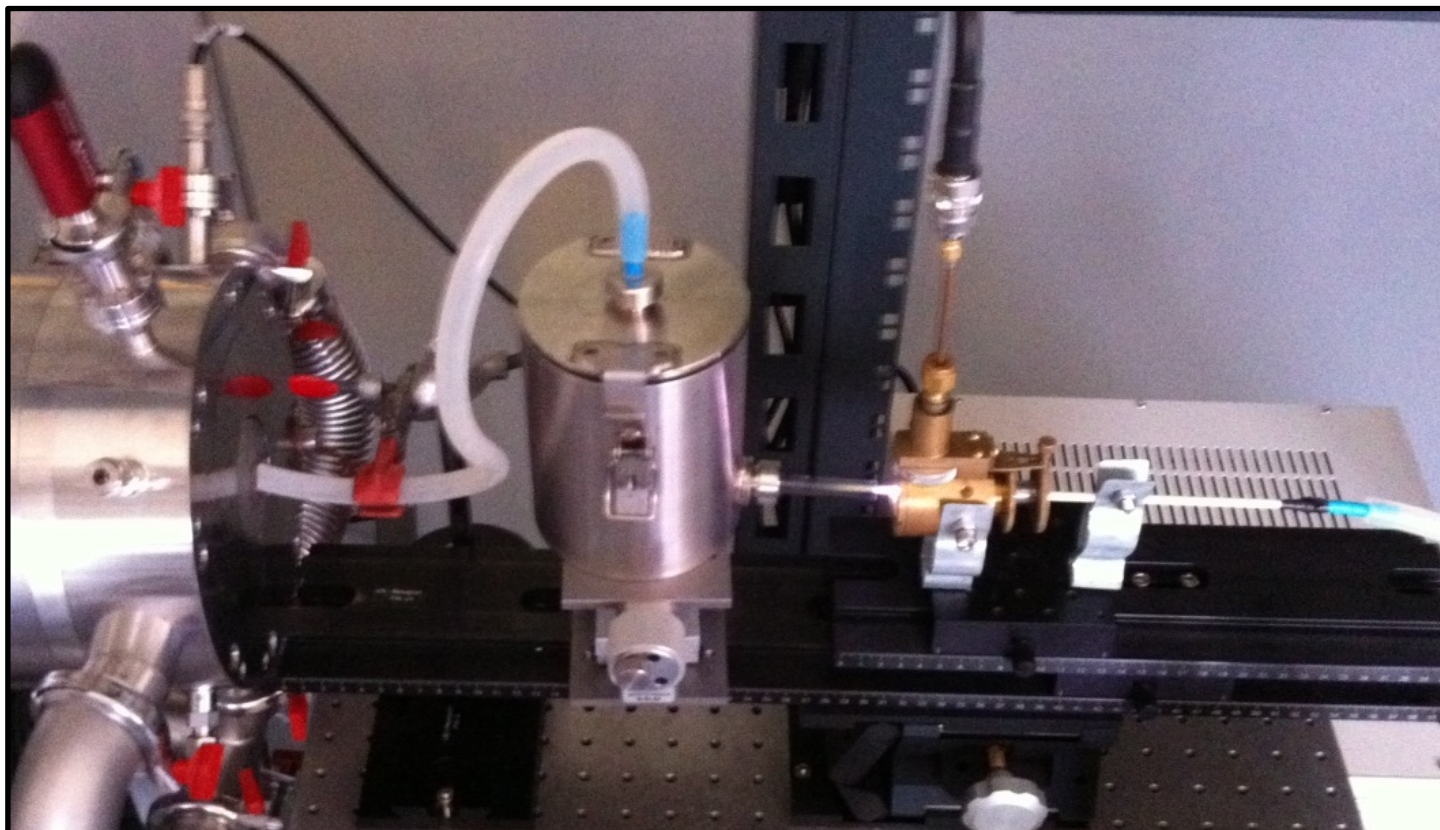
Holzforschung, Vol. 64, pp. 331–336, 2010

**Dielectric Barrier Discharge (DBD):
electron flux on substrate $\sim 10^8 \text{ cm}^{-2} \text{ s}^{-1}$**





Sugarcane Bagasse Plasma Pretreatment



Treatment conditions

- ~ 25 g of dry sugarcane bagasse (50% moisture) – milled at 500 μ m
- Gas flow Mixture: 95% Ar (1.9 SLM) and 5% O₂ (0.1 SML)
- $\Delta t_{\text{treatment}} = 3\text{h}$



Biomass Chemical Analysis

Lignin concentration (%) of raw bagasse and samples related to plasma torch treatment and washing procedure by water and NaOH 1% solution at room temperature.

Samples	Soluble Lignin (%)	Insoluble Lignin (%)	Total of Lignin remaining (%)
raw bagasse	1.58 ± 0.01	20.3 ± 0.1	21.9 ± 0.1
Washed by H ₂ O	2.4 ± 0.9	21.4 ± 0.9	23.8 ± 0.9
Washed by NaOH 1%	1.3 ± 0.9	12.6 ± 0.9	13.9 ± 0.9

About 40% of original lignin was removed!!!

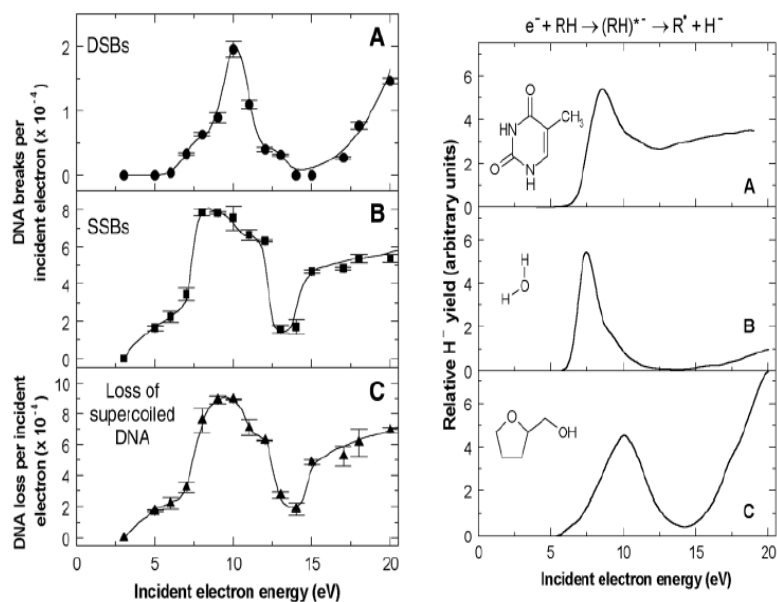
Jayr Amorim, Carlos Oliveira, Jorge A. Souza-Correa, Marco A. Ridenti
Plasma Process. Polym. 2013, DOI: 10.1002/ppap.201200158



Electron-Induced Damage to Biomolecules

Resonant Formation of DNA Strand Breaks by Low-Energy (3 to 20 eV) Electrons

Badia Boudaïffa, Pierre Cloutier, Darel Hunting,
Michael A. Huels,* Léon Sanche



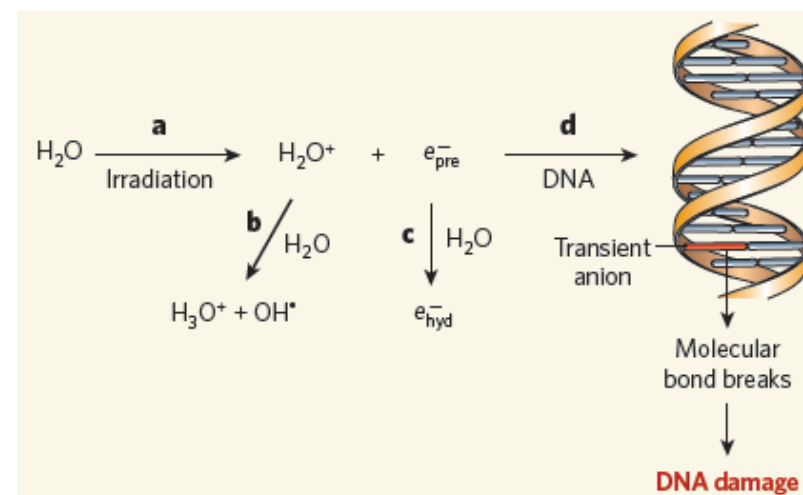
Science, **287** 1658 (2000)

BIOLOGICAL CHEMISTRY

Beyond radical thinking

Léon Sanche

Radiation-induced DNA damage has been attributed to hydroxyl radicals, which form when water absorbs high-energy photons or charged particles. But another product of water's radiolysis might be the real culprit.

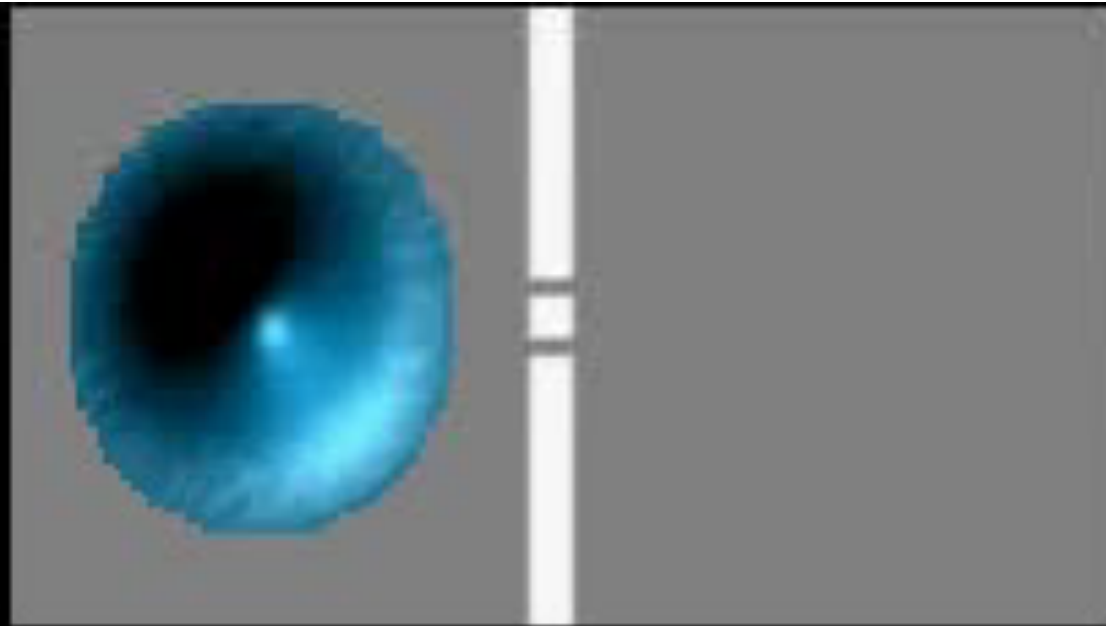
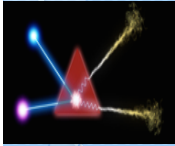


Sanche, Nature **461**, 358 (2009)

J|A|C|S
COMMUNICATIONS

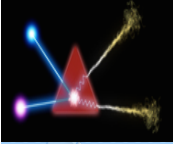
Chun-Rong Wang, Jenny Nguyen, and Qing-Bin Lu*

J. AM. CHEM. SOC. 2009, **131**, 11320–11322



Two slits interference

To see this animations, visit: <http://www.embd.be/quantummechanics/>



Scattering theory

Schrödinger equation

$$H\Psi_{\vec{k}_m}^{(\pm)}(\vec{r}_1, \dots, \vec{r}_{N+1}) = E\Psi_{\vec{k}_m}^{(\pm)}(\vec{r}_1, \dots, \vec{r}_{N+1})$$

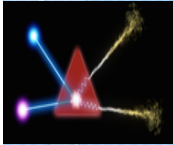
Asymptotic condition

$$\Psi_{\vec{k}_i}^{(\pm)}(\vec{r}_1, \dots, \vec{r}_{N+1}) \xrightarrow{r_{N+1} \rightarrow \infty} S_{\vec{k}_i} + \sum_f^{\text{open}} f_{i \rightarrow f}^B(\vec{k}_i, \vec{k}_f) \Phi_f \frac{e^{\pm i k_f r_{N+1}}}{r_{N+1}}$$

$$S_{\vec{k}_i} = \Phi_i e^{i \vec{k}_i \cdot \vec{r}_{N+1}}$$

Differential cross section

$$\frac{d\sigma^{i \rightarrow f}}{d\Omega}(\vec{k}_i, \vec{k}_f) = \frac{k_f}{k_i} \left| f_{i \rightarrow f}^L(\vec{k}_i, \vec{k}_f) \right|^2$$



Scattering theory

Schrödinger differential equation

$$H\Psi_{\vec{k}_m}^{(\pm)} = [H_N + T_{N+1} + V]\Psi_{\vec{k}_m}^{(\pm)} = E\Psi_{\vec{k}_m}^{(\pm)}$$

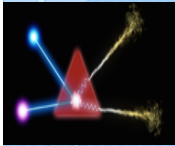
Lippmann-Schwinger integral equation

$$\Psi_{\vec{k}_m}^{(\pm)} = S_{\vec{k}_m} + G_0^{(\pm)}V\Psi_{\vec{k}_m}^{(\pm)}$$

$$S_{\vec{k}_m} = \Phi_m e^{i\vec{k}_m \cdot \vec{r}_{N+1}}$$

Free-particle Green's function (Source of uncertainty – controllable)

$$G_0^{(\pm)} = \frac{1}{E - T_{N+1} - H_N \pm i\epsilon} = \lim_{\epsilon \rightarrow 0} \mathcal{P} \int d^3k \frac{|\Phi_m \vec{k}\rangle \langle \vec{k} \Phi_m|}{\frac{k_m^2}{2} - \frac{k^2}{2} \pm i\epsilon}$$



Schwinger Variational Principle

The Schwinger Variational method serves to get a scattering amplitude free of first order errors for a scattering process that respect the equations

$$A^{(\pm)}|\Psi_{\mathbf{k}}^{(\pm)}\rangle = V|S_{\mathbf{k}}\rangle \text{ and } \begin{cases} f_{\mathbf{k}_f, \mathbf{k}_i} = \langle S_{\mathbf{k}_f} | V | \Psi_{\mathbf{k}_i}^{(+)} \rangle \\ f_{\mathbf{k}_f, \mathbf{k}_i} = \langle \Psi_{\mathbf{k}_f}^{(-)} | V | S_{\mathbf{k}_i} \rangle \\ f_{\mathbf{k}_f, \mathbf{k}_i} = \langle \Psi_{\mathbf{k}_f}^{(-)} | A^{(+)} | \Psi_{\mathbf{k}_i}^{(+)} \rangle \end{cases} \quad \text{and } A^{(\pm)} = V - VG_0^{(\pm)}V$$

The bilinear form of the variational principle for the scattering amplitude is

$$[f_{\mathbf{k}_f, \mathbf{k}_i}] = \langle S_{\mathbf{k}_f} | V | \Psi_{\mathbf{k}_i}^{(+)} \rangle + \langle \Psi_{\mathbf{k}_f}^{(-)} | V | S_{\mathbf{k}_i} \rangle - \langle \Psi_{\mathbf{k}_f}^{(-)} | A^{(+)} | \Psi_{\mathbf{k}_i}^{(+)} \rangle \text{ where arbitrary and}$$

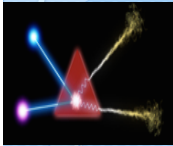
independent variations with respect to

$$\begin{cases} \langle \delta \Psi_{\mathbf{k}_f}^{(-)} | (V | S_{\mathbf{k}_i} \rangle - A^{(+)} | \Psi_{\mathbf{k}_i}^{(+)} \rangle) = 0 \\ (\langle S_{\mathbf{k}_f} | V - \langle \Psi_{\mathbf{k}_f}^{(-)} | A^{(+)} \rangle) | \delta \Psi_{\mathbf{k}_i}^{(+)} \rangle = 0 \end{cases}$$

lead to

$$\begin{cases} V | S_{\mathbf{k}_i} \rangle - A^{(+)} | \Psi_{\mathbf{k}_i}^{(+)} \rangle = 0 \Rightarrow A^{(+)} | \Psi_{\mathbf{k}_i}^{(+)} \rangle = V | S_{\mathbf{k}_i} \rangle \\ \langle S_{\mathbf{k}_f} | V - \langle \Psi_{\mathbf{k}_f}^{(-)} | A^{(+)} = 0 \Rightarrow A^{(-)} | \Psi_{\mathbf{k}_f}^{(-)} \rangle = V | S_{\mathbf{k}_f} \rangle \text{ with } A^{(+)\dagger} = A^{(-)} \end{cases}$$

$A^{(\pm)}|\Psi_{\mathbf{k}}^{(\pm)}\rangle = V|S_{\mathbf{k}}\rangle$ must be equivalent to $H|\Psi_{\mathbf{k}}^{(\pm)}\rangle = E|\Psi_{\mathbf{k}}^{(\pm)}\rangle$ with proper boundary conditions



Schwinger Multichannel Method for electron scattering

In this formalism the operator $A^{(+)}$ was redefined as:

$$A^{(+)} = \frac{1}{2}(PV + VP) - VG_P^{(+)}V + \frac{1}{N+1} \left[\hat{H} - \frac{N+1}{2}(\hat{H}P + P\hat{H}) \right]$$

where $P \equiv \sum_{\ell=1}^{\text{open}} |\Phi_\ell\rangle\langle\Phi_\ell|$ and $\hat{H} = E - H$

 (Channel coupling: source of uncertainty – uncontrollable)

All electrons are identical. So, an expansion of the scattering wave function must be done in a basis $\{\chi_\mu\}$ of anti-symmetric functions (Slater determinants):

$$|\Psi_{\vec{k}_m}^{(\pm)}\rangle = \sum_{\mu} a_{\mu}^{(\pm)}(\vec{k}_m) |\chi_{\mu}\rangle \quad \text{where} \quad \{|\chi_{\mu}\rangle\} = \{a_{N+1} |\Phi_i\rangle \otimes |\varphi_j\rangle\}$$

 (N+1 wave-function expansion: source of uncertainty – controllable)

The final form of the scattering amplitude is equal to the one of the Schwinger Variational principle

$$f_{\vec{k}_i, \vec{k}_f} = -\frac{1}{2\pi} \sum_{mn} \langle S_{\vec{k}_f} | V | \chi_m \rangle (d^{-1})_{mn} \langle \chi_n | V | S_{\vec{k}_i} \rangle$$

with $d_{mn} = \langle \chi_m | A^{(+)} | \chi_n \rangle$ and $S_{\vec{k}_i} \equiv \Phi_i(\vec{r}_1, \dots, \vec{r}_N) e^{i\vec{k}_i \cdot \vec{r}_{N+1}}$



Coupling level

➔ Elastic scattering with and without polarization effects

① Open channel Projector has only one state

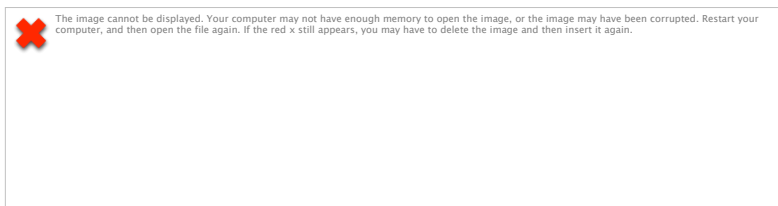
$$P = |\Phi_0\rangle\langle\Phi_0|$$

➔ Φ_0 is molecular target ground state obtained in Hartree-Fock approximation

(Target description: source of uncertainty – program limitation)

② Configuration space is made of

(N+1 wave-function expansion: source of uncertainty – controllable)



Doublet states made of products of target triplet and singlet states by φ_k



$\Phi_j, j \geq 2$ are virtual states obtained from single excitations of the molecular target



φ_i are one-particle wave functions (square integrable molecular orbitals) used in description of the continuum



Coupling level

➔ Inelastic scattering with and without polarization

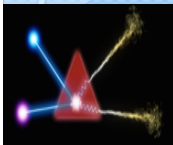
- 1 Open channel projector contains channels of our choice (truncation means approximation)

$$P = \sum_{\ell}^{\text{open}} |\Phi_{\ell}\rangle \langle \Phi_{\ell}| \quad \Rightarrow \quad |\Phi_{\ell}\rangle \text{ are molecular target states obtained with single configuration interaction}$$

- 2 Again the configuration space is made of

$$|\chi_{\mu}\rangle = \begin{cases} \mathbf{a}_{N+1} |\Phi_0\rangle \otimes |\varphi_i\rangle \\ \mathbf{a}_{N+1} |\Phi_j\rangle \otimes |\varphi_k\rangle, j \geq 2 \end{cases} \quad \Rightarrow \quad \text{Doublet states made of products of target triplet and singlet states by } \varphi_k$$

➔ Polarization effects are included with j greater than the number of open channels



Electron scattering by large molecules



Pseudopotential formalism

(source of uncertainty – optional for light molecules)

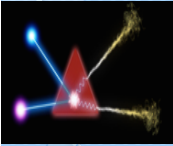
[D. R. Hamann, M. Schlüter and C. Chiang, Phys. Rev. Lett. **43**, 1494 (1979)]

- 1 The pseudo-state energy is equal to the real eigenvalue for a given configuration;
- 2 The pseudopotential is equal to the real potential beyond a certain core radius r_c , and it is soft at the origin;
- 3 The normalized pseudo wave function is equal to the real one beyond the core radius r_c and it is soft and without nodes
- 4 The integrals from 0 to r of the real and pseudo functions agree for $r > r_c$ for each valence state:

$$-\frac{1}{2} \left[(r\Psi)^2 \frac{d}{dE} \frac{d}{dr} \ln \Psi \right]_{r>r_c} = \int_0^{r>r_c} \Psi^2 r^2 dr$$



“Norm Conservation”



Theoretical co-authors



Eliane M. de Oliveira (posdoc)
Alexandra Natalense
Marco A. P. Lima



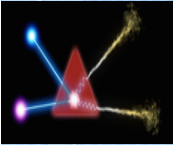
Sergio d'A. Sanchez
Márcio H. F. Bettega



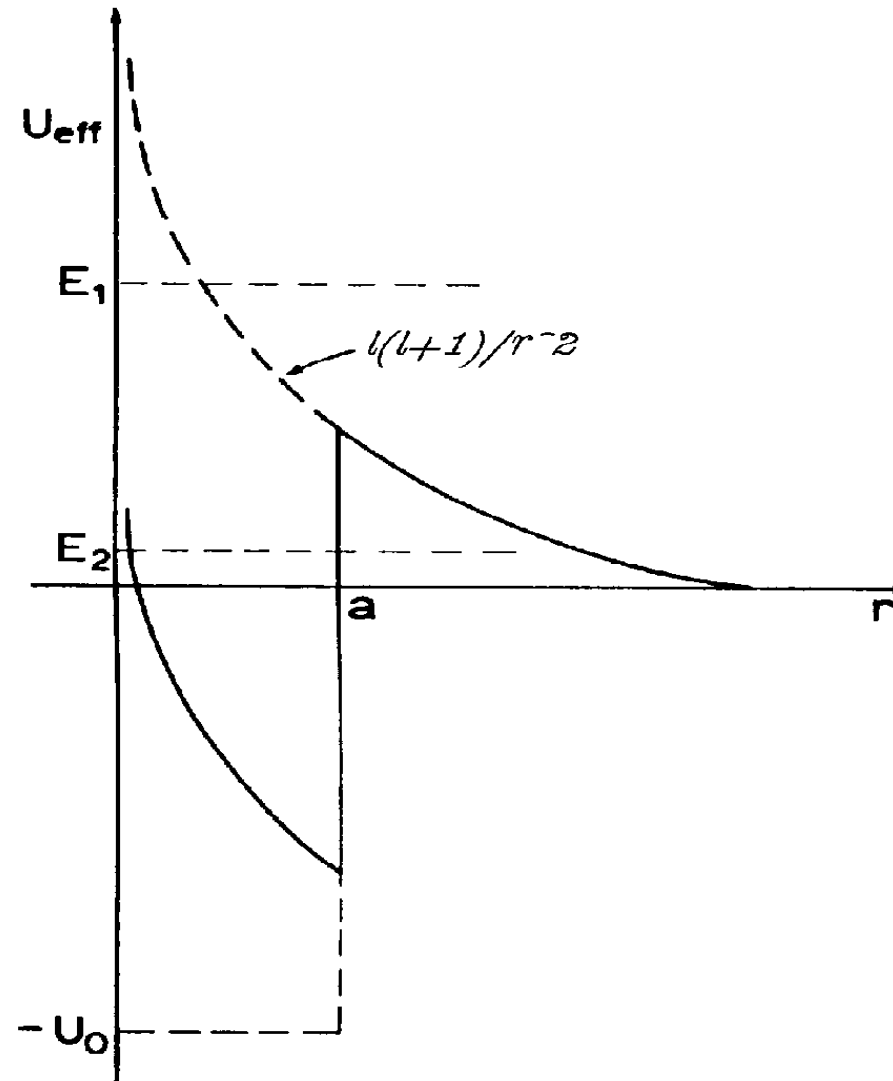
Romarly F. da Costa

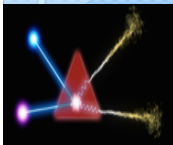


Márcio T. do N. Varella (coordinator)



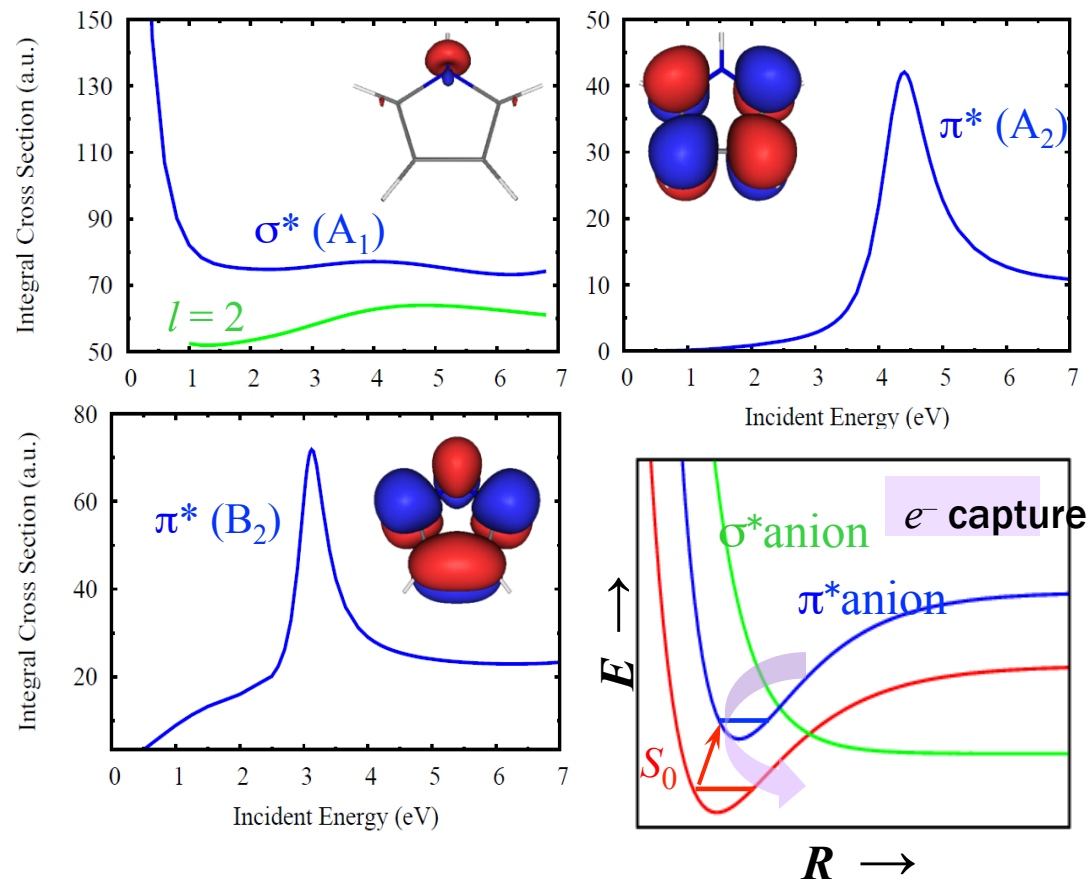
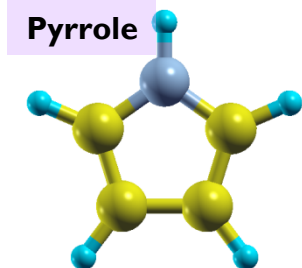
Shape resonances are related to angular momentum traps





Low energy elastic electron scattering from pyrrole

- There are π^* (ring) and σ^* (N-H) shape resonances in pyrrole. Nice prototype!

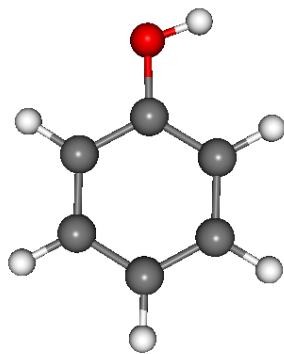


de Oliveira EM, Lima MAP, Bettega MHF, Sanchez SD, da Costa RF, and Varella MTD, J. Chem. Phys. **132**, 204301 (2010)



Lignin Subunits

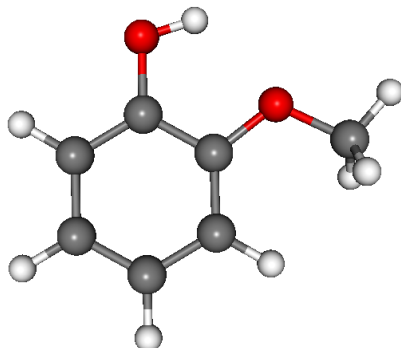
Phenol



MetOH



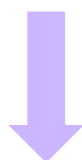
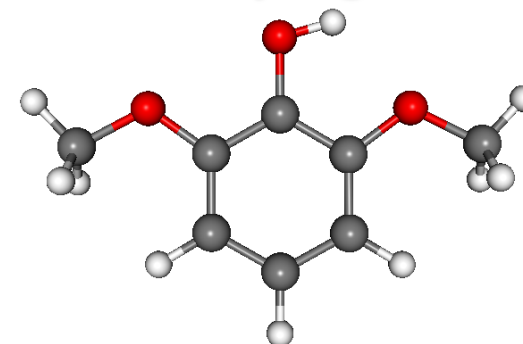
Guaiacol



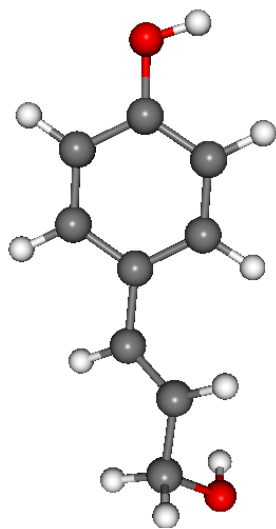
MetOH



Syringol



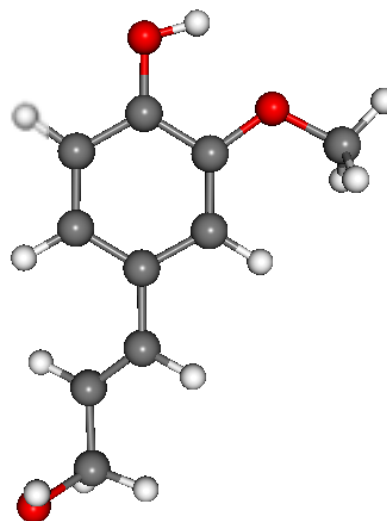
PropenylOH



p-coumaryl alcohol



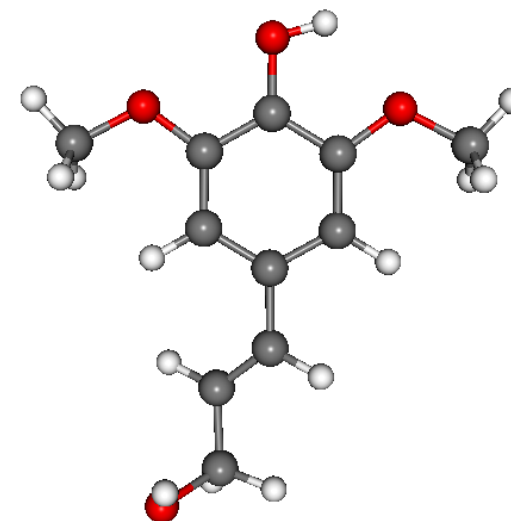
PropenylOH



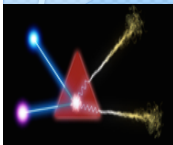
coniferyl alcohol



PropenylOH



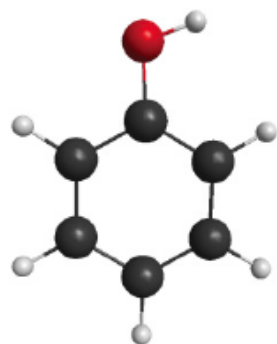
sinapyl alcohol



RAPID COMMUNICATIONS

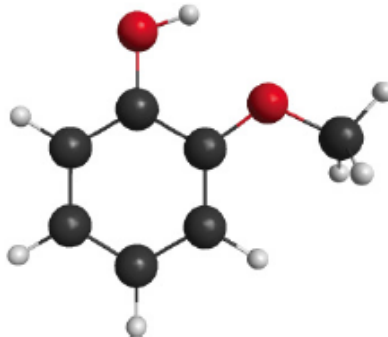
Shape resonance spectra of lignin subunits

PHYSICAL REVIEW A 86, 020701(R) (2012)



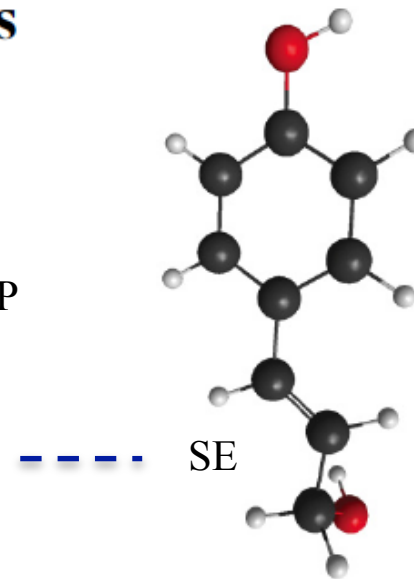
phenol

— SEP
- - - SE



guaiacol

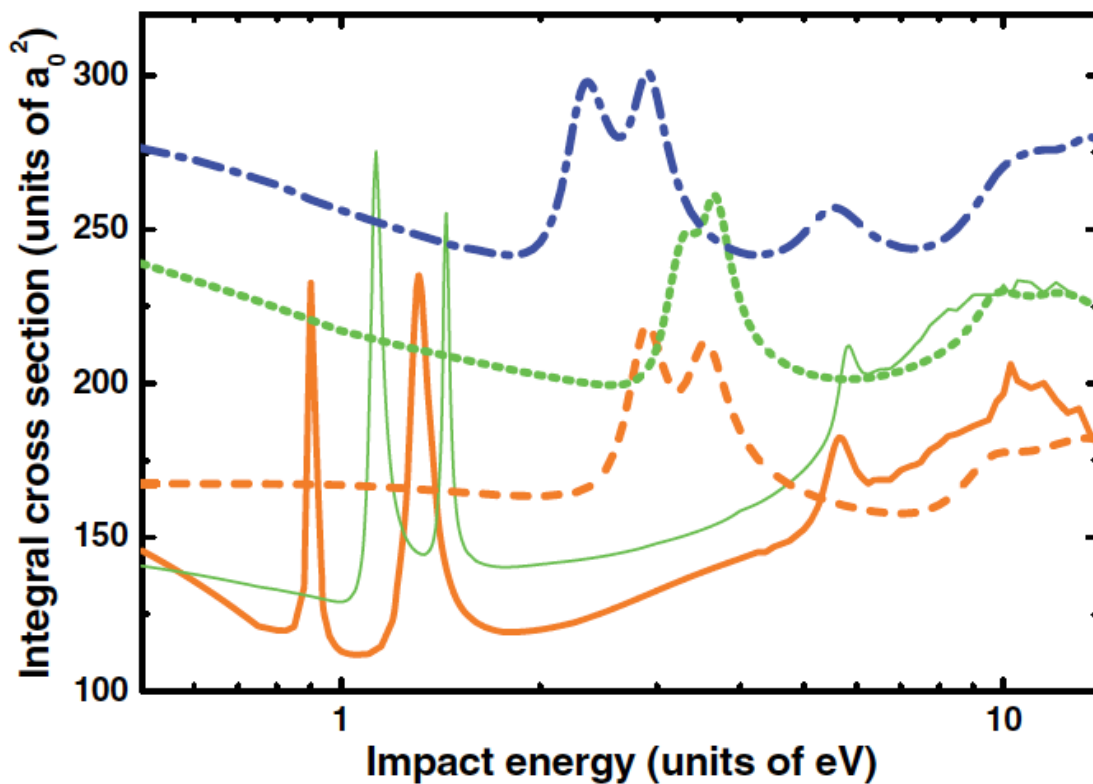
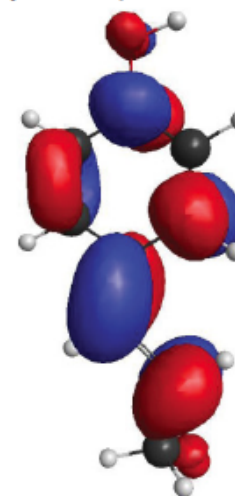
— SEP
- - - SE

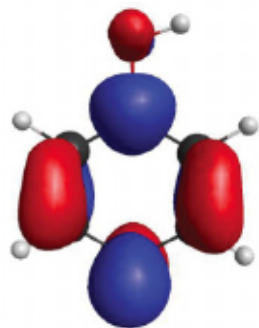


p-coumaryl alcohol

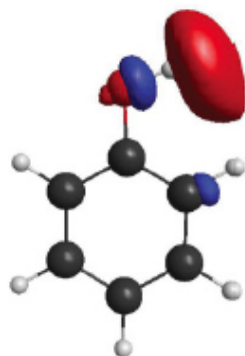
- - - SE

p-Cu (LUMO)



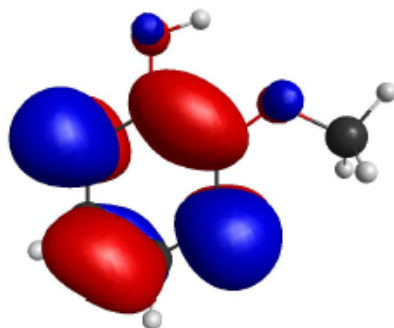


π^* (LUMO+1)

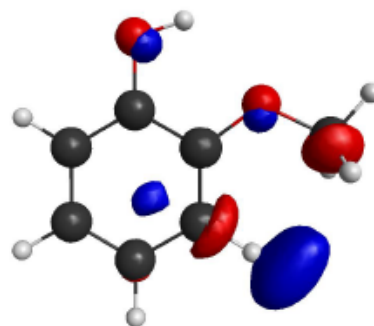


σ^* (LUMO+2)

Phenol: Calculations, ET spectra and DEA data indicate H elimination from π^*/σ^* coupling.

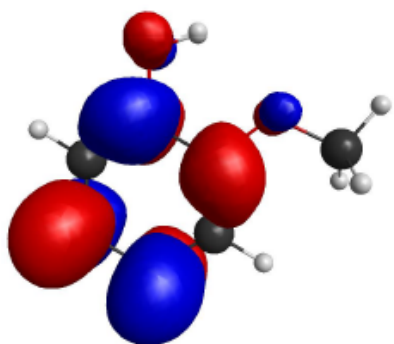


π^* (LUMO)

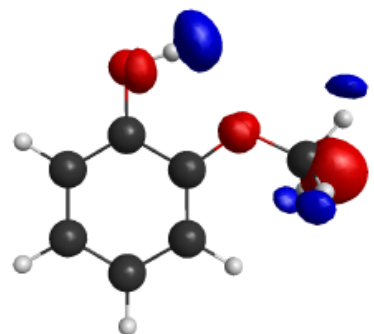


σ^* (LUMO+2)

Guaiacol: Methoxilation is expected to give rise to other dissociation channels. H elimination should be also observed.



π^* (LUMO+1)

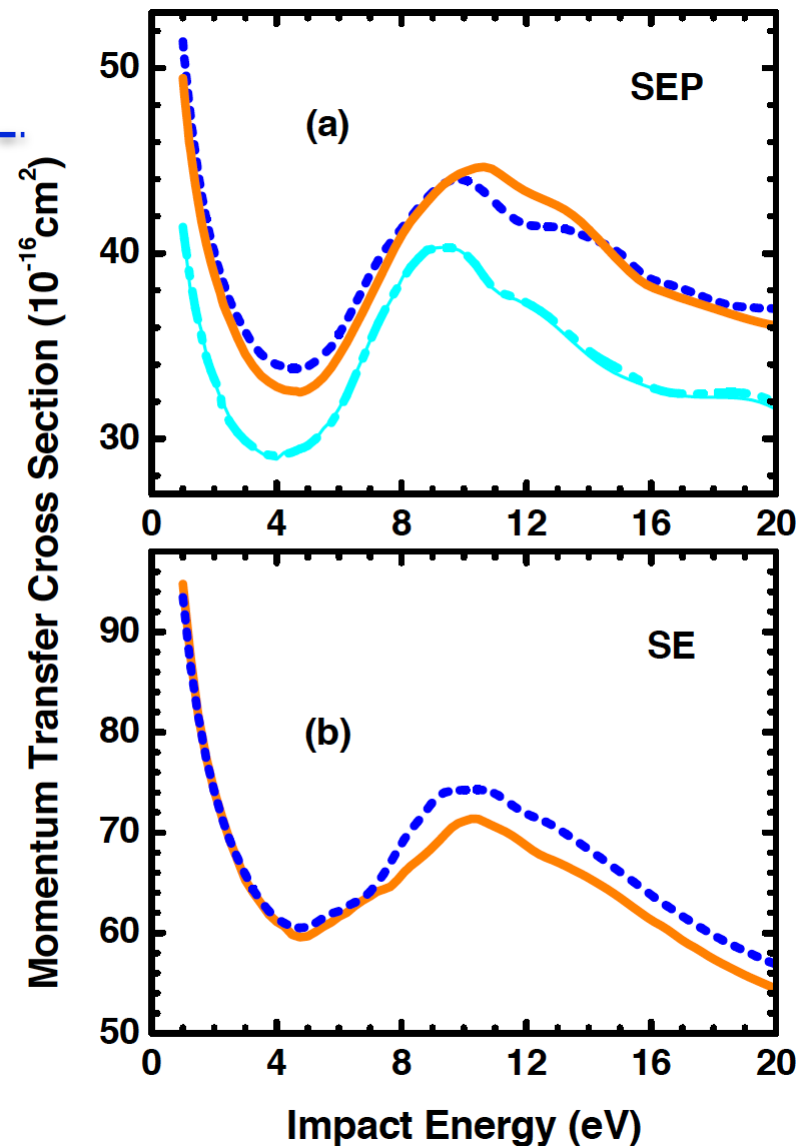
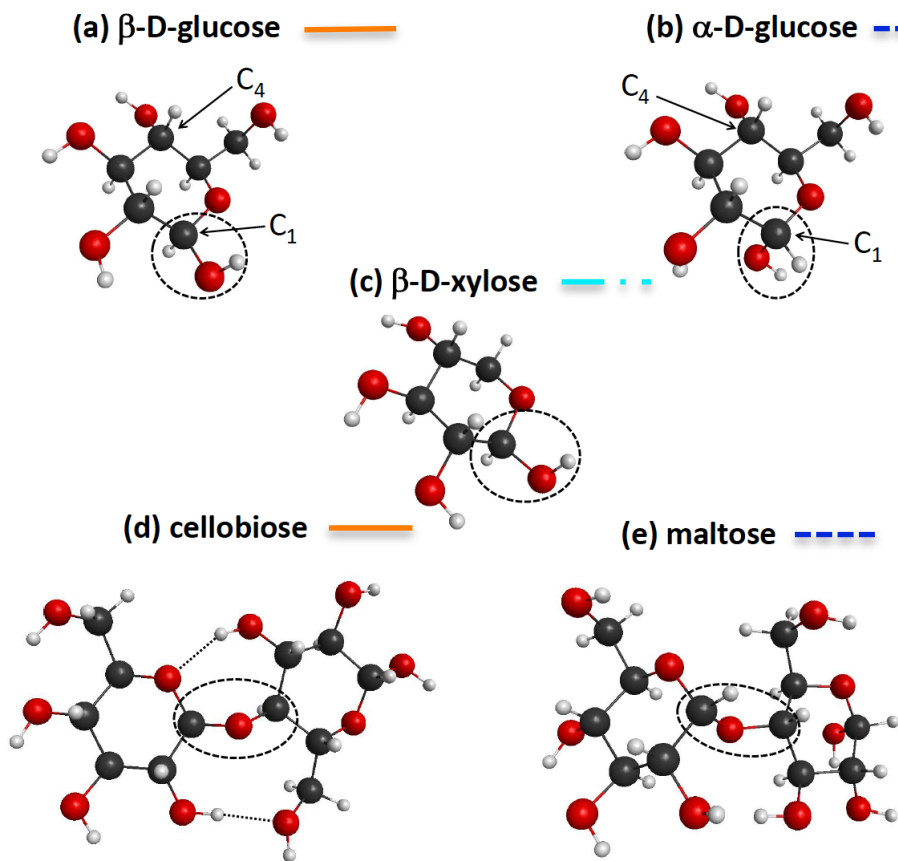


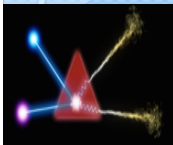
σ^* (LUMO+3)



Low-energy electron scattering by cellulose and Hemicellulose components

Phys. Chem. Chem. Phys. **15**, 1682 (2013).





Theoretical team on electron-scattering of microsolvated molecules



Sylvio Canuto (microsolvation)
Kaline Coutinho (microsolvation)
Márcio T. do N. Varella



Eliane M. de Oliveira (scattering of solvated phenol)
Marco A. P. Lima

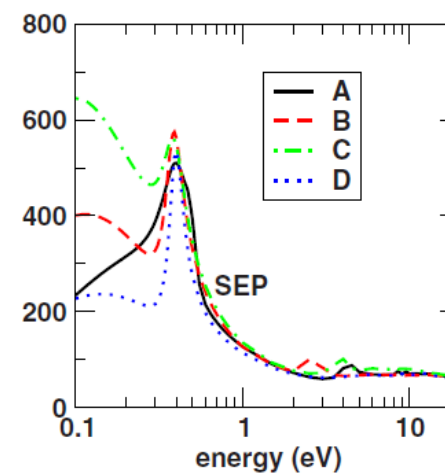
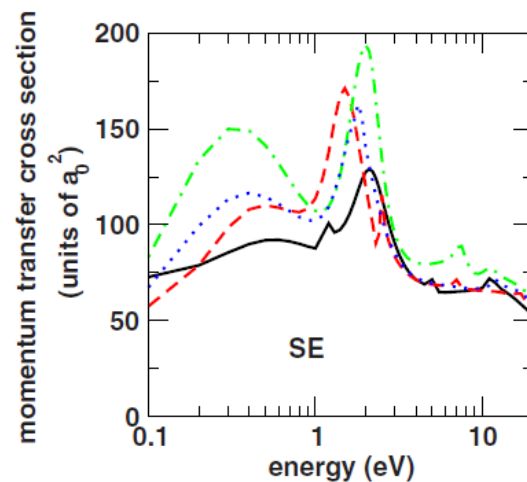
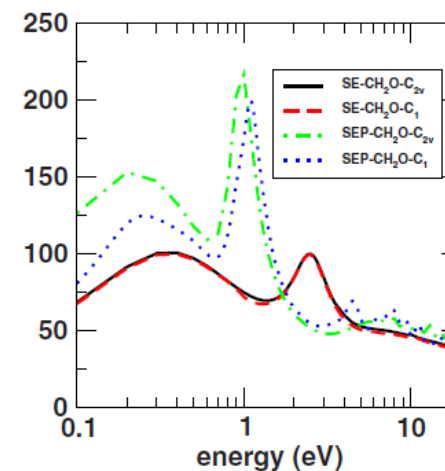
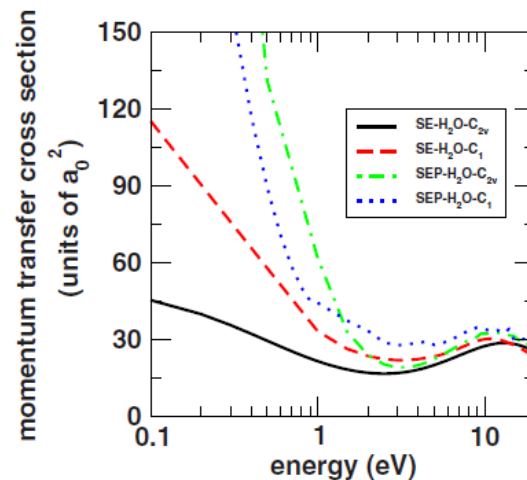
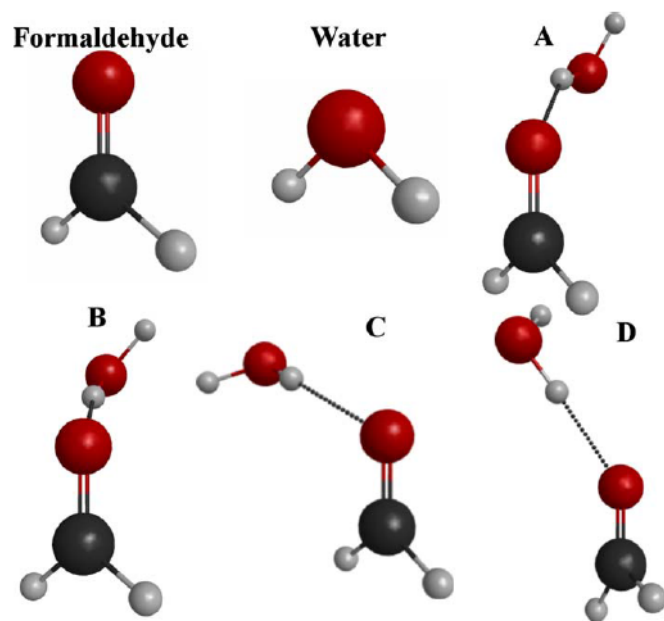


Thiago C. Freitas (his Ph.D. Thesis)
Márcio H. F. Bettega (coordinator)



Electron Collisions with the CH₂O-H₂O complex

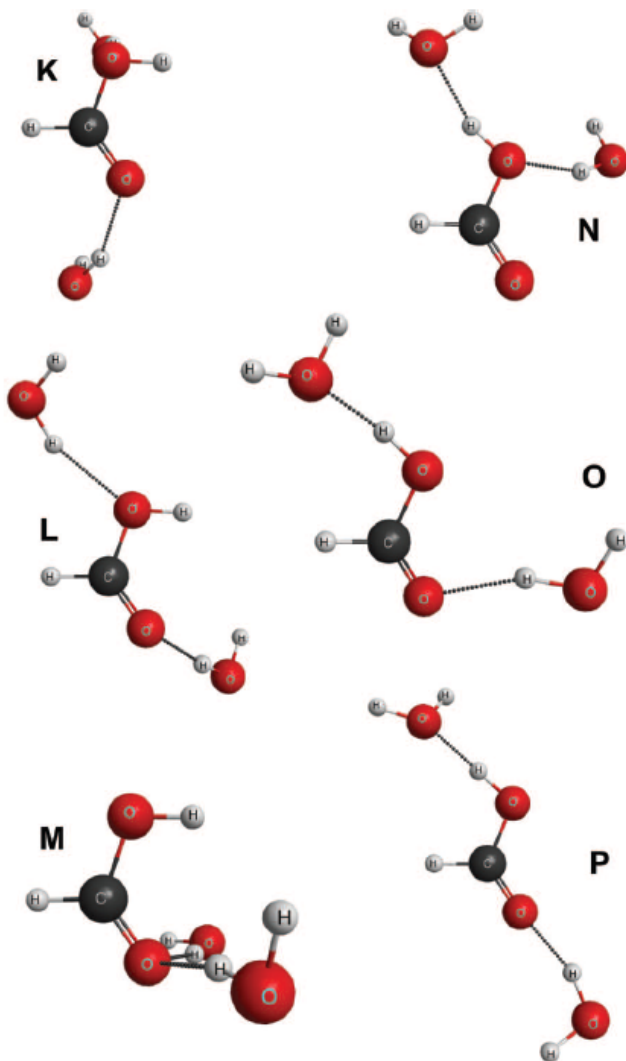
PHYSICAL REVIEW A 80, 062710 (2009)



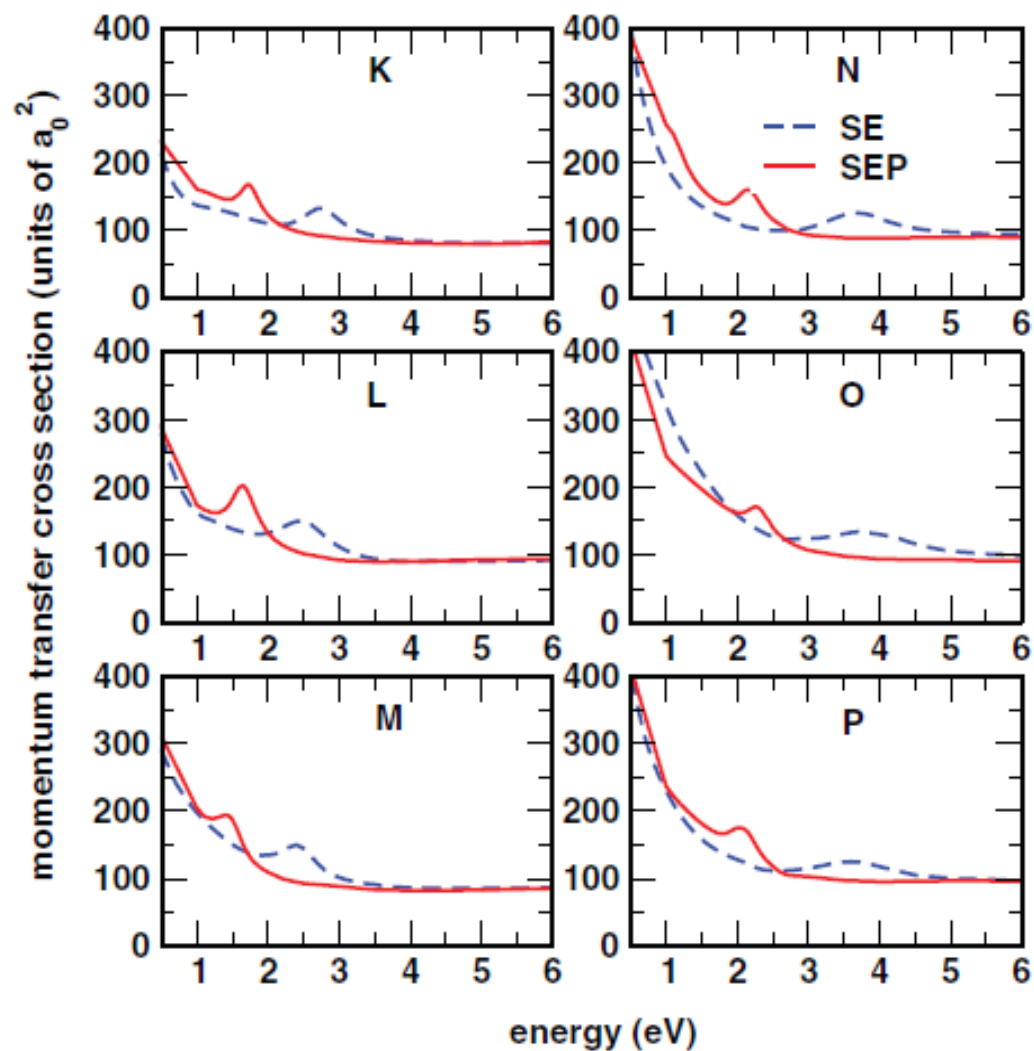


Electron collisions with the HCOOH...(H₂O)_n complexes (n=1, 2) in liquid phase: The influence of microsolvation on the π^* resonance of formic acid

THE JOURNAL OF CHEMICAL PHYSICS 138, 174307 (2013)



π^* shape resonance for HCOOH at around 1.9 eV.

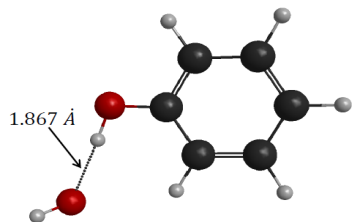




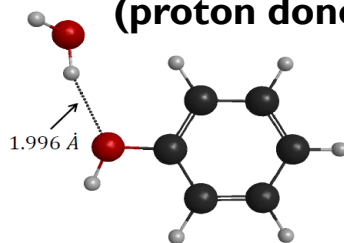
Electron Collisions with Phenol...H₂O

E. M. de Oliveira, T. C. Freitas, K. Coutinho, M. T. do N. Varella, S. Canuto, M. A. P. Lima and M.H.F Bettega, The Journal of Chemical Physics 141, 051105 (2014)

(proton acceptor)

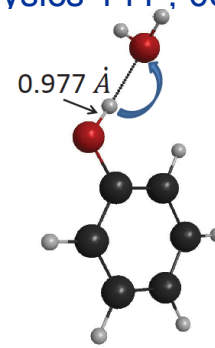


(proton donor)

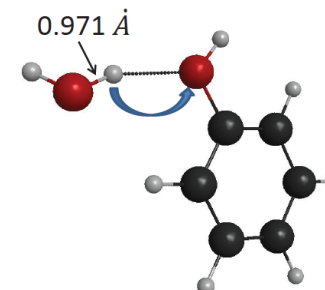


We have studied the microsolvation of Phenol using 4 complexes.

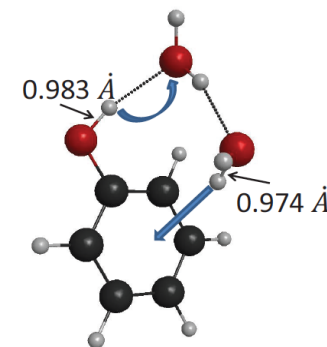
- In Complex A, the water molecule is a proton acceptor.
- In the Complex B, the water is a proton donor.
- Complexes C and D have both situations (one water molecule as acceptor and the other as proton donor).



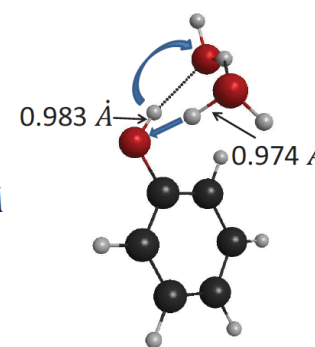
CoA



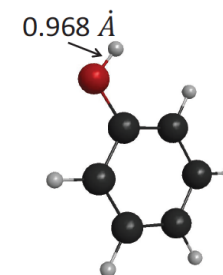
CoB



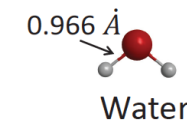
CoC



CoD



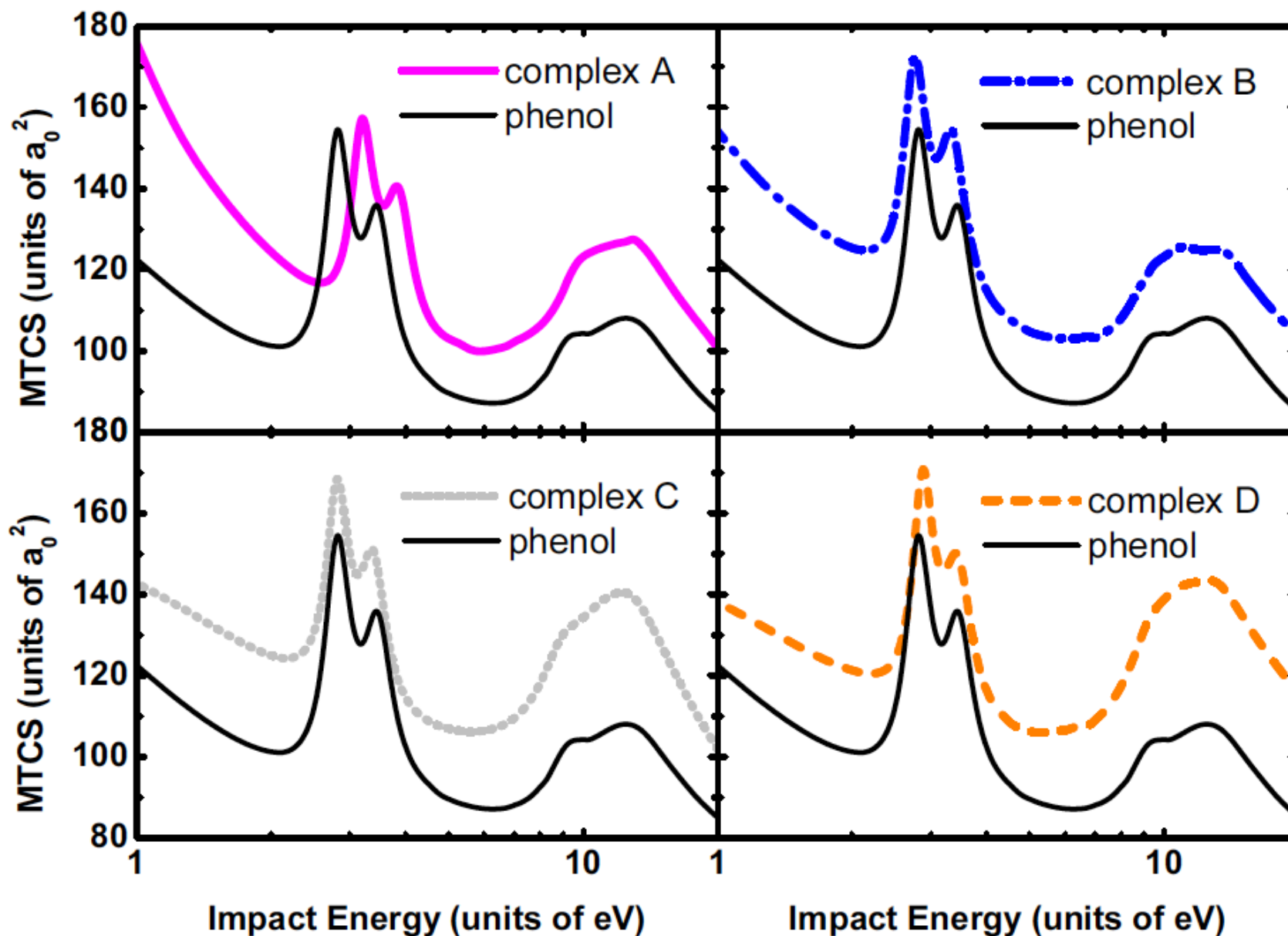
Ph



Water

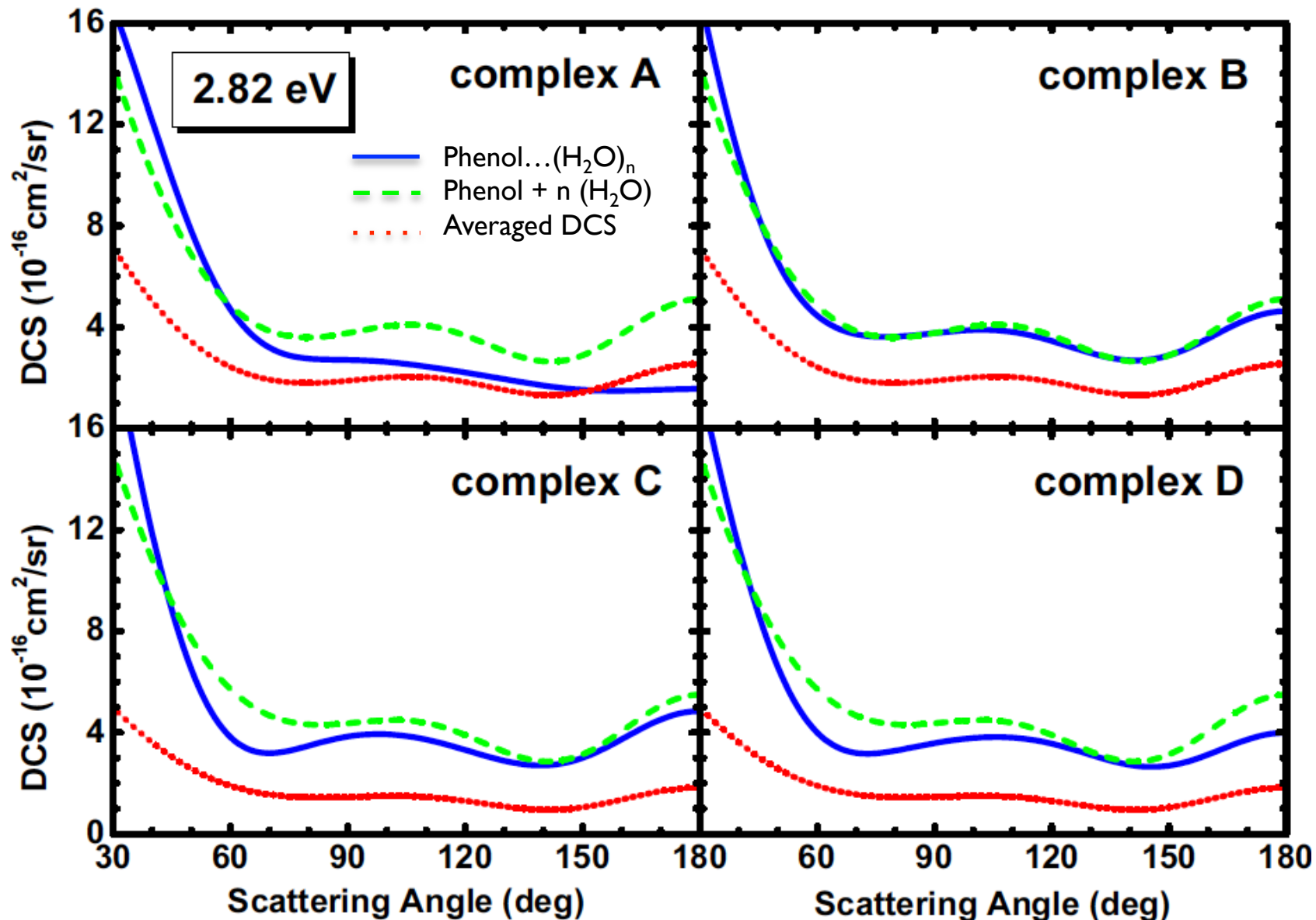


Electron Collisions with Phenol... $(\text{H}_2\text{O})_n$: $n=1,2$



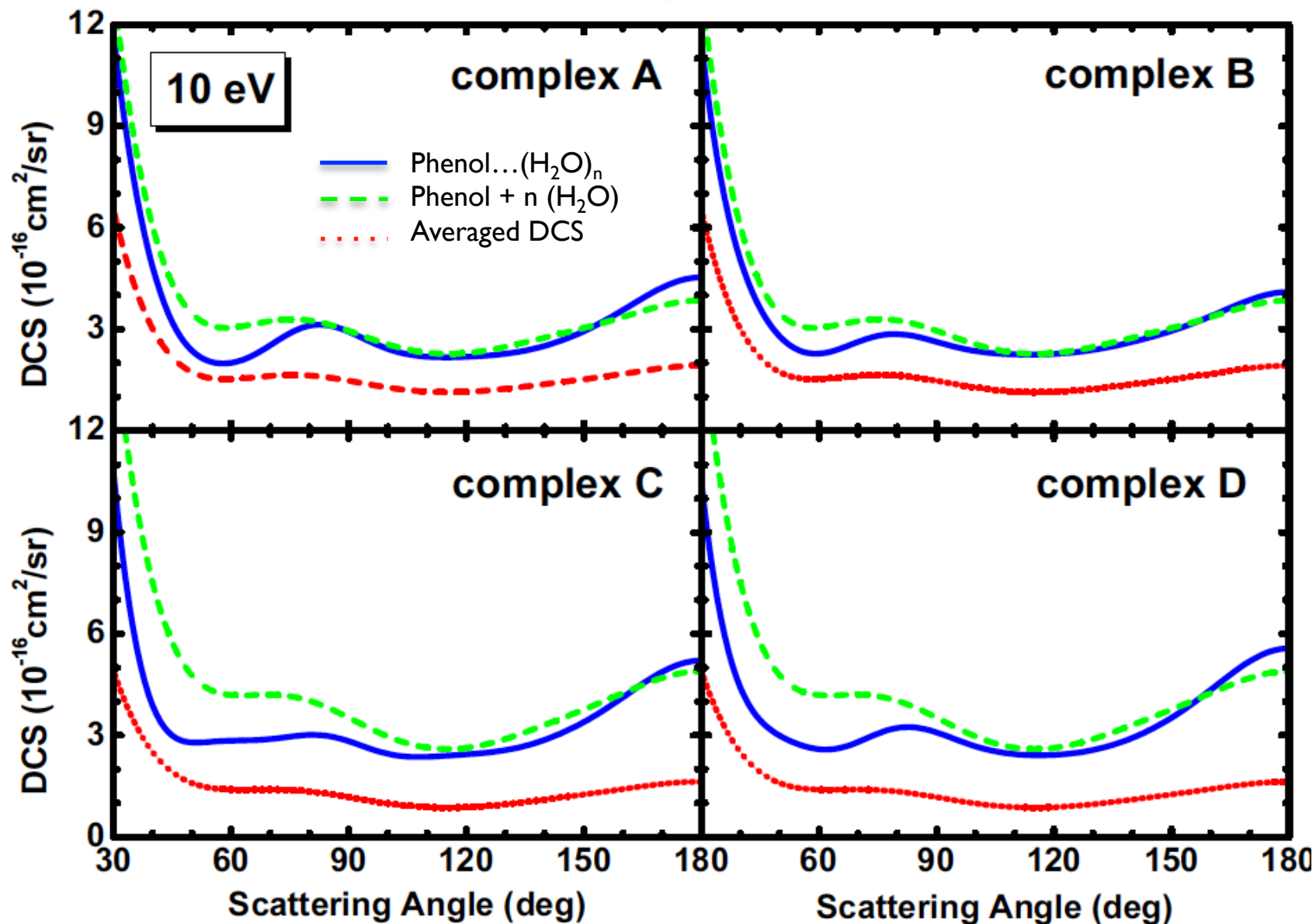


Electron Collisions with Phenol... $(\text{H}_2\text{O})_n$: search for microsolvation signatures in the DCS



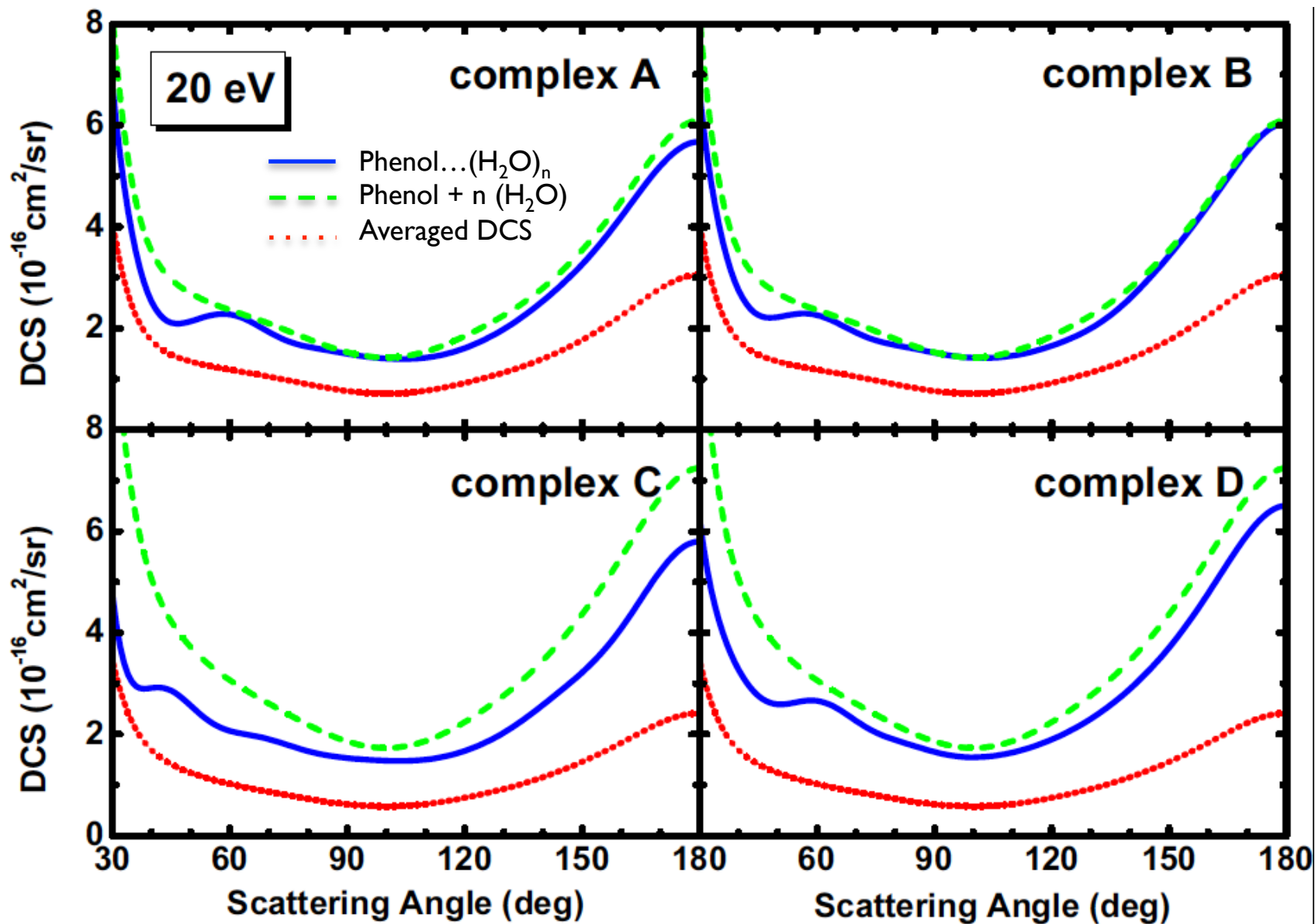


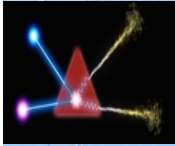
Electron Collisions with Phenol... $(\text{H}_2\text{O})_n$: search for microsolvation signatures in the DCS



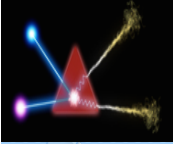


Electron Collisions with Phenol... $(\text{H}_2\text{O})_n$: search for microsolvation signatures in the DCS





ELECTRONIC EXCITATION



Theoretical co-authors



Eliane M. de Oliveira (posdoc)
Marco A. P. Lima



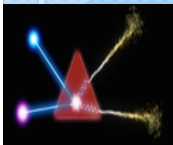
Márcio H. F. Bettega



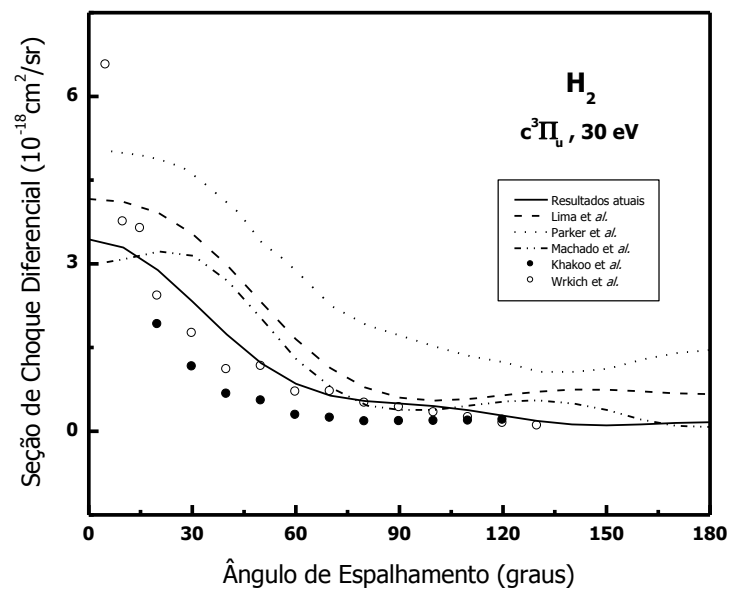
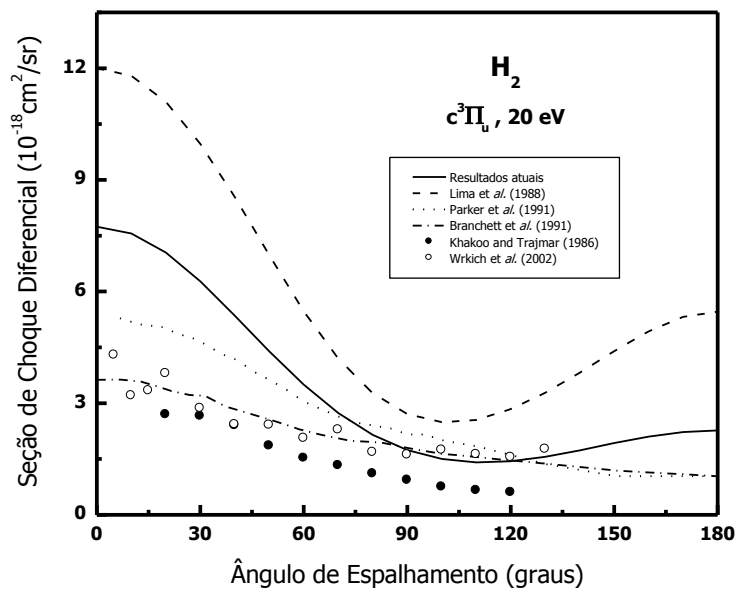
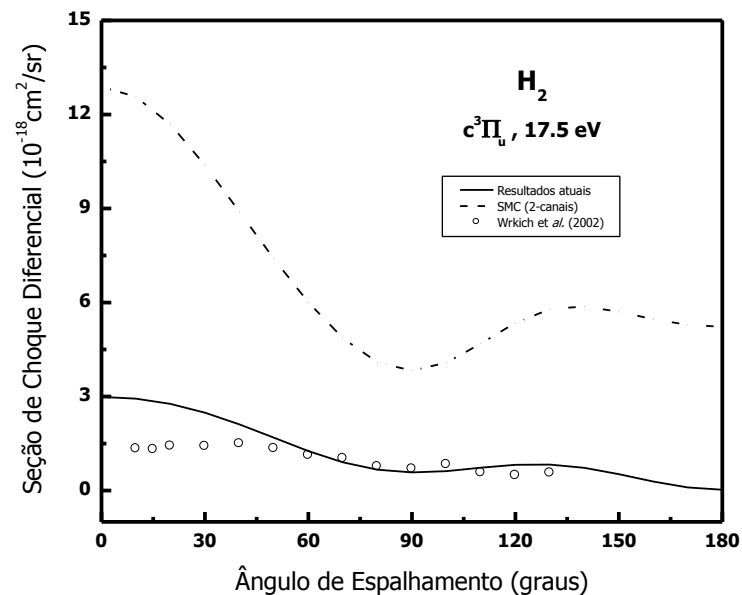
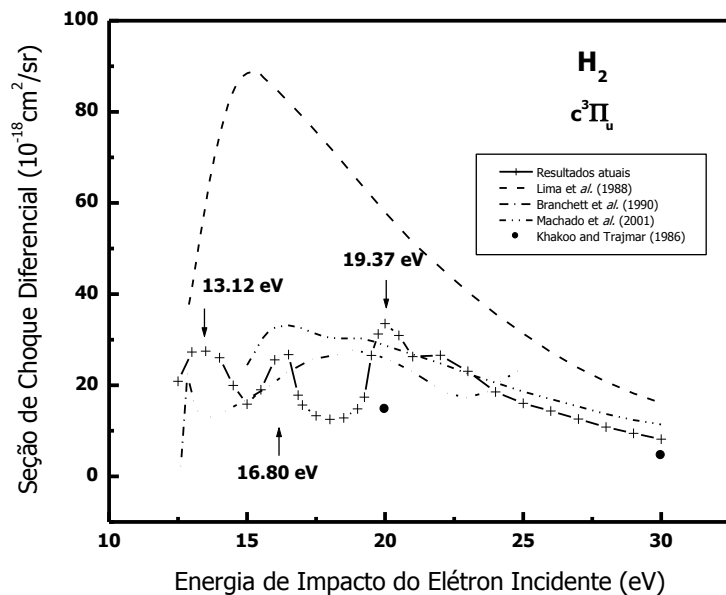
Márcio T. do N. Varella



Romarly F. da Costa (coordinator)



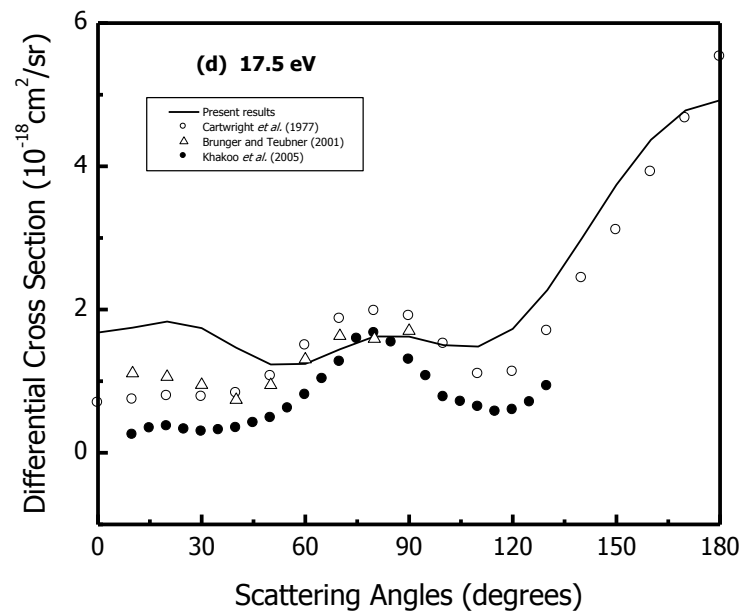
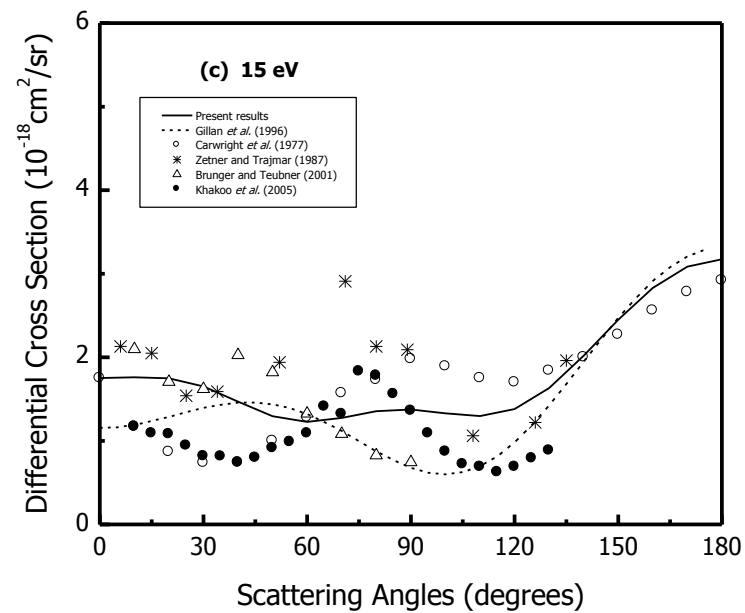
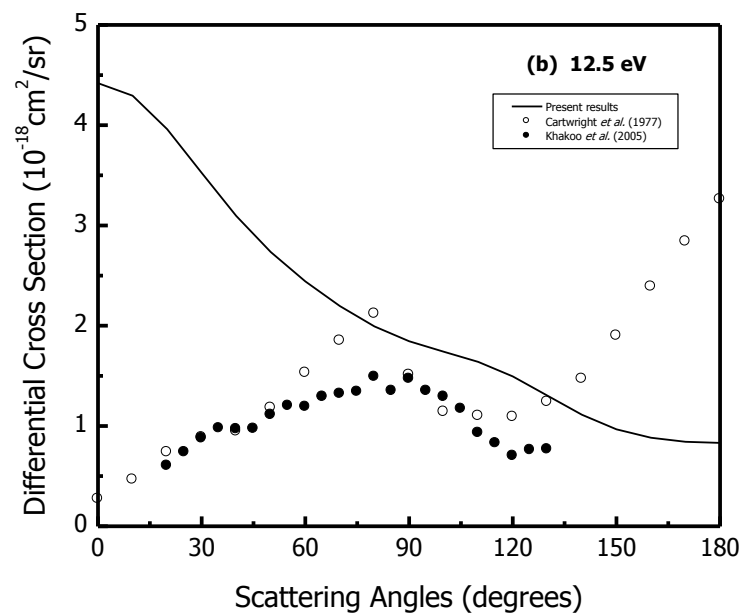
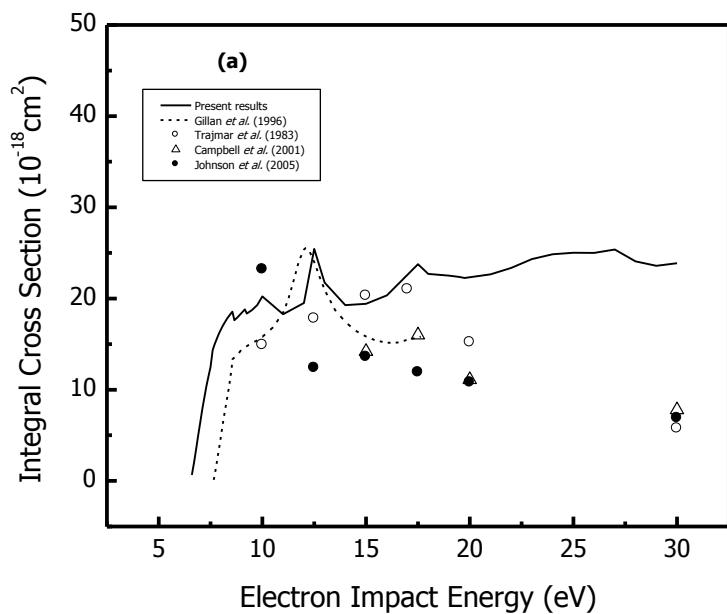
Electronic excitation of H₂ by electron impact





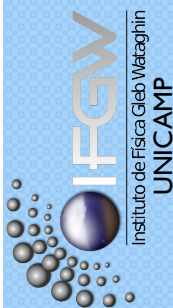
R.F. da Costa and M.A.P. Lima,
Phys. Rev.A **75**, 022705 (2007)

Electronic transition $X \ ^1\Sigma_g \rightarrow A \ ^3\Sigma_u^+$ of N_2 by electron impact

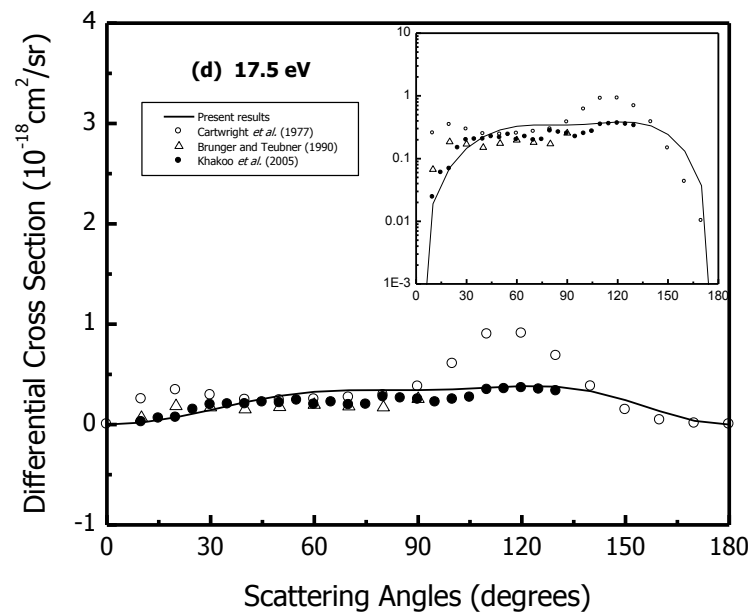
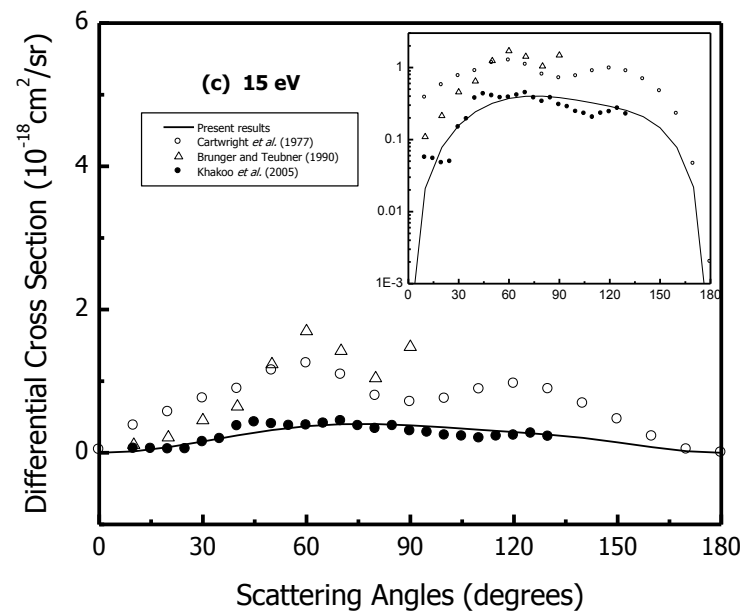
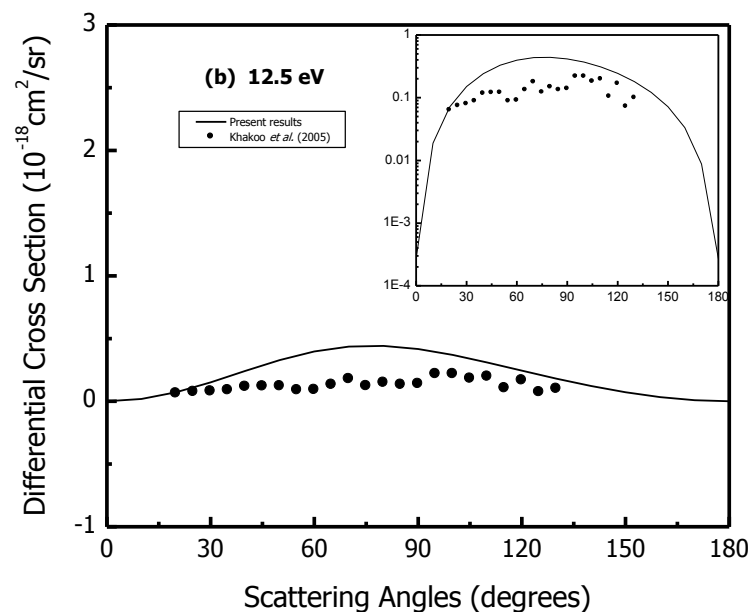
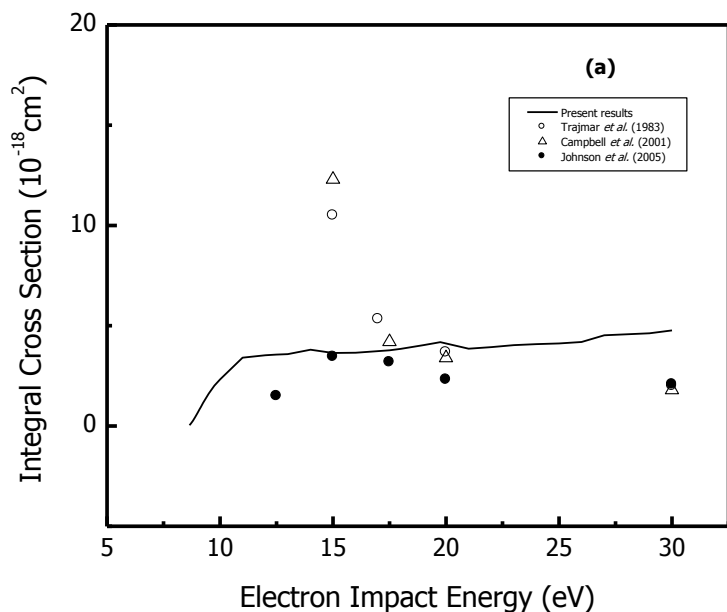


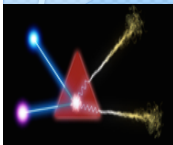


R.F. da Costa and M.A. P. Lima,
Phys. Rev.A **75**, 022705 (2007)



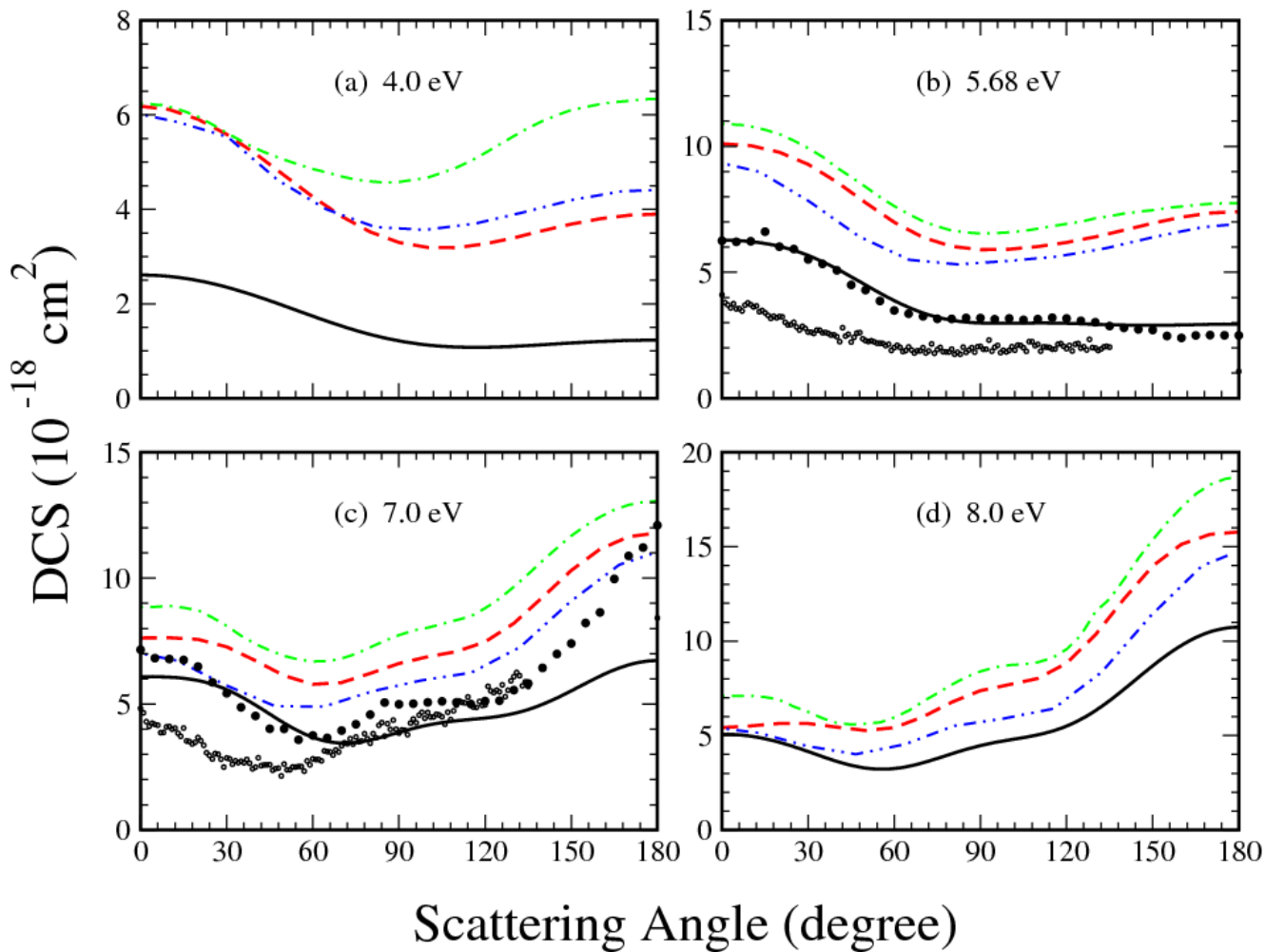
Electronic transition $X^1\Sigma_g \rightarrow a'^1\Sigma_u^-$ of N_2 by electron impact



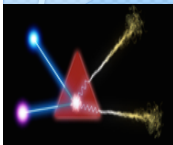


Electronic excitation of \tilde{a}^3B_{1u} state of C_2H_4 by electron impact

R. F. da Costa, M. H. F. Bettega, and M. A. P. Lima,
Phys. Rev. A **77**, 042723 (2008).

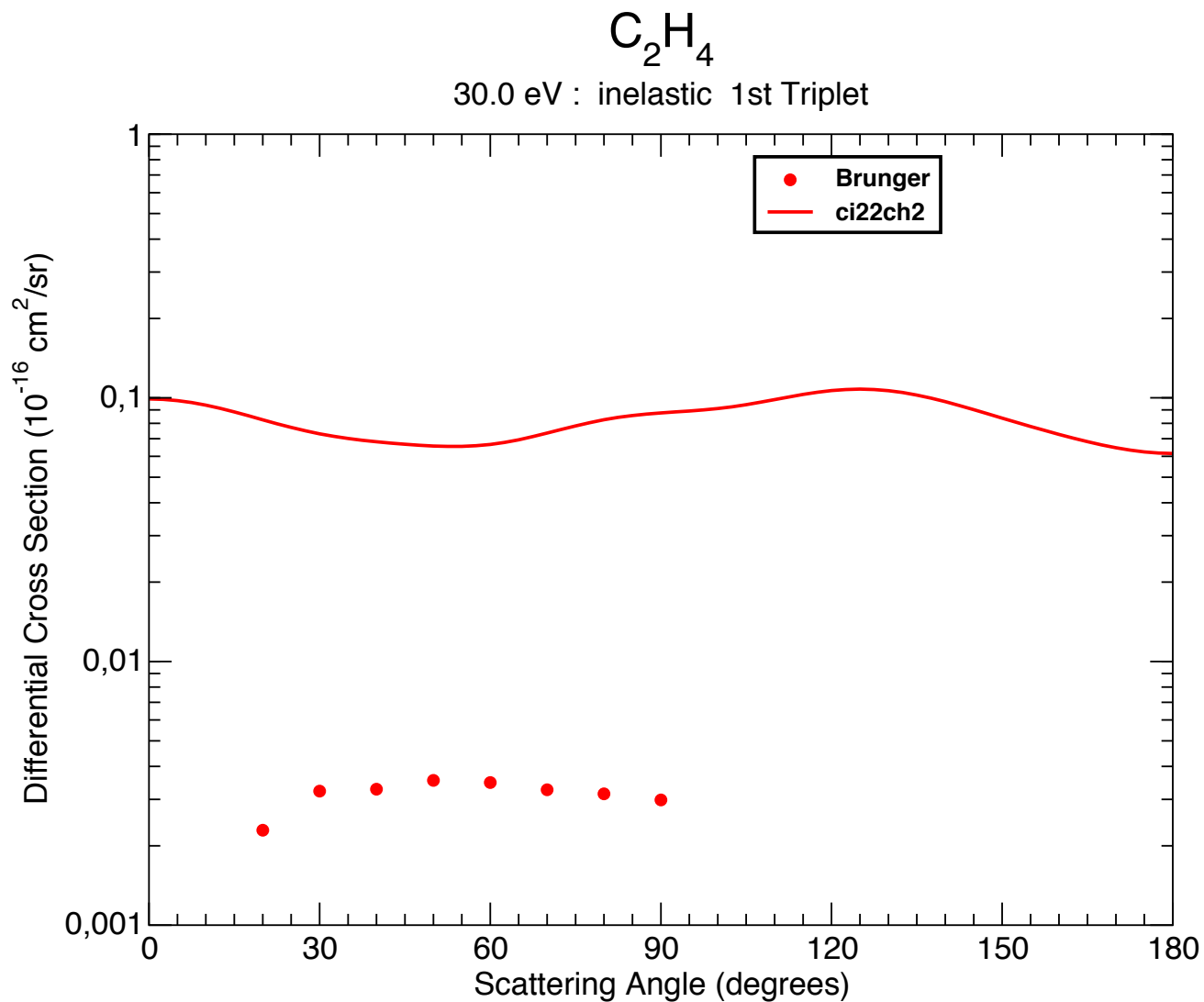


Color lines are
Close-coupling
Calculations
and
bullets are
M.Allan's
data

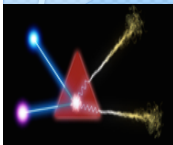


R. F. da Costa, M. H. F. Bettega, M. T. do N. Varella, E. M. de Oliveira and M. A. P. Lima,
Phys. Rev. A **90**, 052707 (2014).

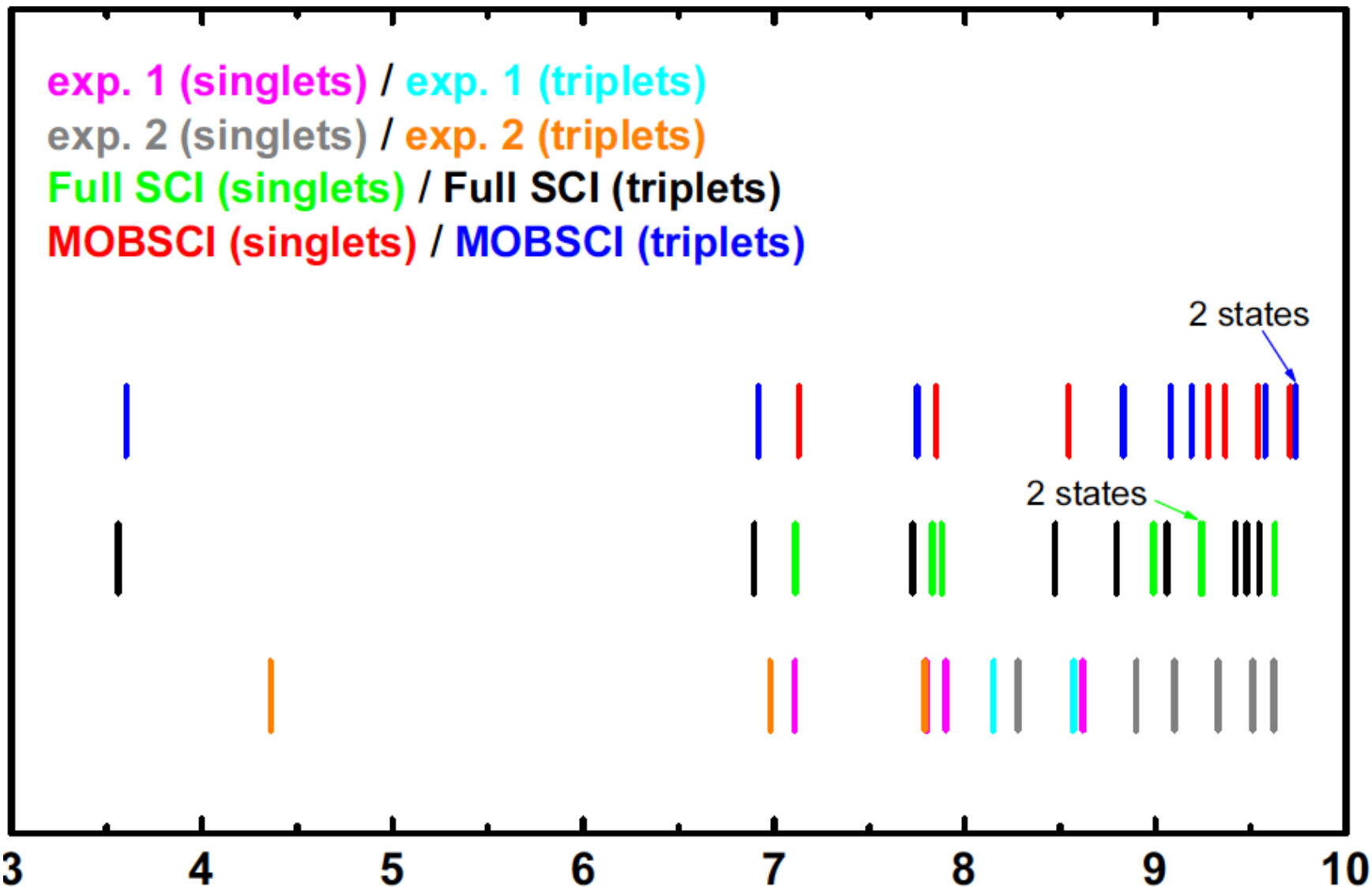
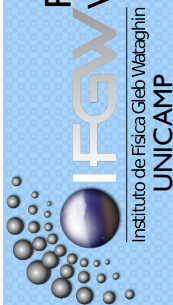
Surprisingly, at higher energies the agreement is not good!!



Multichannel effects on the $e^-C_2H_4$ scattering



R. F. da Costa, M. H. F. Bettega, M. T. do N. Varella, E. M. de Oliveira and M. A. P. Lima, *Phys. Rev. A* **90**, 052707 (2014).



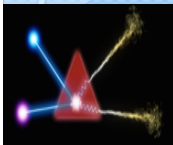


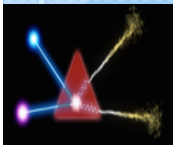
TABLE I. Calculated and experimental excitation energies for ethylene. Up to 20, 30 and 50 eV, the FSCI spectrum is composed by 138, 260 and 402 electronically excited states, respectively. The MOBSCI calculations at these energies were performed with 45 excited states, where 17 of them are physical excited singlets or triplets states and the others are pseudostates.

	Energy (eV)		
	FSCI	MOBSCI	Expt.
TRIPLET	3.56	3.60	4.36 ^a
	6.90	6.92	6.98 ^a
	7.73	7.75	7.79 ^a
	8.48	8.83	8.15 ^b
	8.80	9.08	8.57 ^b
	9.06	9.19	
	9.42	9.58	
	9.48	9.73	
	9.54	9.74	
SINGLET	7.11	7.13	7.11 ^b
	7.83	7.85	7.80 ^b
	7.88	8.55	7.90 ^b
	8.99	9.28	8.28 ^a
	9.24	9.37	8.62 ^b
	9.25	9.54	8.90 ^{a,b}
	9.63	9.71	9.10 ^a
			9.33 ^a
			9.51 ^{a,c}
		9.62 ^a	

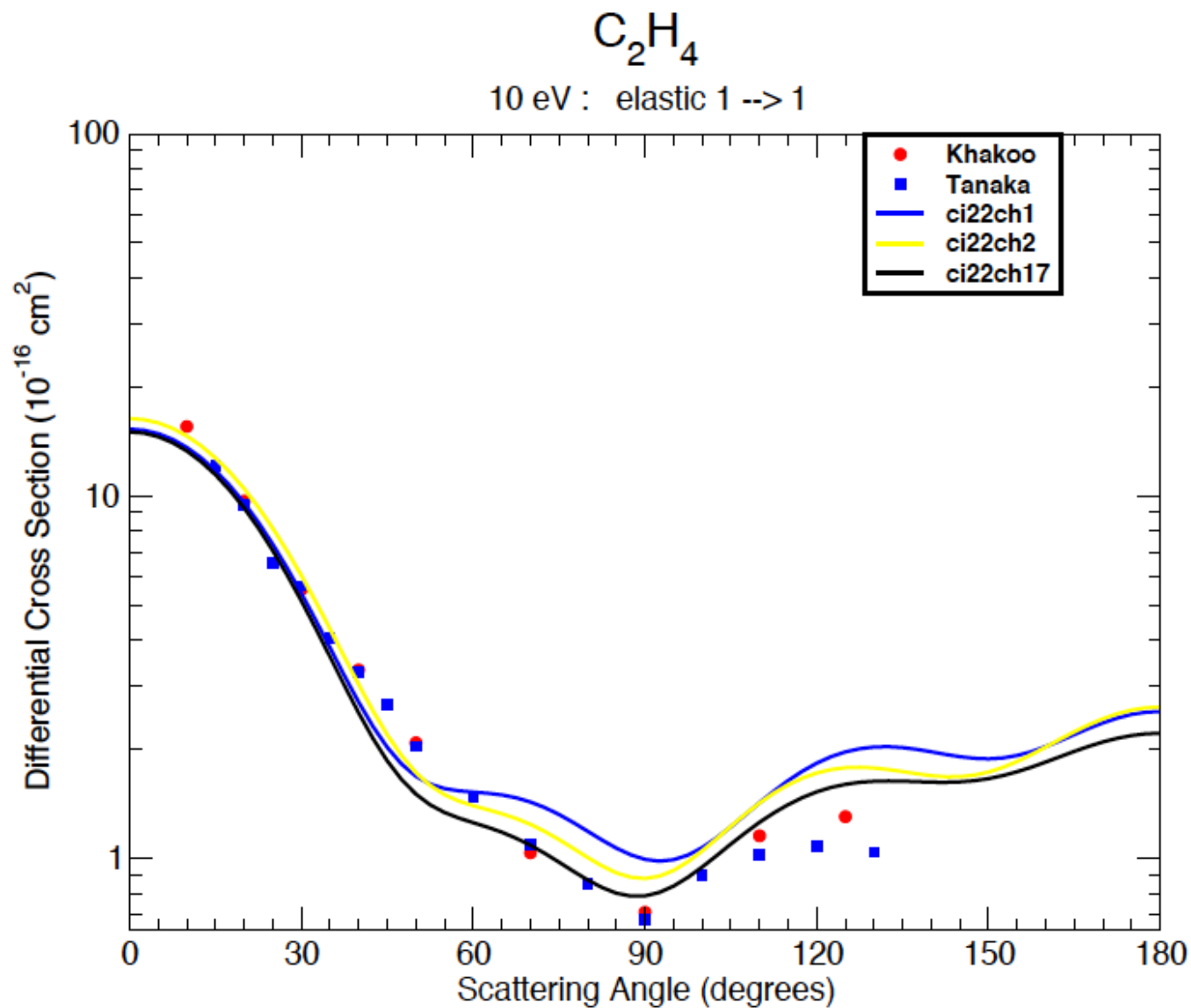
^a Experimental data from Ballard *et al.* [24].

^b Experimental data from Do *et al.* [16].

^c For this energy were found two states.



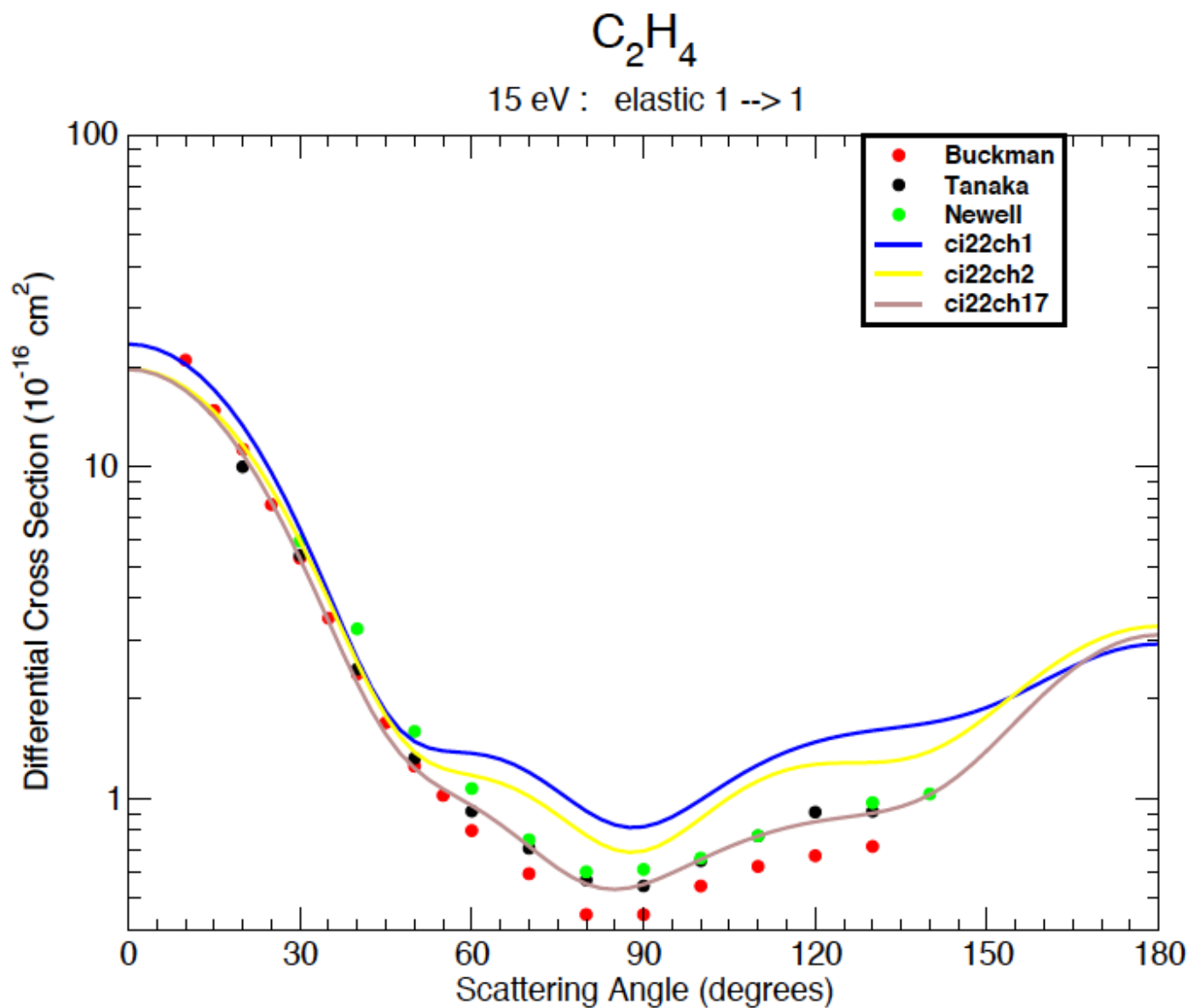
Multichannel effects on the elastic process of $e^-C_2H_4$ scattering

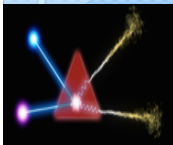




R. F. da Costa, M. H. F. Bettega, M. T. do N. Varella, E. M. de Oliveira and M. A. P. Lima,
Phys. Rev. A **90**, 052707 (2014).

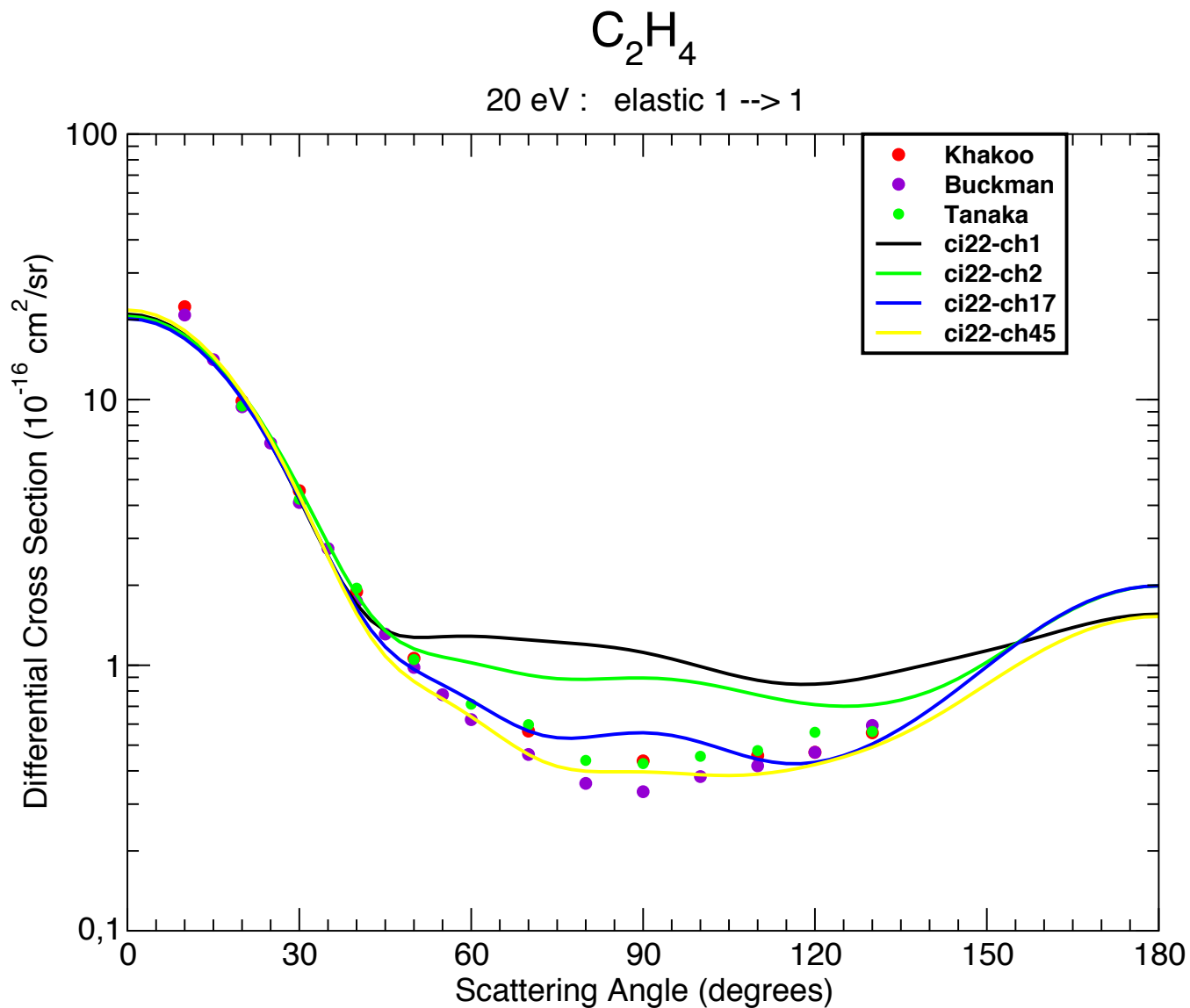
Multichannel effects on the elastic process of $e^-C_2H_4$ scattering

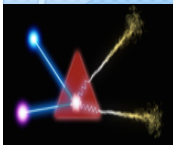




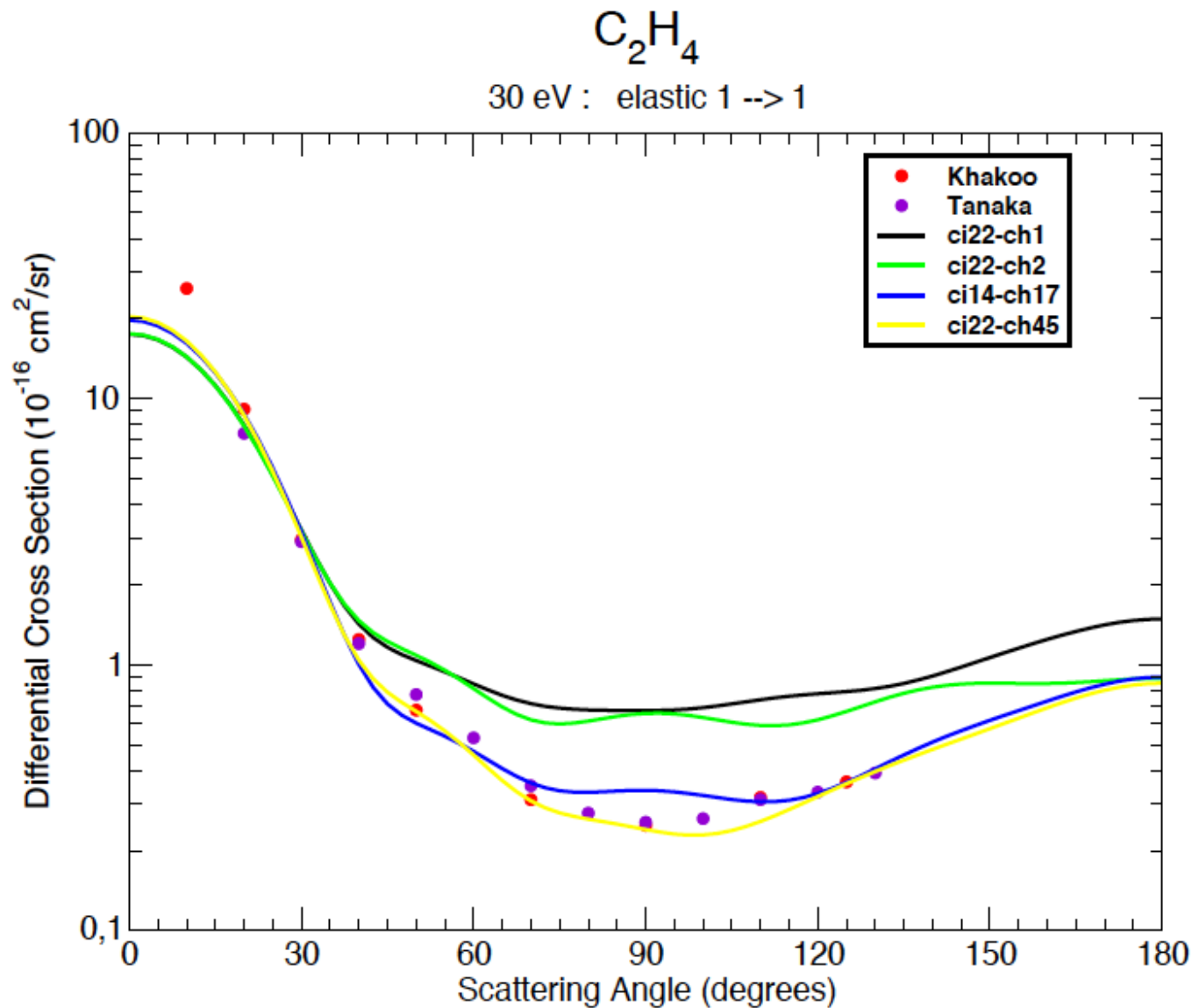
R. F. da Costa, M. H. F. Bettega, M. T. do N. Varella, E. M. de Oliveira and M. A. P. Lima,
Phys. Rev. A **90**, 052707 (2014).

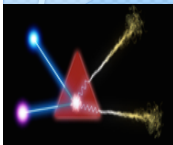
Multichannel effects on the elastic process of e⁻-C₂H₄ scattering





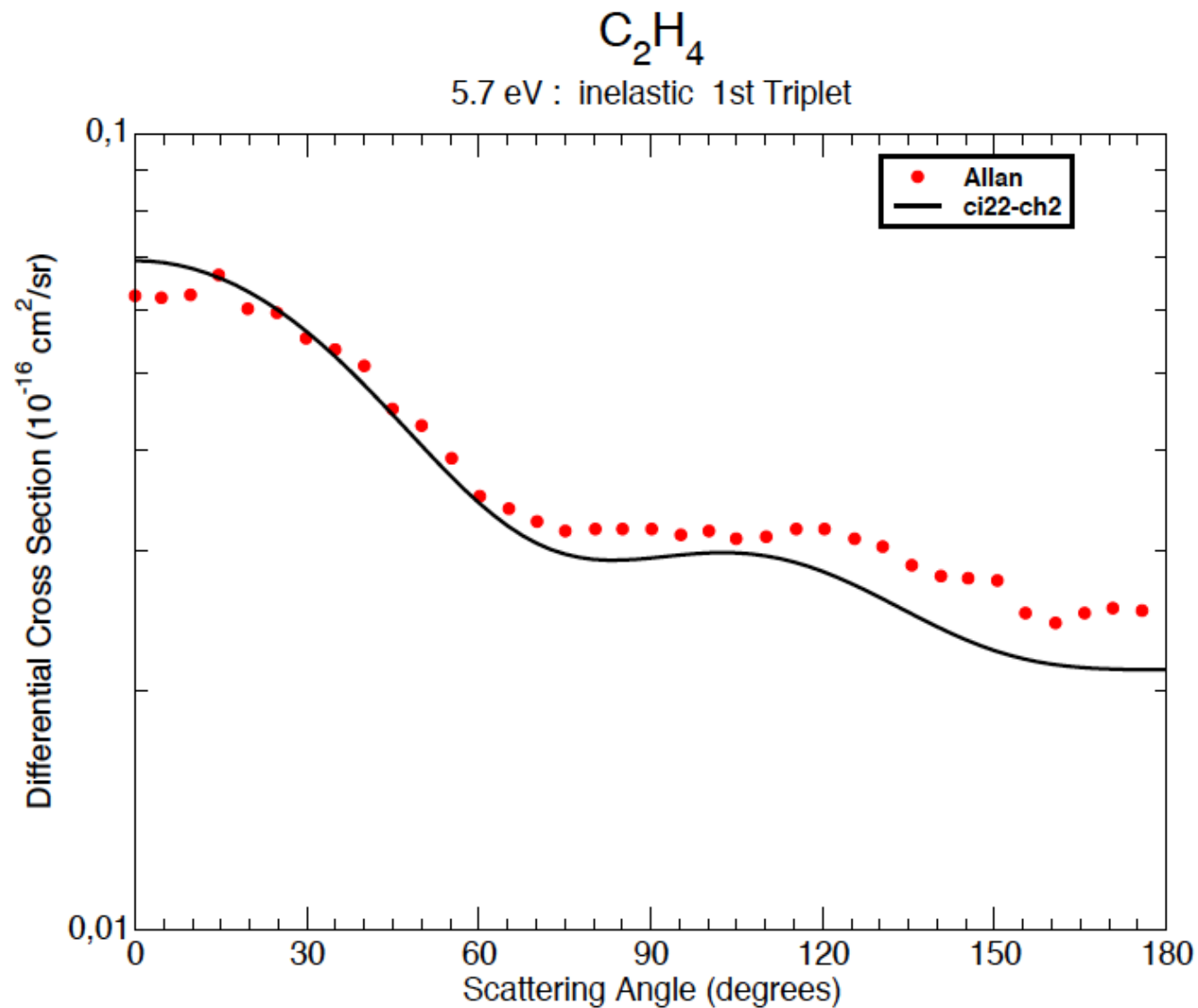
Multichannel effects on the elastic process of $e^-C_2H_4$ scattering

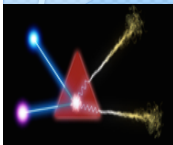




R. F. da Costa, M. H. F. Bettega, M. T. do N. Varella, E. M. de Oliveira and M. A. P. Lima,
Phys. Rev. A **90**, 052707 (2014).

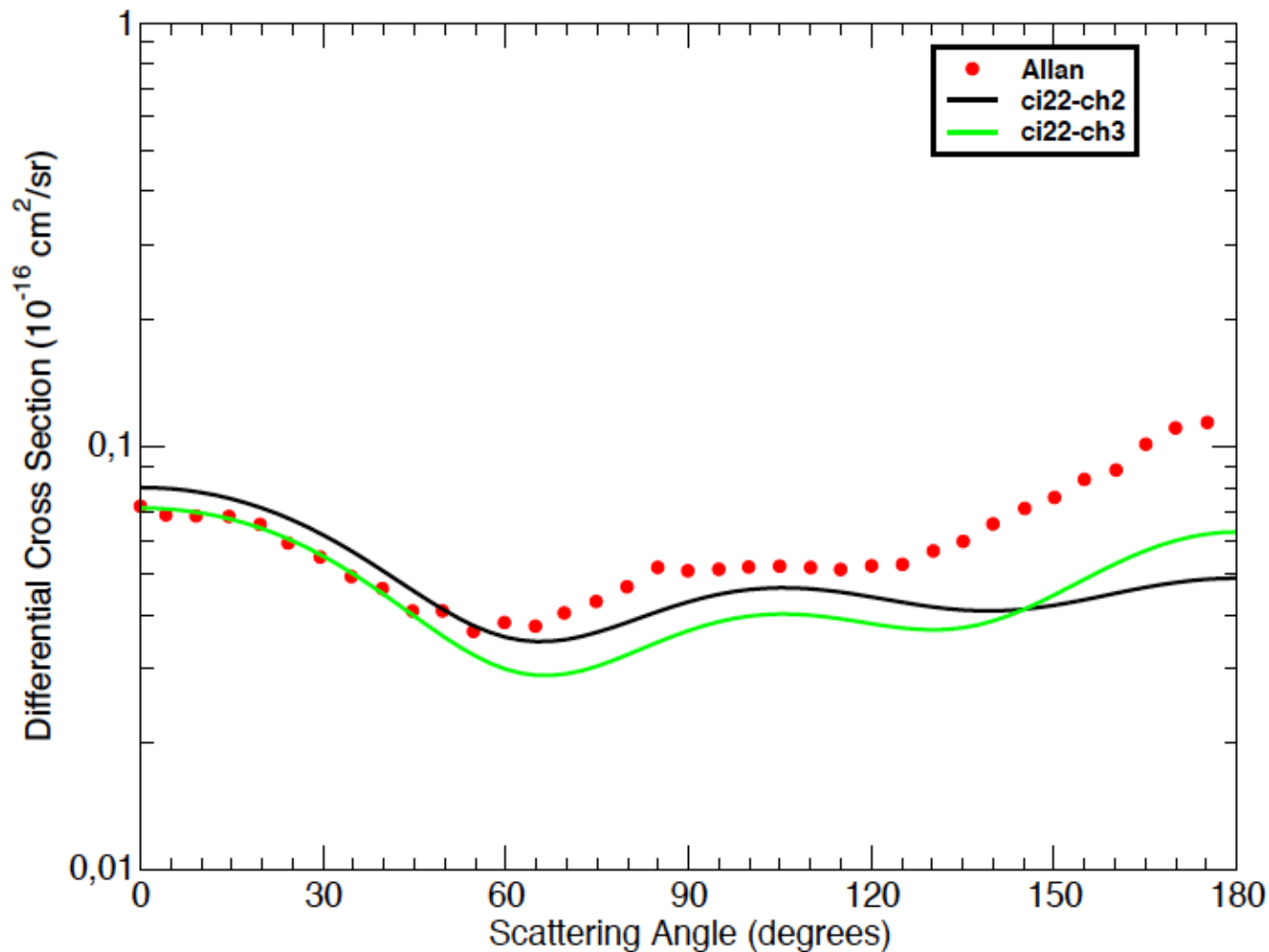
Multichannel effects on the Electronic excitation of \tilde{a}^3B_{1u} state of C_2H_4 by electron impact



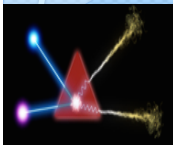


Multichannel effects on the Electronic excitation of \tilde{a}^3B_{1u} state of C_2H_4 by electron impact

C_2H_4
7.0 eV : inelastic 1st Triplet

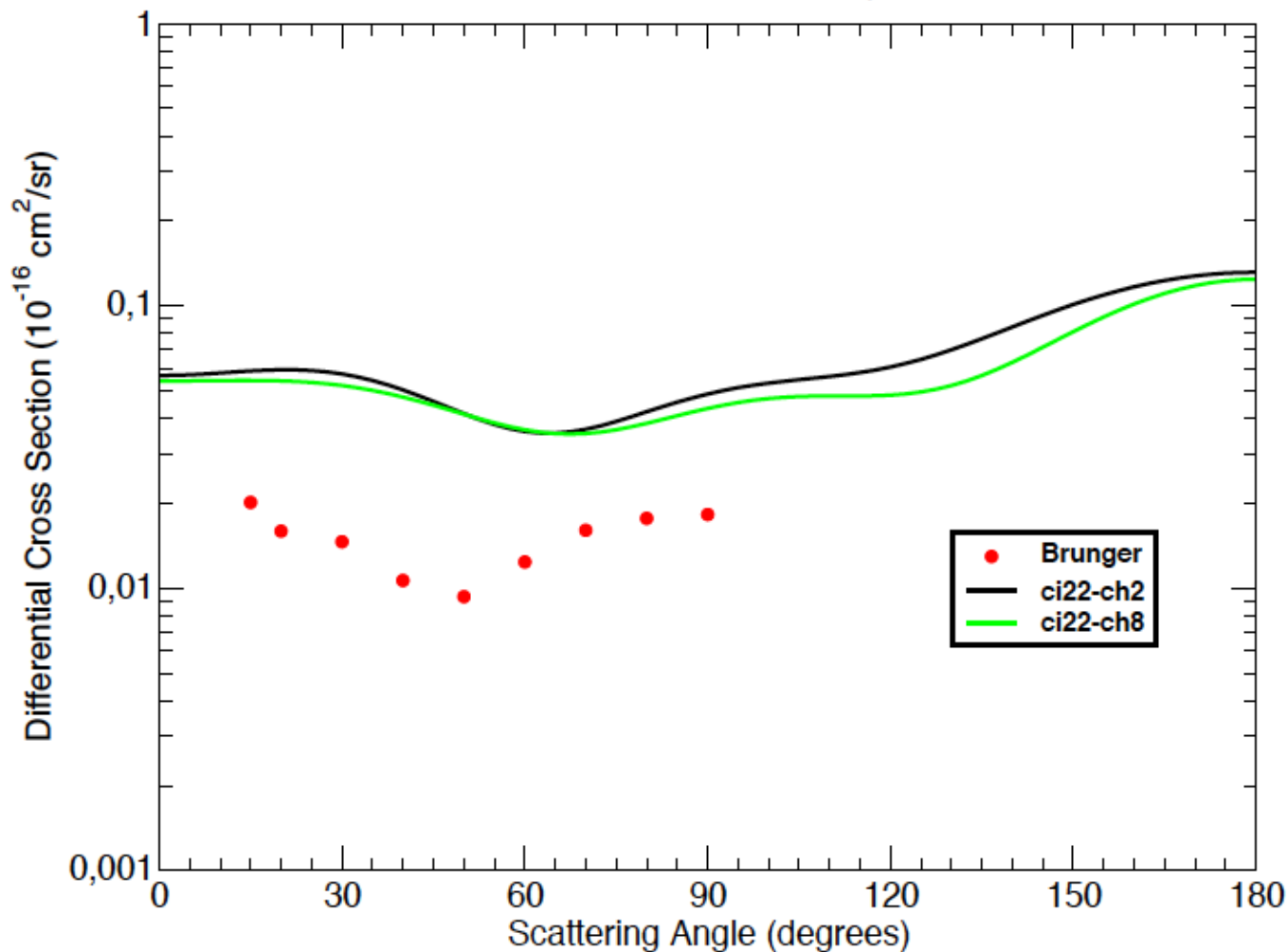


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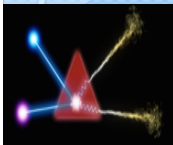


Multichannel effects on the Electronic excitation of \tilde{a}^3B_{1u} state of C_2H_4 by electron impact

C_2H_4
9.0 eV : inelastic 1st Triplet

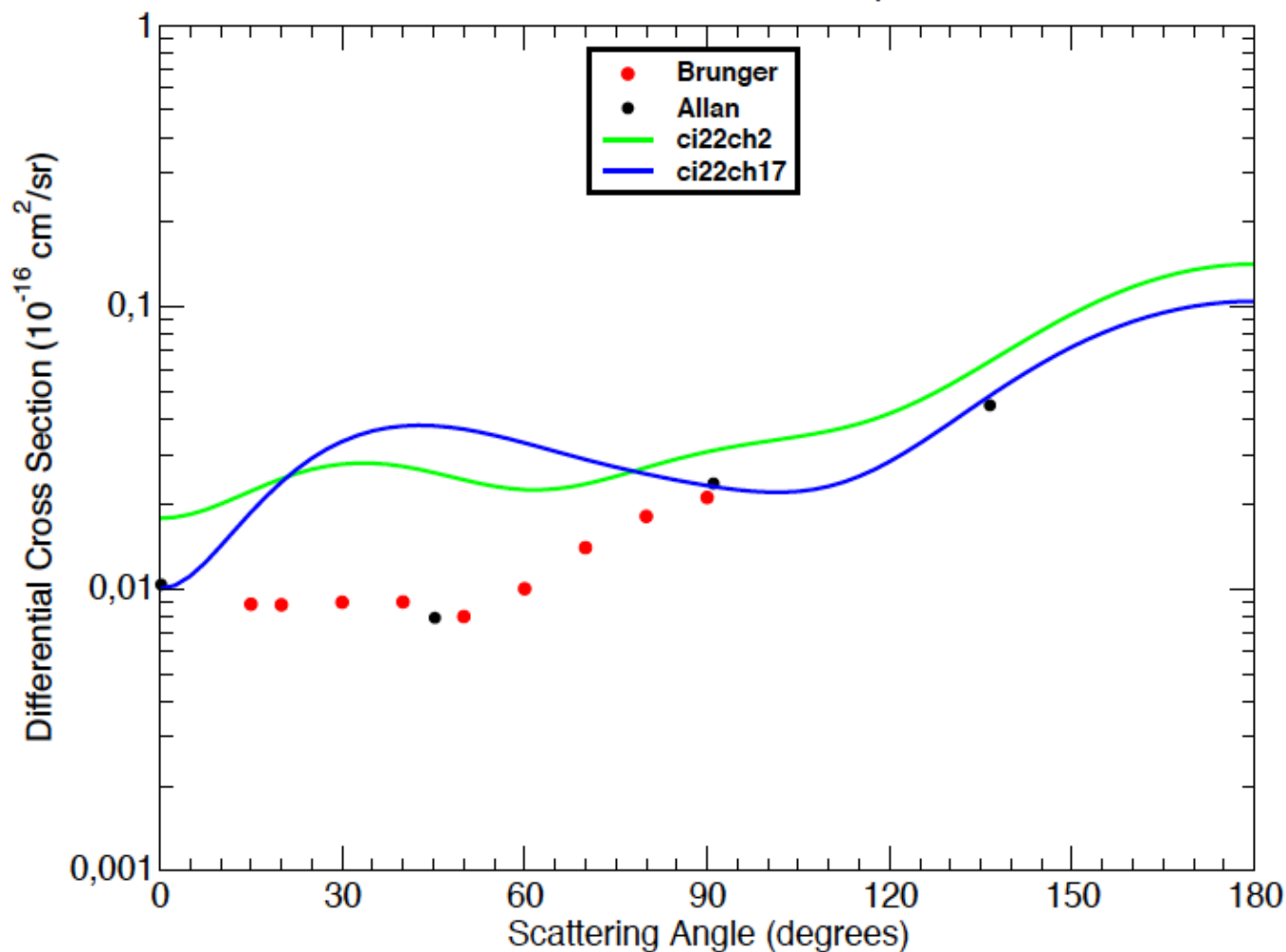


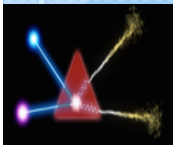
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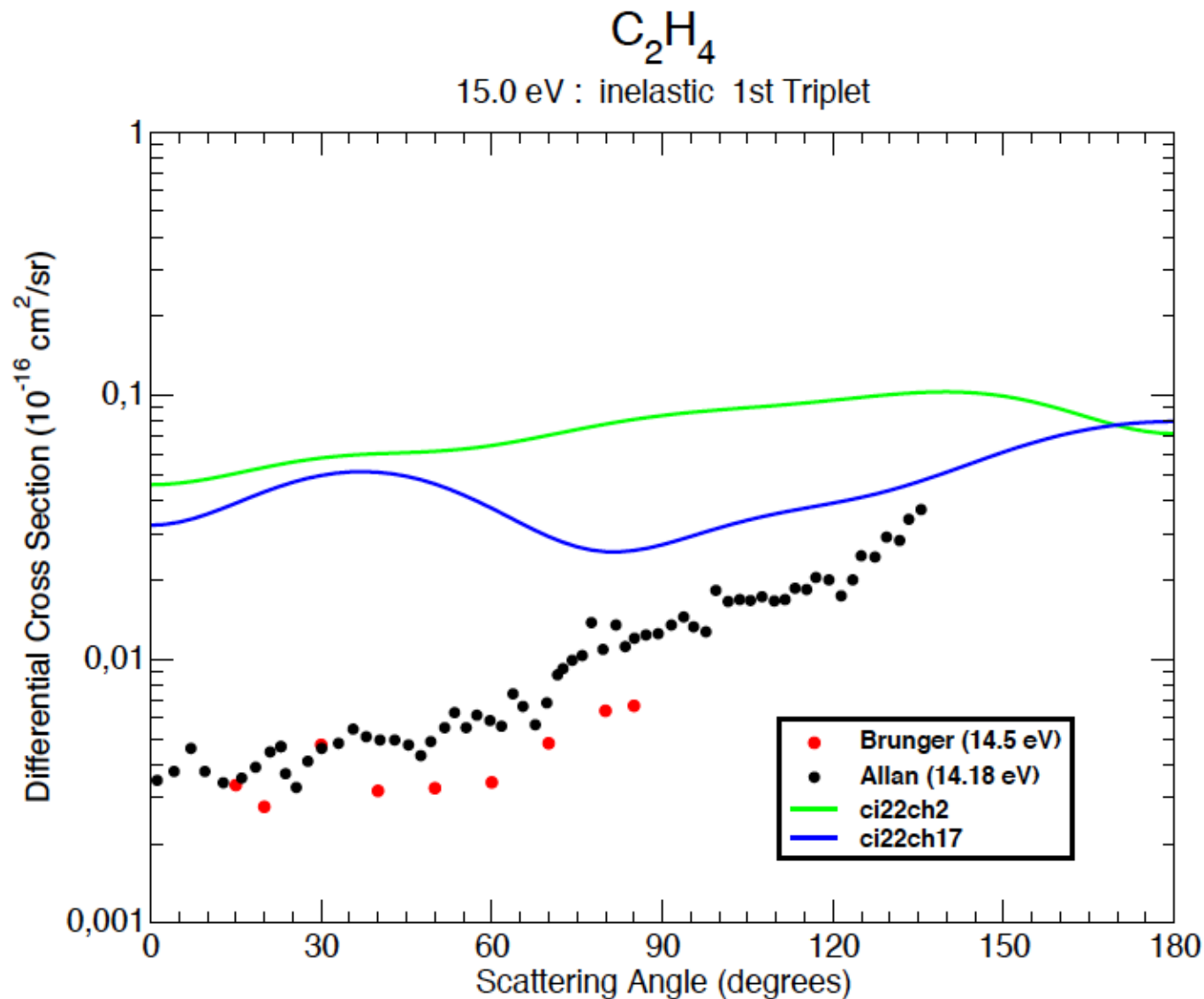
C_2H_4
11.0 eV : inelastic 1st Triplet

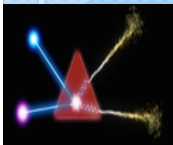




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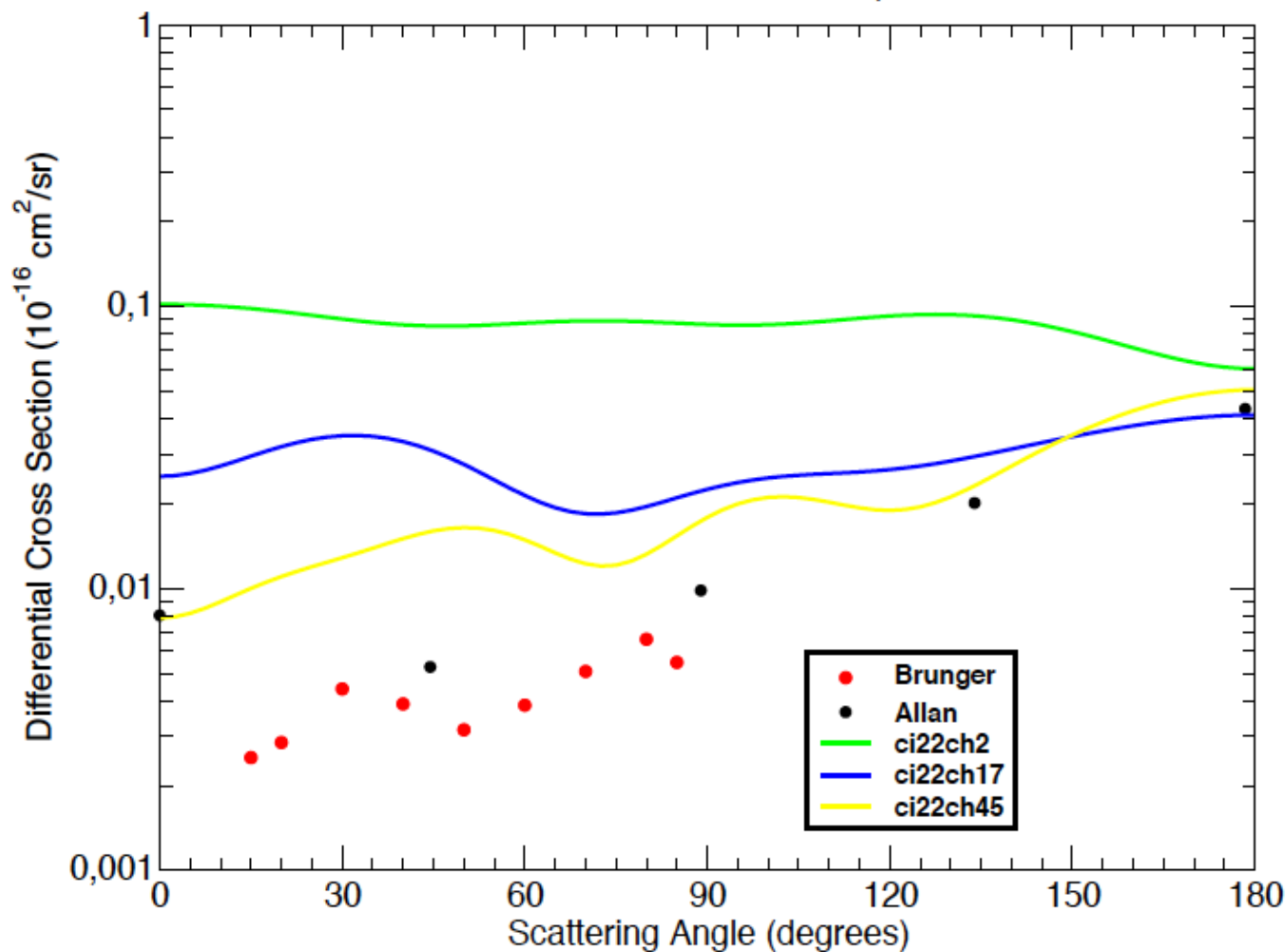
Multichannel effects on the Electronic excitation of \tilde{a}^3B_{1u} state of C_2H_4 by electron impact

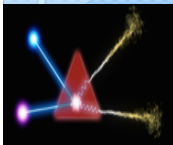




Multichannel effects on the Electronic excitation of \tilde{a}^3B_{1u} state of C_2H_4 by electron impact

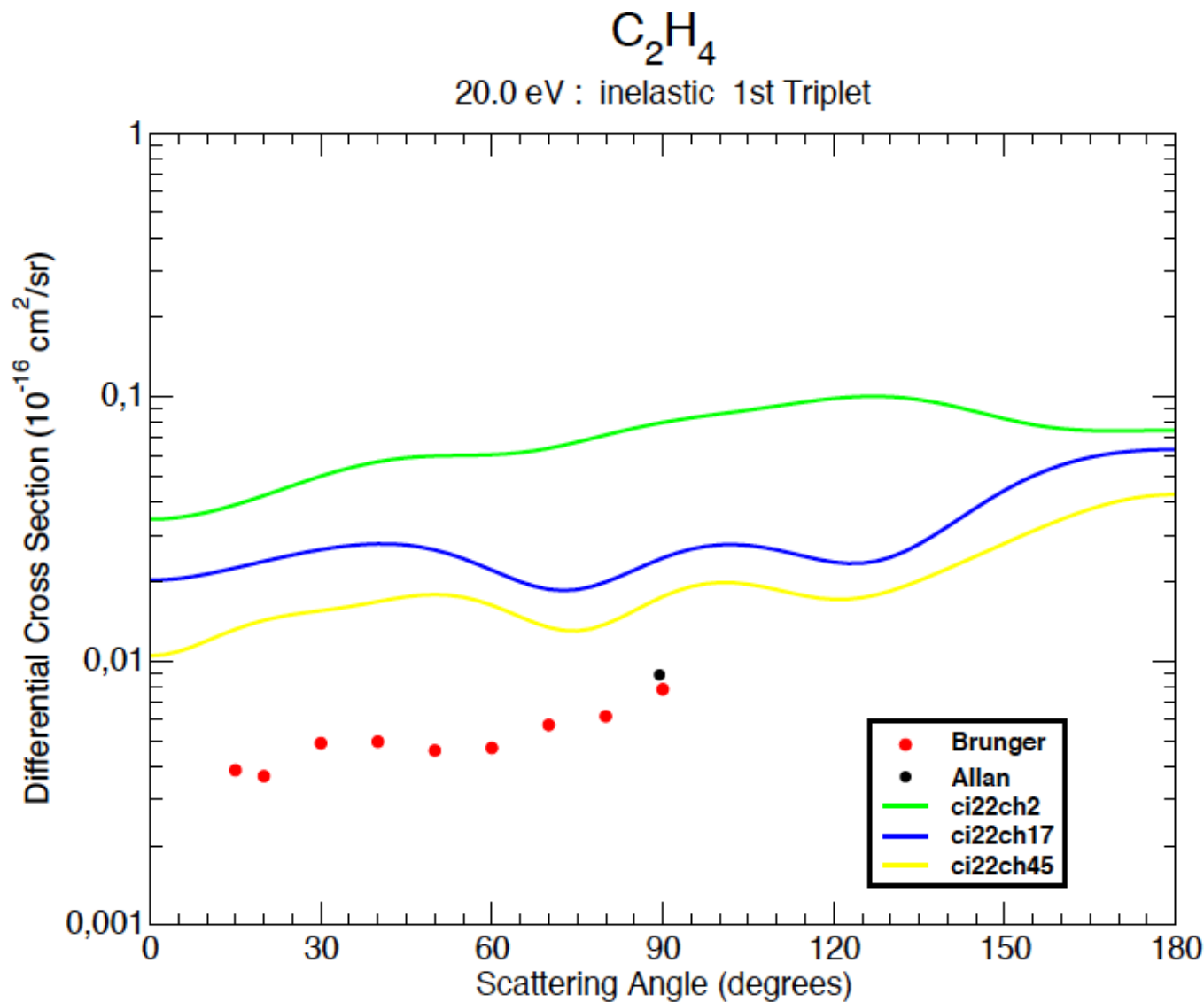
C_2H_4
18.0 eV : inelastic 1st Triplet

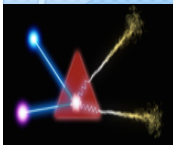




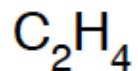
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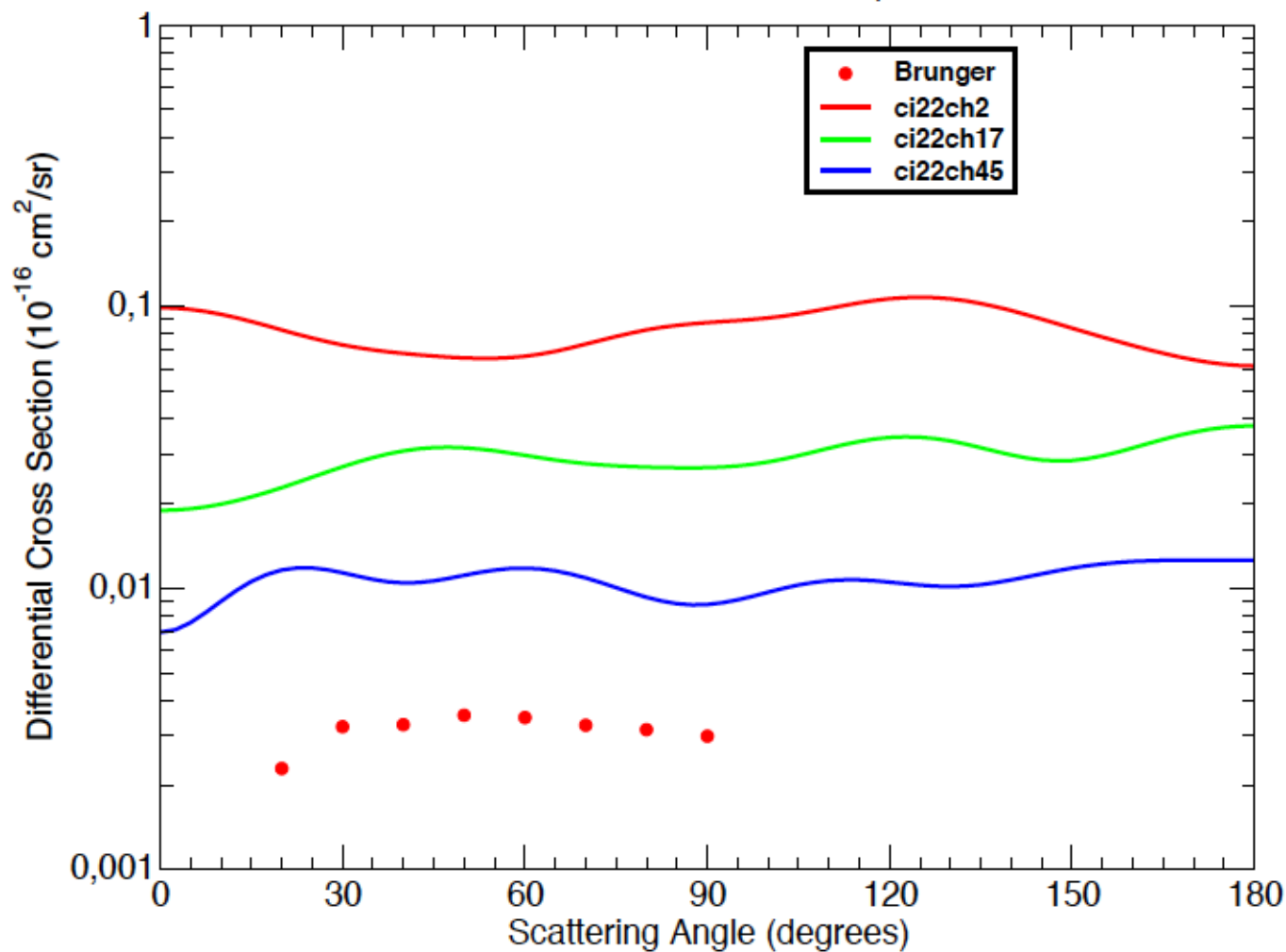




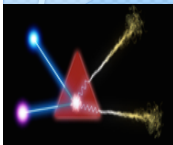
Multichannel effects on the Electronic excitation of \tilde{a}^3B_{1u} state of C_2H_4 by electron impact



30.0 eV : inelastic 1st Triplet

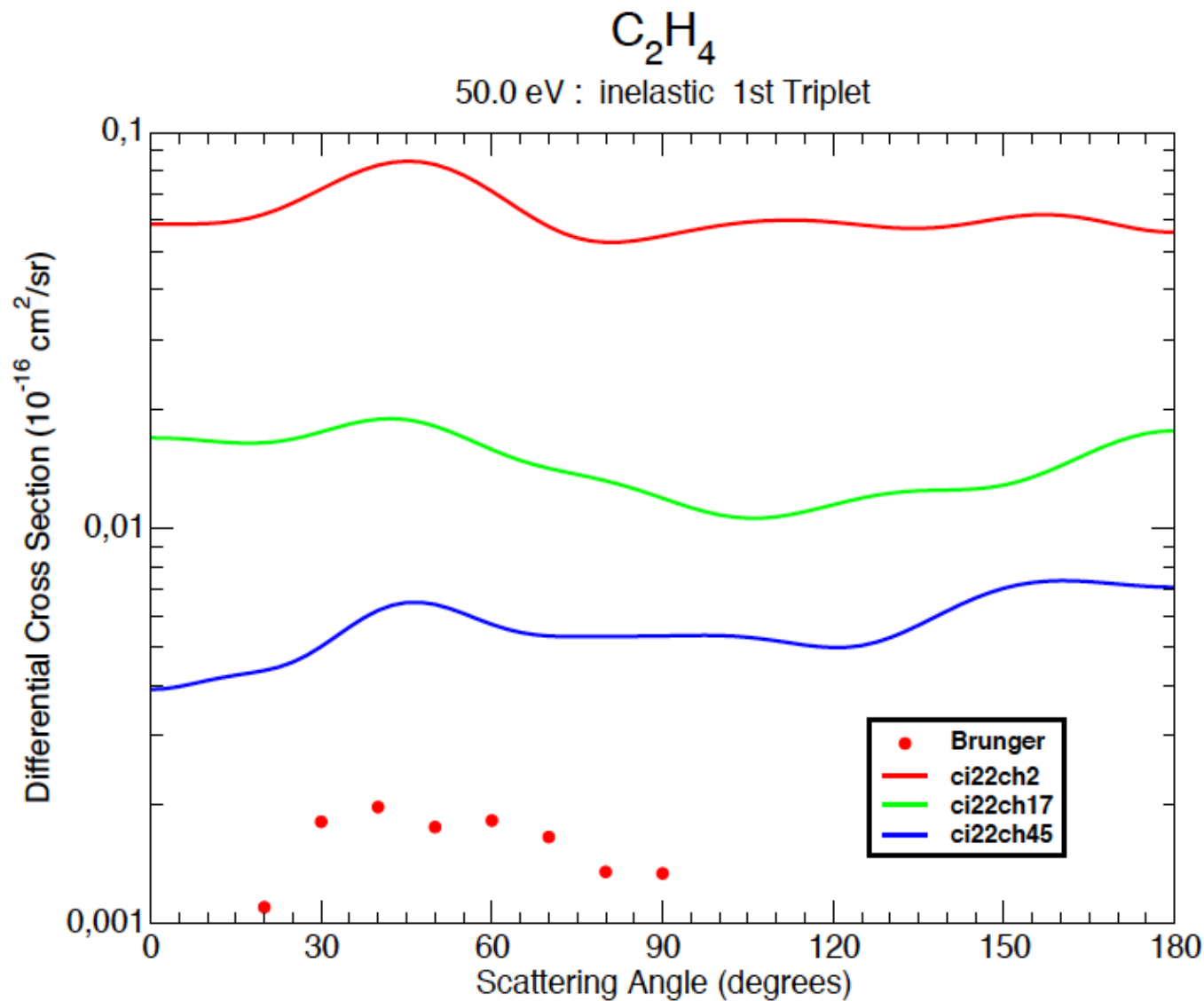


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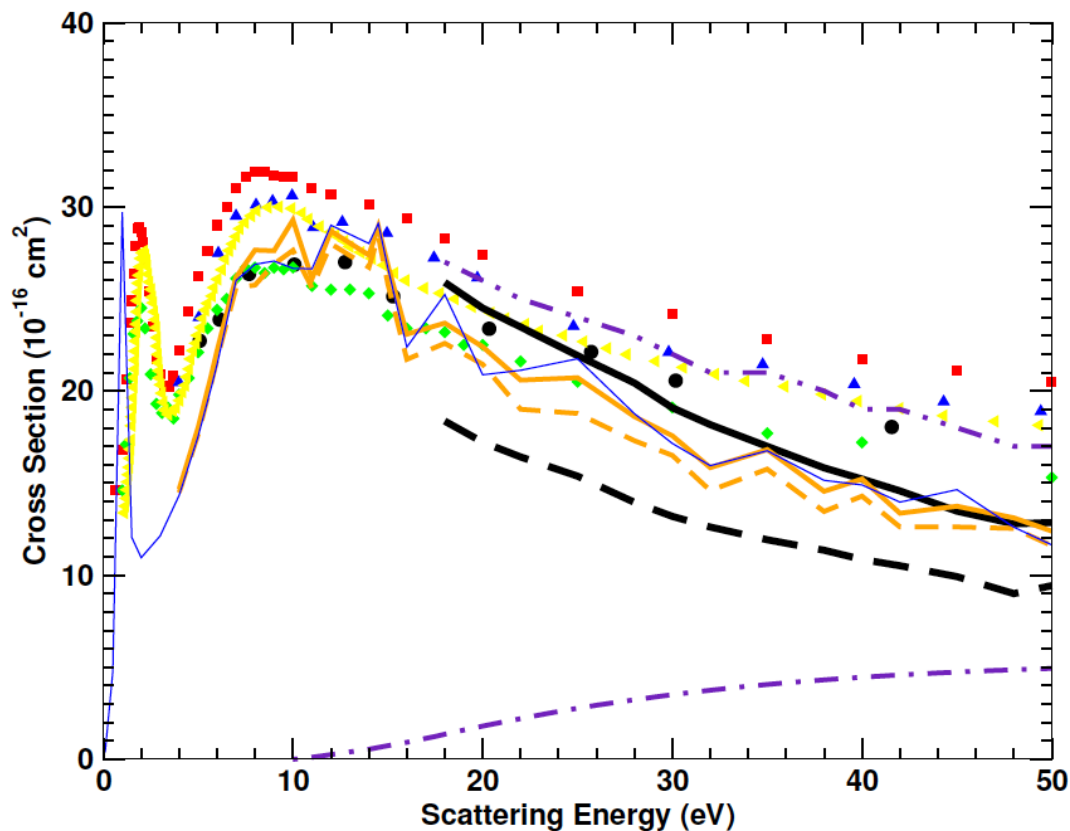


FIG. 4. (Color online) Total and integral cross sections for the energy region 0-50 eV. Thin (blue) line: 1-channel coupling plus polarization (optical theorem gives the same as the elastic transition); dashed (orange) line: elastic transitions for 2-channel coupling plus polarization; full (orange) line: total cross section for 2-channel coupling plus polarization; thick dashed (black) line: elastic transitions for 45-channel coupling plus polarization; thick full (black) line: total cross section for 45-channel coupling plus polarization; dot-dashed (indico) line: ionization cross sections using BEB approximation; dot-dot-dashed (indico) line: total cross sections 45-channel coupling plus polarization plus ionization using BEB approximation; full (black) circles, full (red) squares, full (green) diamonds, full (blue) triangles up and full (yellow) triangles down: experiments of Refs. [1–5], respectively.



Electronic excitation of Phenol by electron impact:

D. B. Jones, G. B. da Silva, R. F. C. Neves, H.V. Duque, L. Chiari, E. M. de Oliveira, **M. C.A. Lopes**, R. F. da Costa, M.T. do N.Varella, M. H. F. Bettega, M.A. P. Lima, and **M. J. Brunger**, *J. Chem. Phys.* **141**, 074314 (2014)

074314-4 Jones *et al.*

J. Chem. Phys. 141, 074314 (2014)

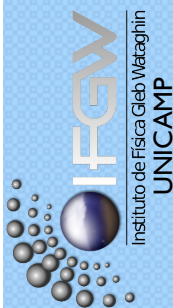
TABLE I. Experimental and calculated excitation energies, assignments, dominant configurations, and optical oscillator strengths (f_0).

Band	Expt. energy (eV)	State	TD-DFT			MOB-SCI		Full-SCI	
			Energy (eV)	Dominant excitation(s)	f_0	Energy (eV)	f_0	Energy (eV)	f_0
I	3.4–4.3	$^3A'$	3.71	$3a'' \rightarrow 5a''$; $4a'' \rightarrow 6a''$	0	3.57	0	3.29	0
		$^3A'$	4.10	$4a'' \rightarrow 5a''$	0	4.73	0	4.49	0
II	4.3–5.4	$^3A'$	4.53	$3a'' \rightarrow 5a''$; $4a'' \rightarrow 6a''$	0	4.90	0	4.78	0
		$^1A'$	4.99	$3a'' \rightarrow 6a''$; $4a'' \rightarrow 5a''$	0.0312	6.09	0.0248	5.82	0.0381
		$^3A''$	5.06	$4a'' \rightarrow 22a'$	0	6.16	0	5.94	0
		$^1A''$	5.13	$4a'' \rightarrow 22a'$	0.0001	6.21	0.0001	6.06	0.0001
		$^3A'$	5.30	$3a'' \rightarrow 6a''$	0	6.03	0	5.73	0
III	5.4–6.3	$^3A''$	5.53	$4a'' \rightarrow 23a'$	0	6.78	0	6.53	0
		$^1A''$	5.57	$4a'' \rightarrow 23a'$	0.0034	6.86	0.0274	6.68	0.0177
		$^1A'$	5.76	$3a'' \rightarrow 5a''$; $4a'' \rightarrow 6a''$	0.0328	6.80	0.0031	6.12	0.0025
		$^3A''$	5.90	$3a'' \rightarrow 22a'$; $4a'' \rightarrow 24a'$	0	6.92	0	6.73	0
		$^1A''$	5.92	$4a'' \rightarrow 24a'$	0				
		$^3A''$	5.95	$3a'' \rightarrow 22a''$; $4a'' \rightarrow 24a'$	0				
		$^1A''$	5.98	$3a'' \rightarrow 22a'$	0.0021	6.99	0	6.86	0.0020
		$^3A''$	6.27	$4a'' \rightarrow 25a'$	0				
IV	6.3–7.3	$^1A''$	6.31	$4a'' \rightarrow 25a'$	0.0115				
		$^3A''$	6.32	$3a'' \rightarrow 23a'$	0				
		$^1A''$	6.35	$3a'' \rightarrow 23a'$	0.0010				
		$^3A''$	6.52	$4a'' \rightarrow 26a'$	0				
		$^1A''$	6.54	$4a'' \rightarrow 26a'$	0				
		$^3A''$	6.63	$3a'' \rightarrow 24a'$	0				
		$^1A'$	6.66	$3a'' \rightarrow 6a''$	0.3744				
		$^1A''$	6.66	$3a'' \rightarrow 24a'$	0.0202				
		$^1A'$	6.71	$3a'' \rightarrow 5a''$; $4a'' \rightarrow 6a''$	0.5827				
		$^3A''$	6.84	$4a'' \rightarrow 27a'$; $4a'' \rightarrow 28a'$	0				
		$^3A'$	6.85	$2a'' \rightarrow 5a''$	0				
		$^3A'$	6.93	$4a'' \rightarrow 7a''$	0				
		$^1A''$	6.93	$4a'' \rightarrow 27a'$; $4a'' \rightarrow 28a'$	0.0009				
		$^1A'$	7.01	$4a'' \rightarrow 7a''$	0.0148				
		$^3A''$	7.07	$3a'' \rightarrow 25a'$	0				
$^1A''$	7.08	$3a'' \rightarrow 25a'$	0						
$^3A'$	7.11	$2a'' \rightarrow 6a''$	0						
$^3A''$	7.19	$4a'' \rightarrow 27a'$; $4a'' \rightarrow 28a'$	0						
$^1A''$	7.22	$4a'' \rightarrow 27a'$; $4a'' \rightarrow 28a'$	0.0005						
$^3A''$	7.27	$21a' \rightarrow 5a''$	0						
$^3A''$	7.29	$3a'' \rightarrow 26a'$	0						
V	7.3–8.6	$^1A''$	7.32	$3a'' \rightarrow 26a'$	0.0002				
		$^1A''$	7.57	$21a' \rightarrow 5a''$	0.0043				
		$^3A''$	7.58	$4a'' \rightarrow 29a'$	0				
		$^3A''$	7.59	$3a'' \rightarrow 27a'$; $3a'' \rightarrow 28a'$	0				

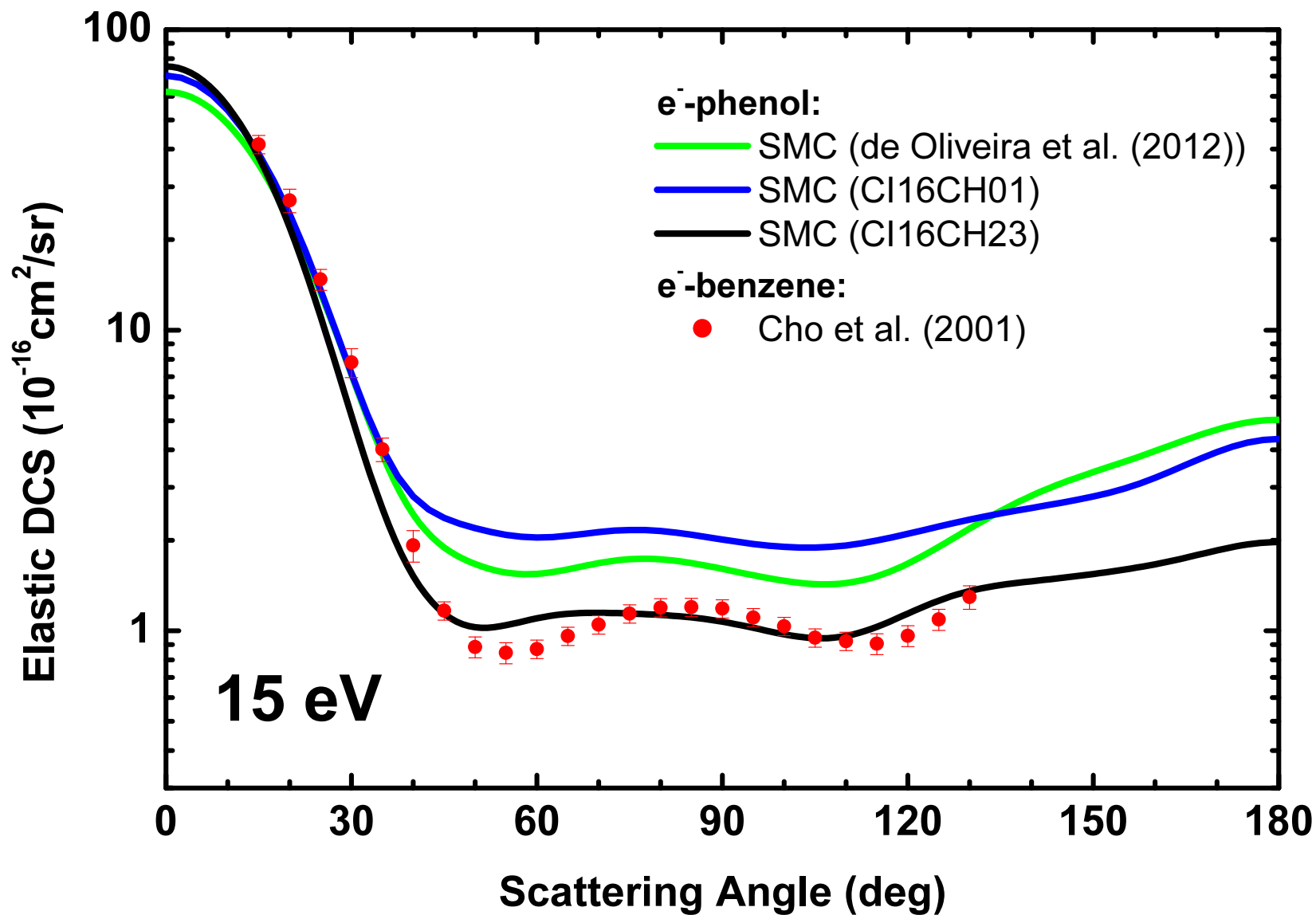
Our scattering calculations have 5 singlets and 7 triplets below 7eV in good agreement with the full single configuration interaction. We also included 20 additional pseudo states as possible open channels (a total of 33 channels)



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Electronic excitation of Phenol by electron impact: Effects on the Elastic channel

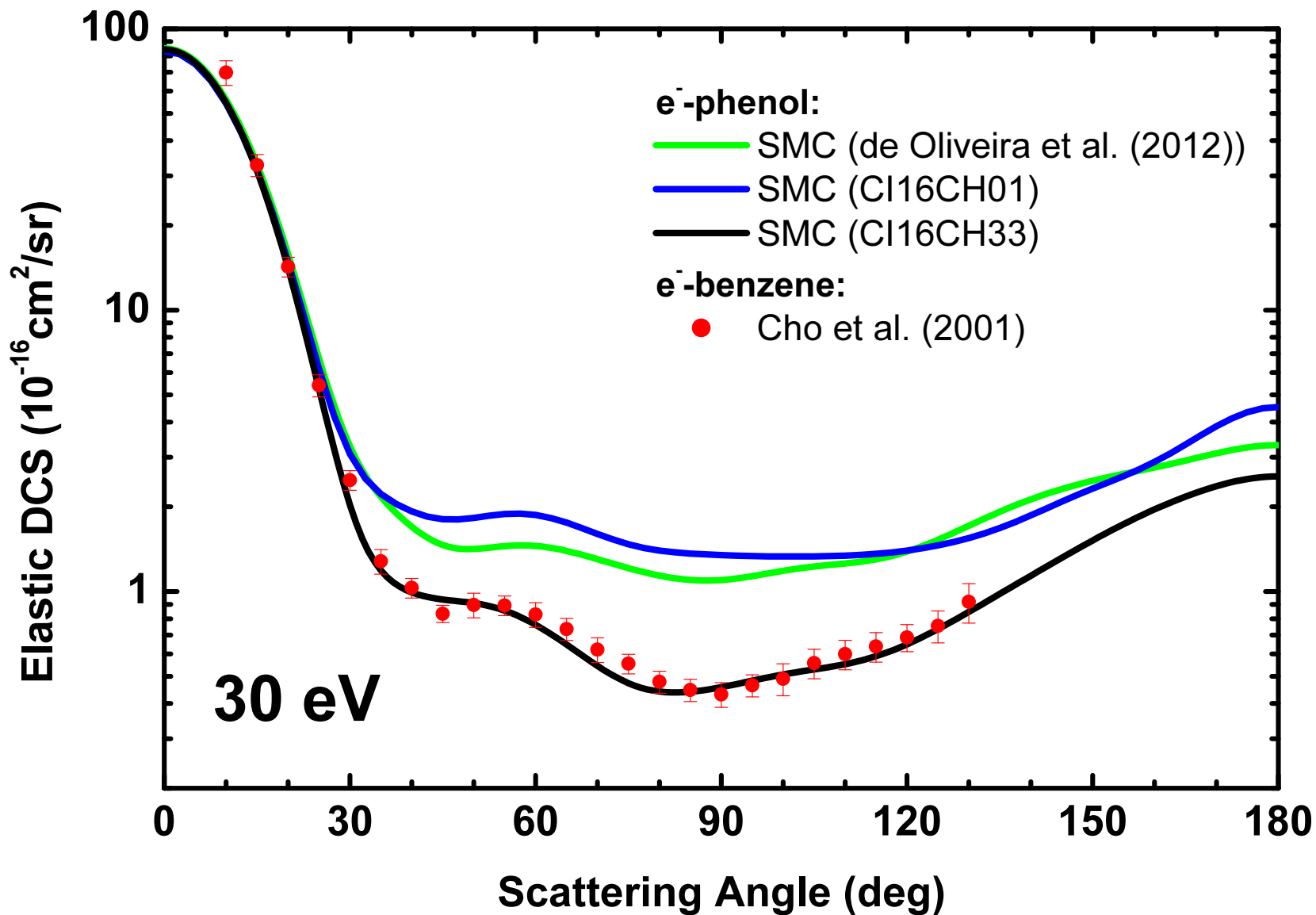




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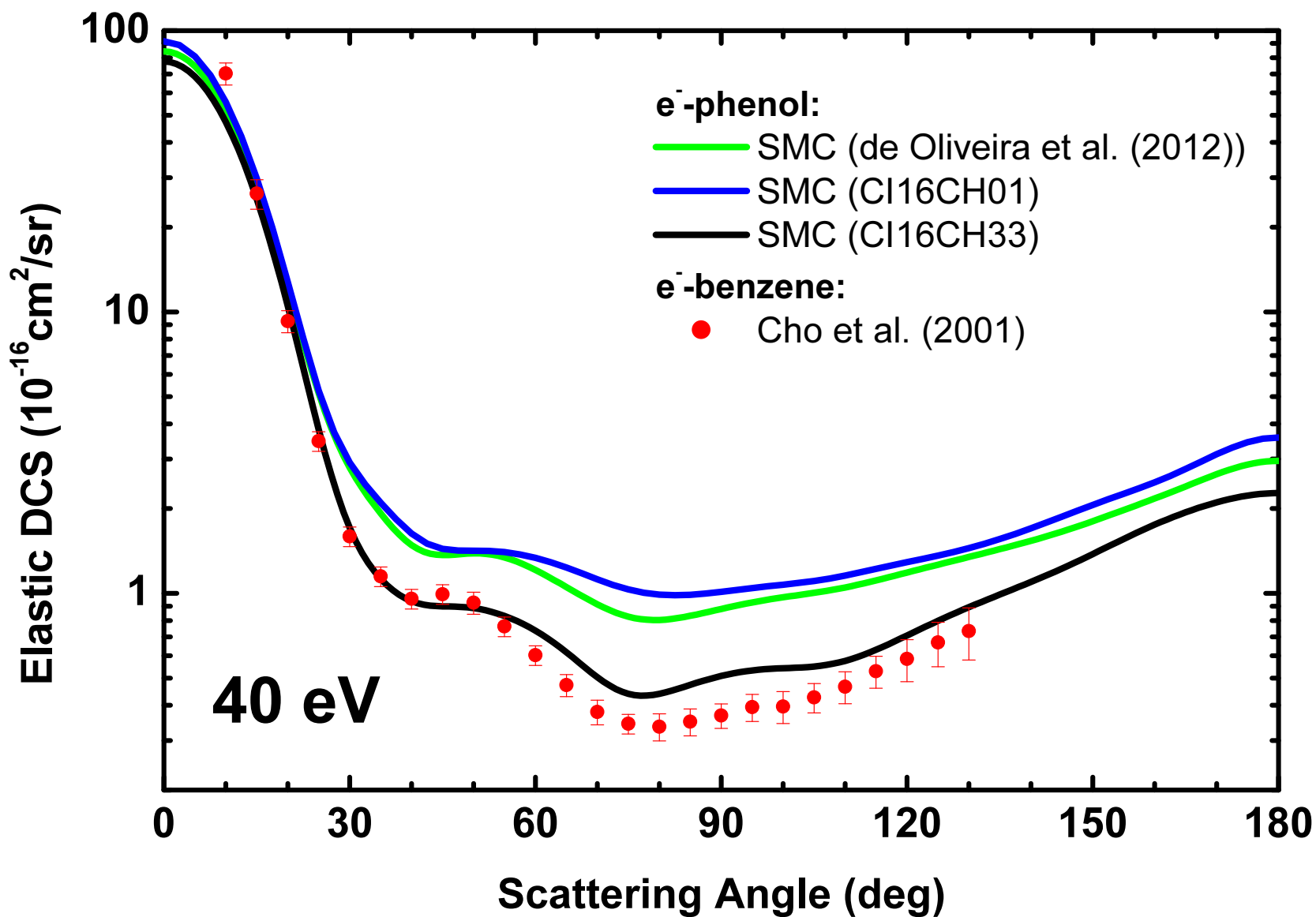
Electronic excitation of Phenol by electron impact: Effects on the Elastic channel





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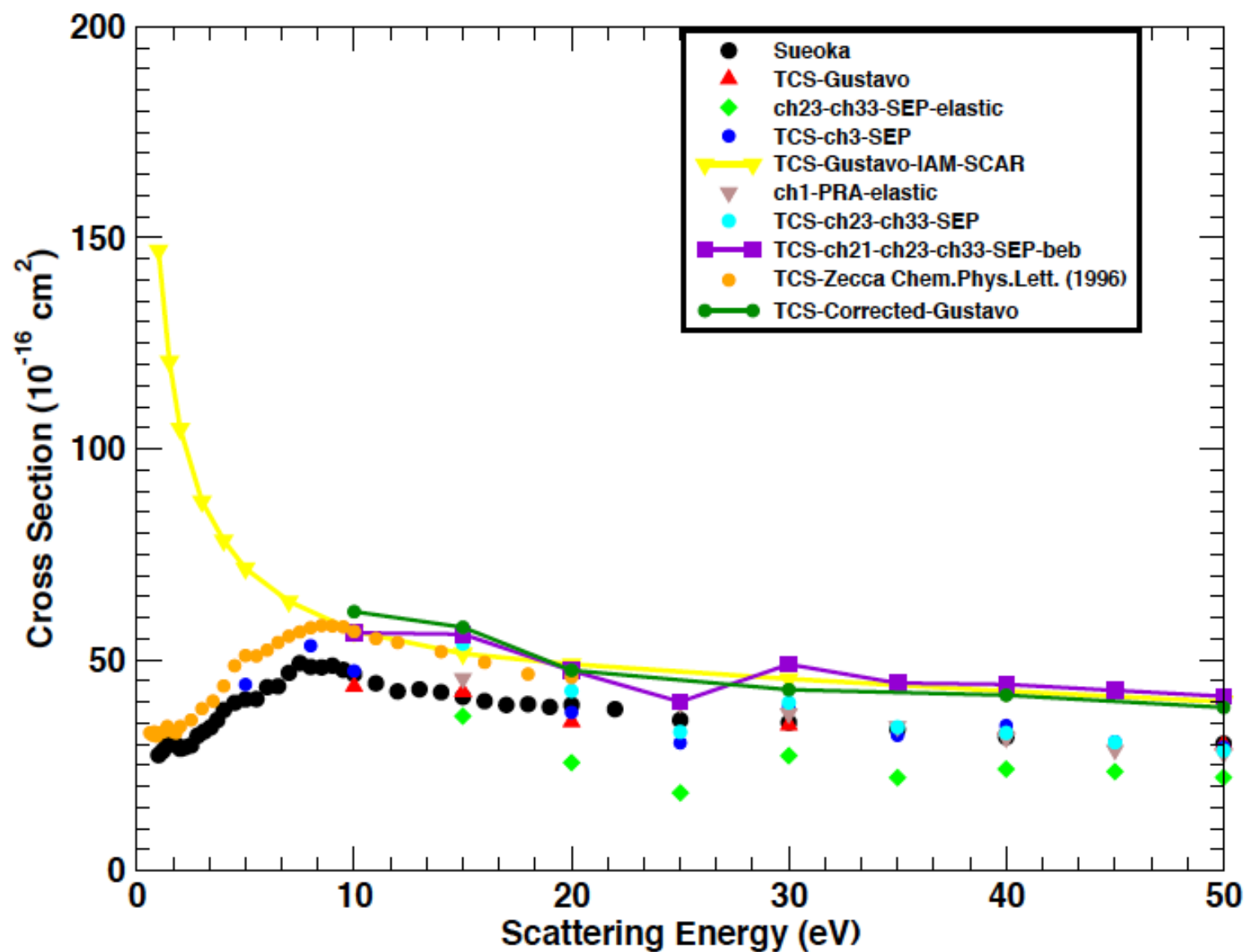
Electronic excitation of Phenol by electron impact: Effects on the Elastic channel





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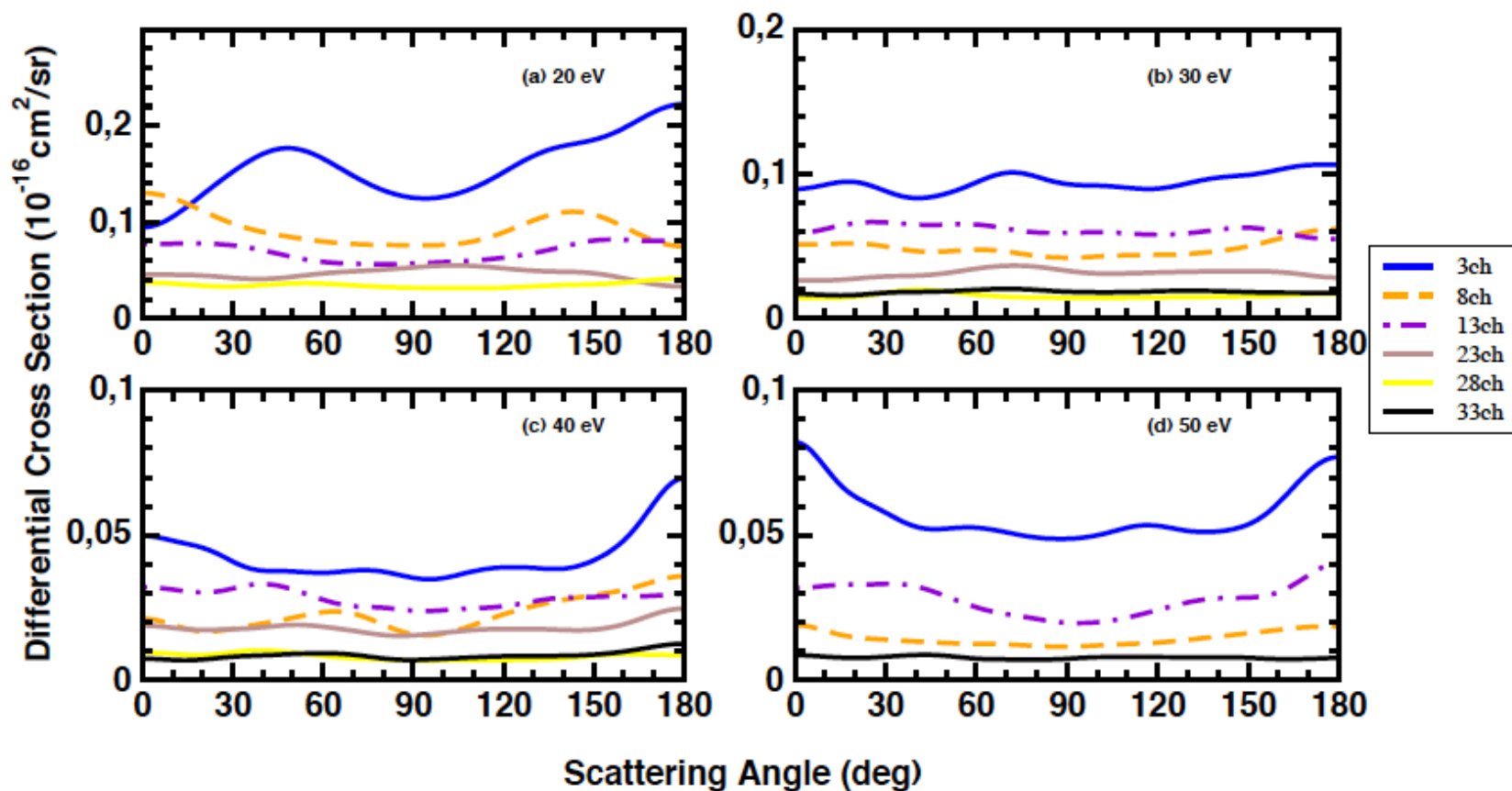
Phenol: Integral and Total Cross Sections





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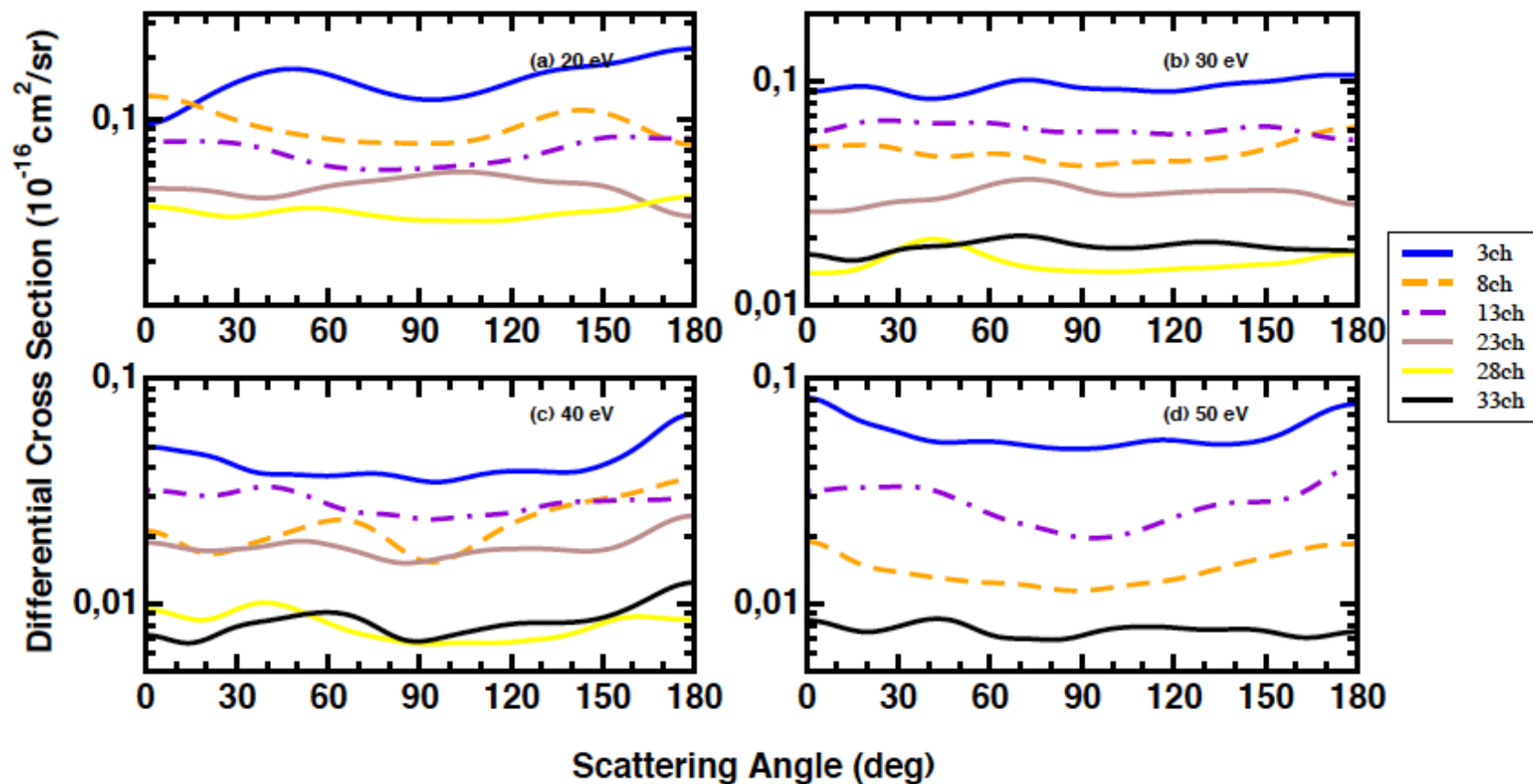
Phenol: electronic excitation of the 1st triplet





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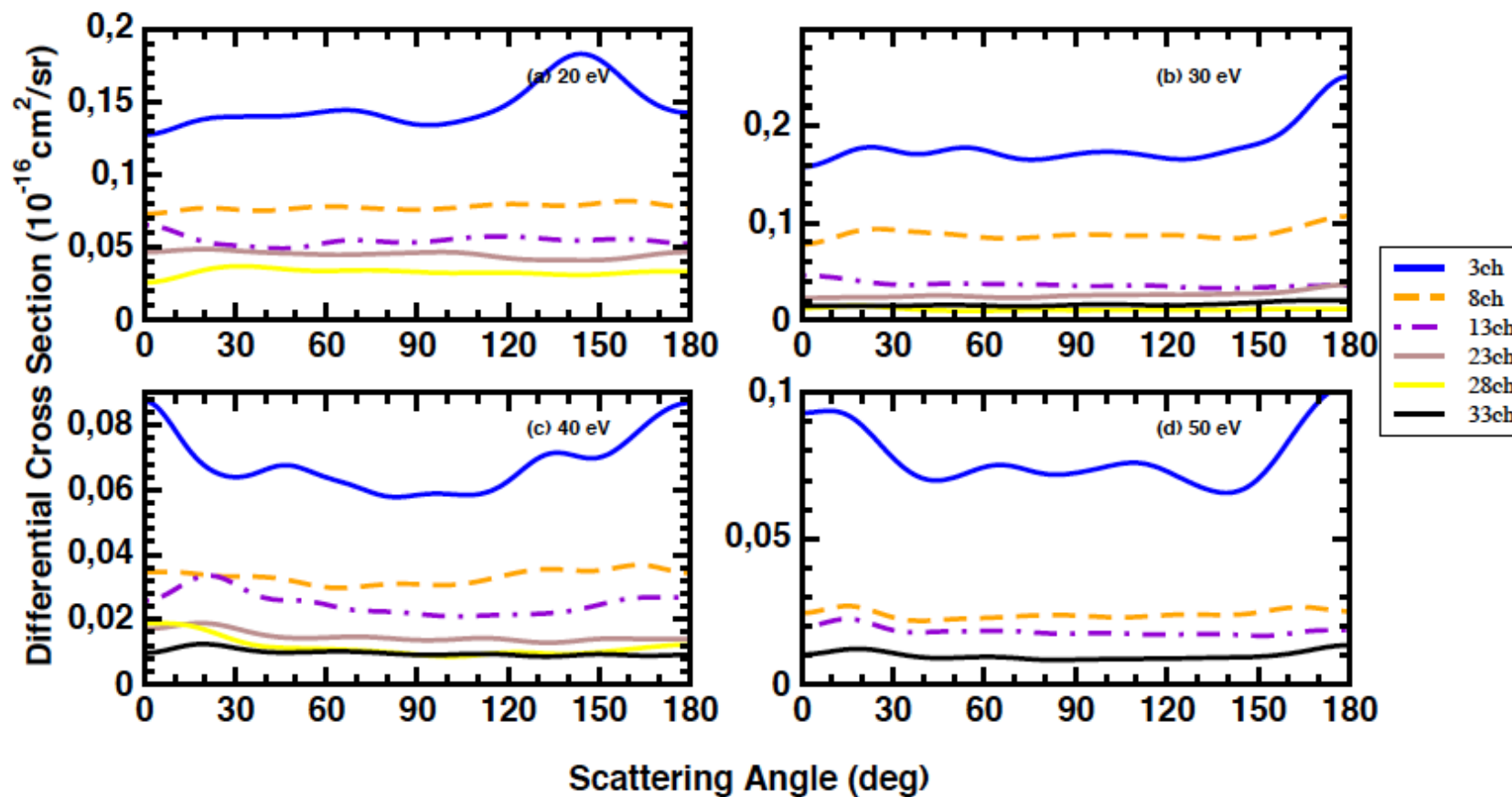
Phenol: electronic excitation of the 1st triplet





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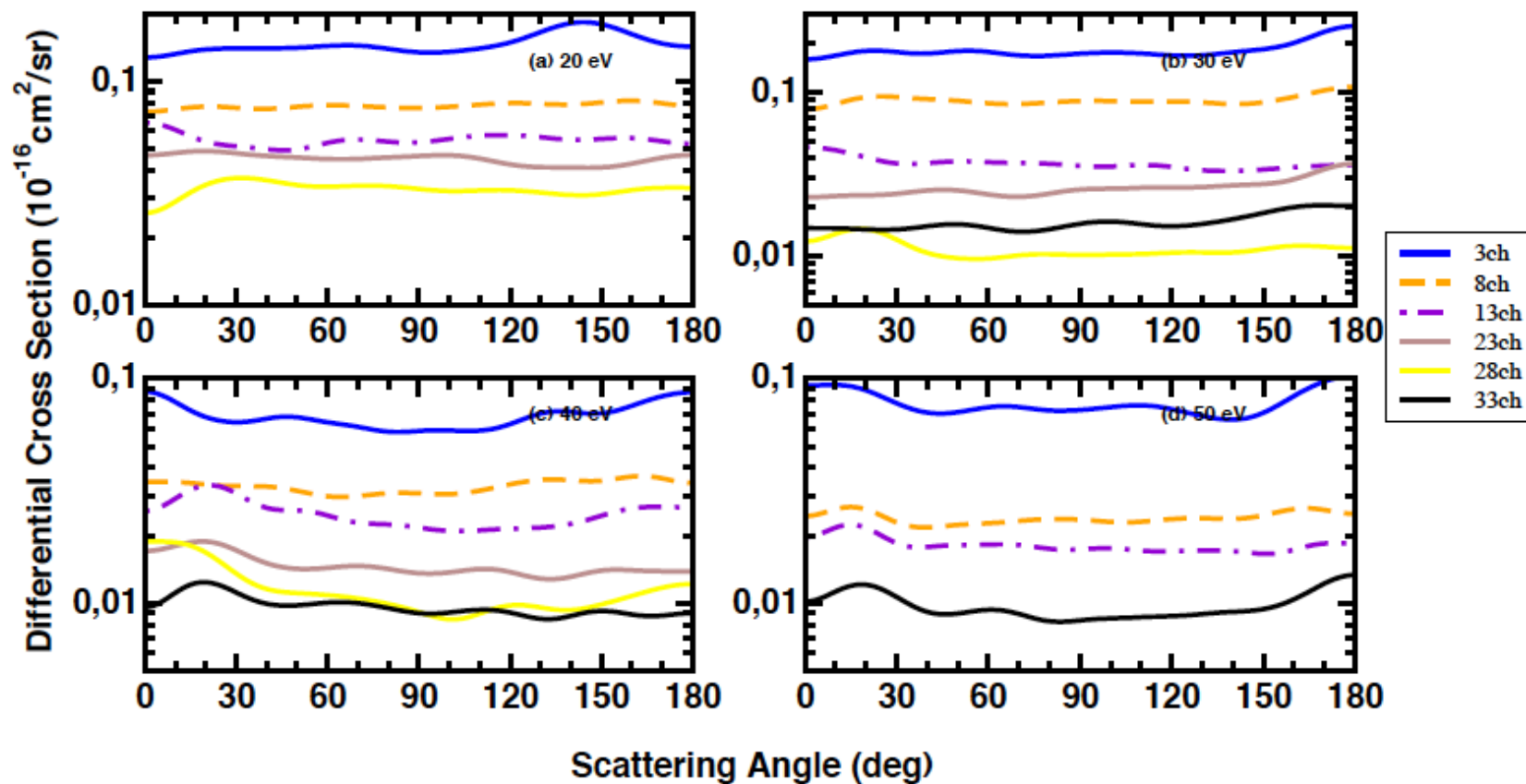
Phenol: electronic excitation of the 2nd triplet

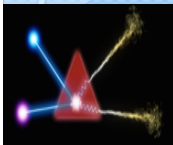




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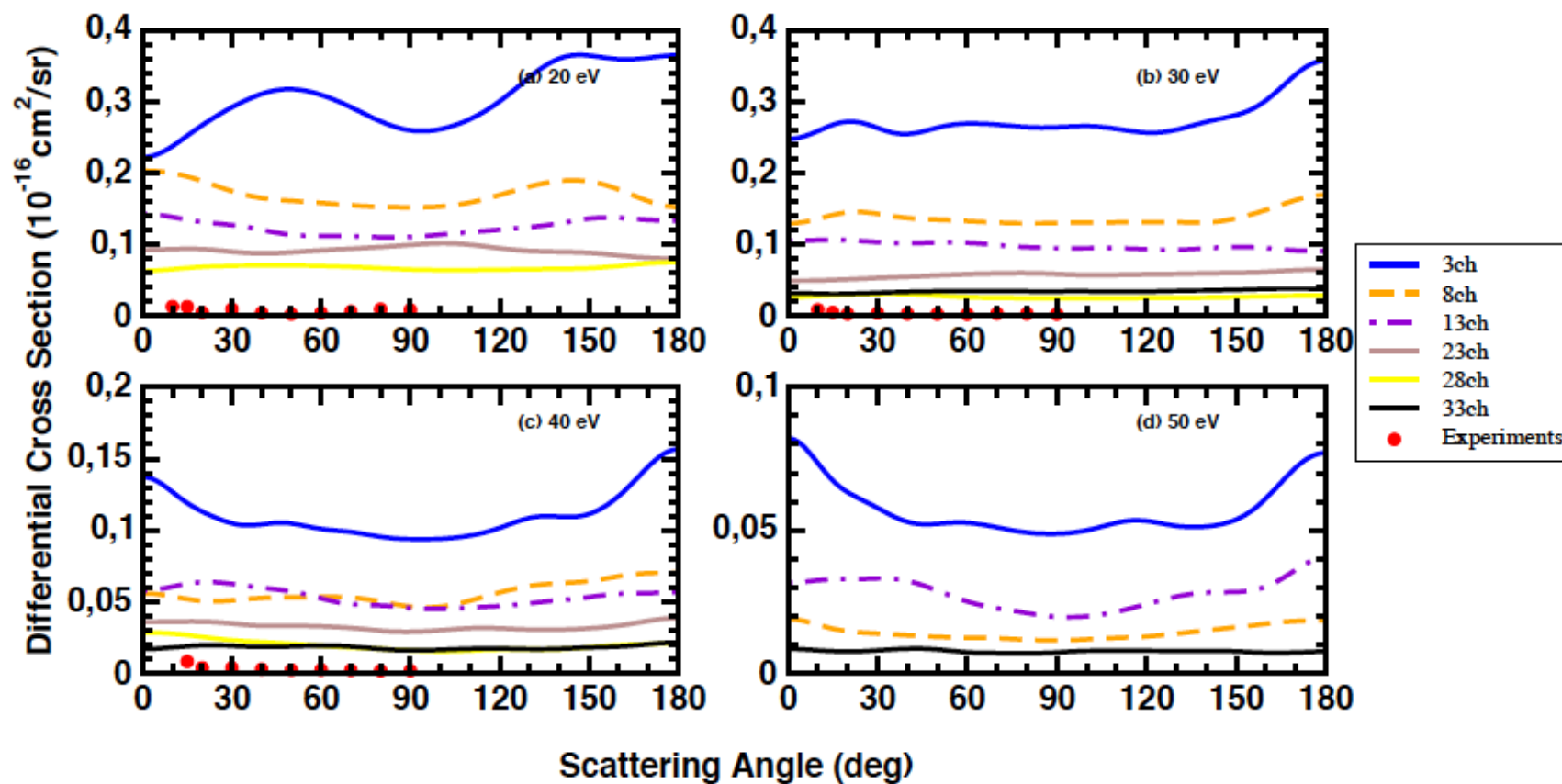
Phenol: electronic excitation of the 2nd triplet

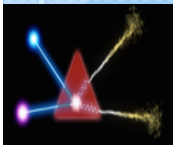




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J. Chem. Phys. **142**, 104305 (2015).

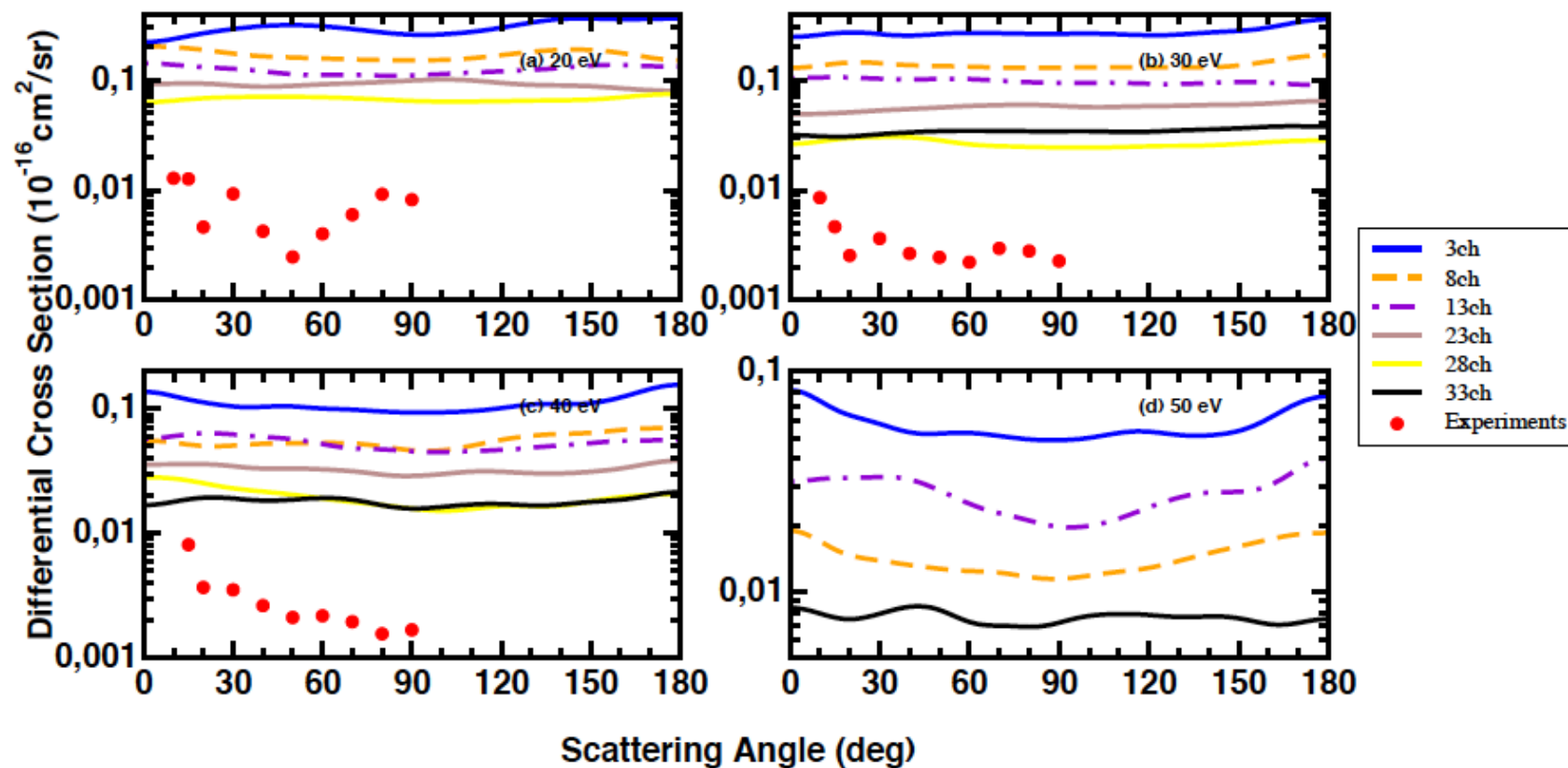
Phenol: electronic excitation of Band I

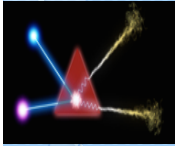




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Phenol: electronic excitation of Band I





Thank you very much for your attention

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