Exercise Class 1

Consider a particle of mass m subject to the potential

$$V(x) = \begin{cases} \infty & x \le 0, \\ -V_0 \delta(x-a) & x > 0, \end{cases}$$
(1)

where $V_0 > 0$.

a) Show that the derivative of the eigenfunction $\psi(x)$ presents a discontinuity at x = a and determine it in terms of V_0 , m and $\psi(a)$.

b) Discuss the existence of bound states in terms of the size of a.

Solution

a) Let's integrate the Schrödinger equation for x > 0 from $a - \varepsilon$ to $a + \varepsilon$:

$$-\frac{\hbar^2}{2m}\int_{a-\varepsilon}^{\varepsilon+a} dx \left[\frac{d^2\psi(x)}{dx^2}\right] - V_0 \int_{a-\varepsilon}^{a+\varepsilon} dx \delta(x-a)\psi(x) = E \int_{a-\varepsilon}^{a+\varepsilon} dx\psi(x),$$
$$-\frac{\hbar^2}{2m} \left[\frac{d\psi(a+\varepsilon)}{dx} - \frac{d\psi(a-\varepsilon)}{dx}\right] - V_0\psi(a) = E \int_{a-\varepsilon}^{a+\varepsilon} dx\psi(x),$$

taking $\varepsilon \to 0$ we obtain

$$\frac{d\psi(a^{+})}{dx} - \frac{d\psi(a^{-})}{dx} = -\frac{2mV_0}{\hbar^2}\psi(a).$$
 (2)

b) For x > 0 the Schrödinger equation is written as

$$-\frac{\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} - V_0\delta(x-a)\psi(x) = E\psi(x).$$
 (3)

For x > a we have

$$-\frac{\hbar^2}{2m}\frac{d^2\psi_R(x)}{dx^2} = E\psi_R(x),\tag{4}$$

$$\frac{d^2\psi_R(x)}{dx^2} = k\psi_R(x),\tag{5}$$

and for x < a:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi_L(x)}{dx^2} = E\psi_L(x),\tag{6}$$

$$\frac{d^2\psi_L(x)}{dx^2} = k^2\psi_L(x).$$
(7)

Note that in both equations, (5) and (7),

$$k = \frac{\sqrt{-2mE}}{\hbar},\tag{8}$$

and as we are interested in the bound states of the system, we can write

$$E = -|E|,\tag{9}$$

and therefore,

$$k = \frac{\sqrt{2m|E|}}{\hbar}.$$
(10)

Based on (5) and (7) the solutions of the equations are of the form:

$$\psi(x) = \begin{cases} \psi_R(x) = A_R e^{kx} + B_R e^{-kx}, & x > a\\ \psi_L(x) = A_L e^{kx} + B_L e^{-kx}, & x < a \end{cases}.$$
 (11)

Let's apply the boundary conditions to obtain the eigenvalues of the problem. (i) $x \to \infty \implies \psi(x) \to 0 \implies \psi_R(x) \to 0$:

$$\Rightarrow A_R = 0 \tag{12}$$

(ii) $\psi(0) = 0 \implies \psi_L(0) = 0$:

$$\Rightarrow B_L = -A_L \tag{13}$$

Using (12) and (13) in (11), we can write

$$\begin{cases} \psi_R(x) = B_R e^{-kx}, & x > a\\ \psi_L(x) = A_L \left(e^{kx} - e^{-kx} \right), & x < a \end{cases}.$$
(14)

(iii) $\psi_L(a) = \psi_R(a)$:

$$\Rightarrow B_R e^{-ka} = A_L \left(e^{ka} - e^{-ka} \right) \tag{15}$$

(iv) Descontinuity of the derivative calculated in (2):

$$\frac{d\psi_R(x)}{dx} = -kB_R e^{-kx},\tag{16}$$

$$\frac{d\psi_L(x)}{dx} = kA_L \left(e^{kx} + e^{-kx}\right). \tag{17}$$

Substituting (16) and (17) in (2) we obtain

$$B_R k e^{-ka} + k A_L \left(e^{ka} + e^{-ka} \right) = \frac{2mV_0}{\hbar^2} B_R e^{-ka}.$$
 (18)

Substituting (15) in (18) we find that

$$kA_L\left(e^{ka} - e^{-ka}\right) + kA_L\left(e^{ka} + e^{-ka}\right) = \frac{2mV_0}{\hbar^2}A_L\left(e^{ka} - e^{-ka}\right)$$
$$kA_Le^{ka} = \frac{mV_0}{\hbar^2}A_L\left(e^{ka} - e^{-ka}\right)$$

$$\Rightarrow A_L \left[k - \frac{mV_0}{\hbar^2} \left(1 - e^{-2ka} \right) \right] = 0.$$
(19)

We are looking for nontrivial solutions, i.e., $A_L \neq 0$:

$$\Rightarrow k = \frac{mV_0}{\hbar^2} \left(1 - e^{-2ka} \right) \tag{20}$$

The equation (20) is a transcedental equation (very common in the solution of physical problems!!). The solutions of this equation is given by the intersection points between the function f(k) = k and the function $g(k) = \frac{mV_0}{\hbar^2} \left(1 - e^{-2ka}\right)$. One way of determining the solutions (in most problems we can just solve this kind of equation numerically) is note that this equation will just have a solution if at the point x = 0:

$$\left[\frac{dg(k)}{dk} \right]_{k=0} > \left[\frac{df(k)}{dk} \right]_{k=0}$$

$$\Rightarrow a > \frac{\hbar^2}{2mV_0}$$

$$(21)$$

Thus, using (20) and the condition (21), we obtain that the energy of the bound state is given by

$$\Rightarrow E = -\frac{mV_0^2}{2\hbar^2} \left(1 - e^{-2ka}\right)^2, \quad a > \frac{\hbar^2}{2mV_0} \tag{22}$$