

F-315 (Mecânica Geral I)

Aula 22

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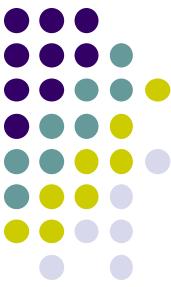
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Slides do prof. Antonio Vidiella Barranco:

<http://www.ifi.unicamp.br/~vidiella/aulas.html>



Cálculo Variacional

TM seção 6.5

Funções com várias variáveis

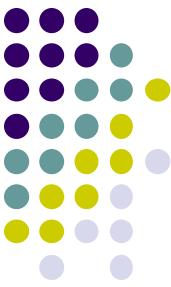
$$f = f\{y_i(x), y'_i(x); x\} \quad i = 1, 2, \dots, n$$

Analogamente,

$$\frac{\partial J}{\partial \alpha} = \int_{x_1}^{x_2} \sum_i \left[\frac{\partial f}{\partial y_i} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_i} \right) \right] \eta_i(x) dx$$

Resulta no conjunto de equações de Euler

$$\frac{\partial f}{\partial y_i} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_i} \right) = 0 \quad \text{TM eq. (6.57)}$$



Cálculo Variacional

TM seção 6.6

Equações de Euler com vínculos

TM eq. (6.69)

$$\frac{\partial f}{\partial y_i} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_i} \right) + \sum_j \lambda_j(x) \frac{\partial g_j}{\partial y_i} = 0$$

Onde os vínculos $g_j\{y_i; x\} = 0$ e $i = 1, 2, \dots, n$
 $j = 1, 2, \dots, m$

As funções λ_j são os multiplicadores de Lagrange

TM seção 6.6, pgs. 219–222; G seção 2.4, pgs. 45–49

$f = f\{y_i, y'_i; x\}$: múltiplas variáveis dependentes $\{y_i(x)\}$
 $i = 1, \dots, n$, com vínculos $g_k = g_k\{y_i; x\} = 0$, $k = 1, \dots, m$

Exemplo: $g_1 = \sum_{i=1}^3 x_i^2 - r^2 = 0$ (\mathbf{r} sobre a esfera de raio r)

$$y_i(\alpha, x) = y_i(0, x) + \alpha \eta_i(x), \quad \frac{\partial y_i}{\partial \alpha} = \eta_i(x), \quad \frac{\partial y'_i}{\partial \alpha} = \frac{d\eta_i(x)}{dx} = \eta'_i(x)$$

$$\frac{\partial J}{\partial \alpha} = \int_{x_1}^{x_2} \sum_{i=1}^n \left(\frac{\partial f}{\partial y_i} \frac{\partial y_i}{\partial \alpha} + \frac{\partial f}{\partial y'_i} \frac{\partial y'_i}{\partial \alpha} \right) dx = \int_{x_1}^{x_2} \sum_{i=1}^n \left[\frac{\partial f}{\partial y_i} \eta_i(x) + \frac{\partial f}{\partial y'_i} \eta'_i(x) \right] dx$$

$$= \int_{x_1}^{x_2} \sum_{i=1}^n \left[\frac{\partial f}{\partial y_i} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_i} \right) \right] \eta_i(x) dx = 0.$$

Agora os $\{\eta_i(x)\}$ não são mais independentes, mas estão vinculados através da

$$\text{condição sobre } dg_k = \sum_{i=1}^n \frac{\partial g_k}{\partial y_i} \frac{\partial y_i}{\partial \alpha} d\alpha = \sum_{i=1}^n \frac{\partial g_k}{\partial y_i} \eta_i(x) d\alpha = 0.$$

Exemplo: um vínculo ($m = 1$) com $n = 2$, $g(y_1, y_2; x)$

$$\frac{dg}{da} = \frac{\partial g}{\partial y_1} \eta_1(x) + \frac{\partial g}{\partial y_2} \eta_2(x) = 0 \rightarrow \frac{\eta_2(x)}{\eta_1(x)} = -\frac{\partial g/\partial y_1}{\partial g/\partial y_2}$$

$$\begin{aligned}\frac{\partial J}{\partial a} &= \int_{x_1}^{x_2} \left\{ \left[\frac{\partial f}{\partial y_1} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_1} \right) \right] \eta_1(x) + \left[\frac{\partial f}{\partial y_2} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_2} \right) \right] \eta_2(x) \right\} dx \\ &= \int_{x_1}^{x_2} \left\{ \left[\frac{\partial f}{\partial y_1} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_1} \right) \right] - \left[\frac{\partial f}{\partial y_2} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_2} \right) \right] \frac{\partial g/\partial y_1}{\partial g/\partial y_2} \right\} \eta_1(x) dx = 0.\end{aligned}$$

Agora o integrando entre colchetes se anula $\forall \eta_1(x) \rightarrow$

$$\lambda(x) \equiv - \left[\frac{\partial f}{\partial y_1} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_1} \right) \right] \left(\frac{\partial g}{\partial y_1} \right)^{-1} = - \left[\frac{\partial f}{\partial y_2} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_2} \right) \right] \left(\frac{\partial g}{\partial y_2} \right)^{-1}$$

é o chamado ***multiplicador de Lagrange.***

$$\frac{\partial f}{\partial y_1} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_1} \right) + \lambda(x) \frac{\partial g}{\partial y_1} = \frac{\partial f}{\partial y_2} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_2} \right) + \lambda(x) \frac{\partial g}{\partial y_2} = 0$$

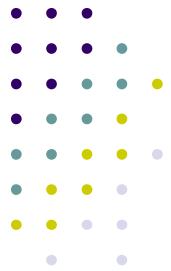
Para m vínculos **holonômicos** $g_k = g_k\{y_i; x\} = 0$, isto é equivalente a extremizar $\tilde{J} \equiv \int_{x_1}^{x_2} [f + \sum_{k=1}^m \lambda_k(x) g_k] dx$ como se $\{y_i(x), \lambda_k(x)\}$ fossem variáveis independentes entre si.

Equações de Euler com vínculos:

TM eq. (6.69)

$$\frac{\partial f}{\partial y_i} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'_i} \right) + \sum_{k=1}^m \lambda_k(x) \frac{\partial g_k}{\partial y_i} = 0$$

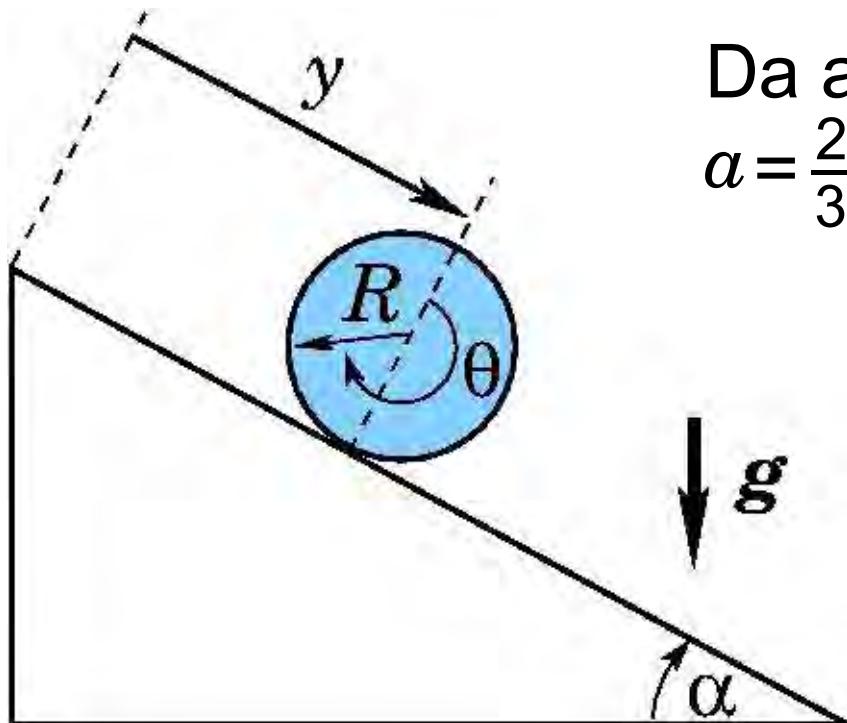
O caso geral **não-holonômico** $g_k = g_k\{y_i, y'_i; x\} = 0$ não é trivial: ver FLANNERY, *Am. J. Phys.* **73**, 265–272 (2005). No entanto, se o vínculo está na forma integral $K\{y_i\} = \int_{x_1}^{x_2} g\{y_i, y'_i; x\} dx$, então extremiza-se $\int_{x_1}^{x_2} (f + \lambda g) dx$.



Problema

TM exemplo 6.5, pgs. 221-222

Disco de raio R rolando sem escorregar em um plano inclinado. Obter a equação de vínculo $g(y, \theta)$.

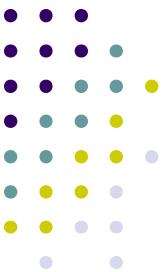


Da aula 16, aceleração:
 $a = \frac{2}{3}g \sin \alpha$

$$y = R\theta$$

$$g(y, \theta) = y - R\theta = 0$$

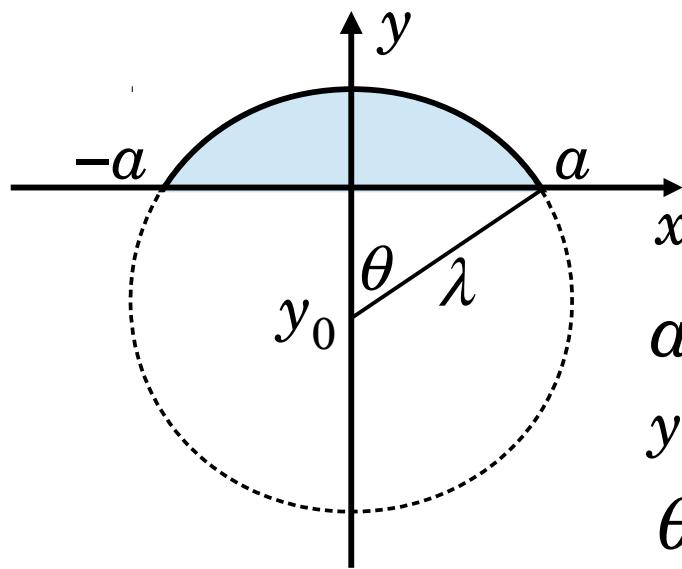
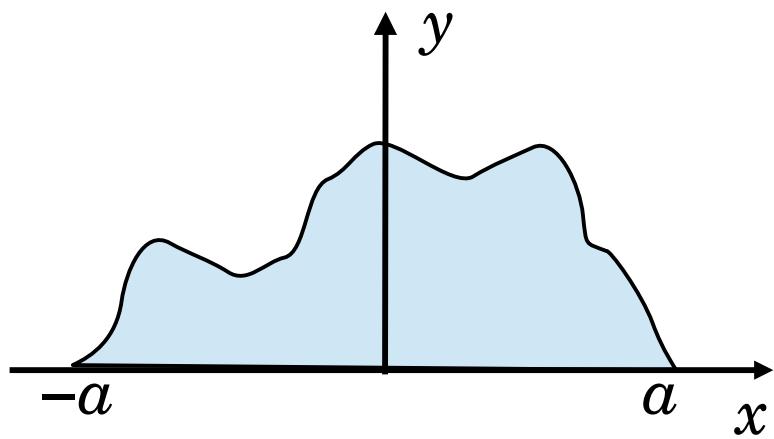
$$\frac{\partial g}{\partial y} = 1, \quad \frac{\partial g}{\partial \theta} = -R$$



Problema

TM exemplo 6.6, pgs. 222-224

Problema isoperimétrico: Encontrar a curva $y(x)$, de comprimento ℓ , limitada inferiormente pelo eixo x , que passa pelos pontos $(-a,0)$ e $(a,0)$ de modo que a área circunscrita seja máxima.



$$a = \lambda \sin \theta$$

$$y_0 = \lambda \cos \theta$$

$$\theta = \ell / 2\lambda$$

TM exemplo 6.6 (problema de Dido), pgs. 222–224

Problema de maximizar a área $A = \int_{-a}^a y \, dx$ com comprimento fixo

$\ell = \int_{-a}^a ds = \int_{-a}^a \sqrt{1+y'^2} \, dx$. Condições de contorno: $y(x = \pm a) = 0$.

Maximizar $\tilde{J} = \int_{-a}^a \left(y + \lambda \sqrt{1+y'^2} \right) dx \equiv \int_{-a}^a \tilde{f} \, dx$. A função $\tilde{f}(y, y'; x)$

satisfaz a equação de Euler: $\frac{\partial \tilde{f}}{\partial y} - \frac{d}{dx} \left(\frac{\partial \tilde{f}}{\partial y'} \right) = 0$

$$\frac{\partial \tilde{f}}{\partial y} = 1 = \frac{d}{dx} \left(\frac{\partial \tilde{f}}{\partial y'} \right) = \frac{d}{dx} \left[\frac{\lambda y'}{(1+y'^2)^{1/2}} \right] \rightarrow \frac{\lambda y'}{(1+y'^2)^{1/2}} = x + \text{constante} = x - x_0 \rightarrow$$

$$\lambda^2 y'^2 = (x - x_0)^2 (1 + y'^2) \rightarrow y'(x) = \frac{\pm(x-x_0)}{[\lambda^2 - (x-x_0)^2]^{1/2}} \rightarrow y(x) = y_0 \mp [\lambda^2 - (x-x_0)^2]^{1/2}$$

Caso geral, circunferência de raio λ : TM eq. (6.85)

$$(x - x_0)^2 + (y - y_0)^2 = \lambda^2, \quad x_0 = 0, \quad y_0 = \lambda \cos \frac{\ell}{2\lambda}, \quad a = \lambda \operatorname{sen} \frac{\ell}{2\lambda}$$

$$\text{Equação para } a \text{ provém de } \ell = \int_{-a}^a ds = \lambda \int_{-a/\lambda}^{a/\lambda} \frac{du}{\sqrt{1-u^2}} = 2\lambda \operatorname{arcsen} \frac{a}{\lambda}.$$

G problema 2.9 (curva catenária), pg. 65

Problema da curva catenária: minimizar potencial gravitacional

$$V = \rho g \int_{-a}^a y \, ds = \rho g \int_{-a}^a y \sqrt{1+y'^2} \, dx \text{ com comprimento fixo } \ell = \int_{-a}^a ds \\ = \int_{-a}^a \sqrt{1+y'^2} \, dx. \text{ Condições de contorno: } b \equiv y(x=a) - y(x=-a).$$

Minimizar $\tilde{J} = \int_{-a}^a (y + \lambda) \sqrt{1+y'^2} \, dx \equiv \int_{-a}^a \tilde{f} \, dx$. A função $\tilde{f}(y, y'; x)$ satisfaaz a equaçao de Euler: $\frac{\partial \tilde{f}}{\partial y} - \frac{d}{dx} \left(\frac{\partial \tilde{f}}{\partial y'} \right) = 0$

$$\frac{\partial \tilde{f}}{\partial y} = (1+y'^2)^{1/2} = \frac{d}{dx} \left(\frac{\partial \tilde{f}}{\partial y'} \right) = \frac{d}{dx} \left[\frac{(y+\lambda)y'}{(1+y'^2)^{1/2}} \right] = \frac{y'^2 + (y+\lambda)y''}{(1+y'^2)^{1/2}} - \frac{(y+\lambda)y'^2 y''}{(1+y'^2)^{3/2}}$$

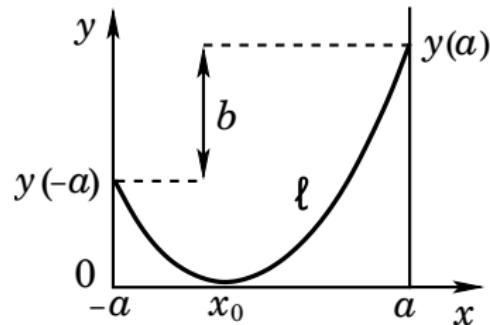
$$\rightarrow (y+\lambda)y'' = 1+y'^2 \rightarrow \frac{2y'y''}{1+y'^2} = \frac{2y'}{y+\lambda} \rightarrow \frac{d}{dx} \ln(1+y'^2) = \frac{d}{dx} \ln(y+\lambda)^2$$

$$\rightarrow \ln(1+y'^2) = C + \ln(y+\lambda)^2 \rightarrow 1+y'^2 = \left(\frac{y+\lambda}{\lambda c}\right)^2 \rightarrow y' = \frac{dy}{dx} =$$

$$= \pm \frac{1}{\lambda c} \sqrt{(y+\lambda)^2 - (\lambda c)^2} \rightarrow \Delta x \equiv x - x_0 = \pm \lambda c \int_{y_0/\lambda}^{y/\lambda} \frac{du}{\sqrt{(1+u)^2 - c^2}}$$

G problema 2.9 (curva catenária), pg. 65

Escolhendo x_0 tal que $y_0 \equiv y(x_0) = 0$ e $y'(x_0) = 0 \rightarrow c = 1$, ou seja, x_0 representa um mínimo da função $y(x)$, conforme mostra a figura ao lado:



$$\Delta x = \lambda c \int_{y_0/\lambda}^{y/\lambda} \frac{du}{\sqrt{(1+u)^2 - c^2}} = \lambda c \ln[1 + u + \sqrt{(1+u)^2 - c^2}] \Big|_{y_0/\lambda}^{y/\lambda} \stackrel{y_0=0}{=} \lambda \operatorname{arccosh}\left(\frac{y+\lambda}{\lambda}\right)$$

$$b \equiv y(a) - y(-a) = \lambda \cosh \frac{\Delta x}{\lambda} \Big|_{-a}^a = -2\lambda \operatorname{senh} \frac{a}{\lambda} \operatorname{senh} \frac{x_0}{\lambda}, \quad y'(x) = \operatorname{senh} \frac{\Delta x}{\lambda}$$

$$\ell = \int_{-a}^a ds = \int_{-a}^a \cosh \frac{\Delta x}{\lambda} dx = \lambda \operatorname{senh} \frac{\Delta x}{\lambda} \Big|_{-a}^a = 2\lambda \operatorname{senh} \frac{a}{\lambda} \operatorname{cosh} \frac{x_0}{\lambda}$$

Curva catenária com mínimo em $(x_0, y_0 \equiv 0)$

$$y(x) = \lambda \left[\cosh \left(\frac{x-x_0}{\lambda} \right) - 1 \right], \quad \sqrt{\ell^2 - b^2} = 2\lambda \operatorname{senh} \frac{a}{\lambda}, \quad \operatorname{tgh} \frac{x_0}{\lambda} = -\frac{b}{\ell}$$

TM seção 6.4 (equação de Euler II), pgs. 216–217

Derivada total: $\frac{df}{dx} = \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} + \frac{\partial f}{\partial x} = y' \frac{\partial f}{\partial y} + y'' \frac{\partial f}{\partial y'} + \frac{\partial f}{\partial x}$

Temos também $\frac{d}{dx} \left(y' \frac{\partial f}{\partial y'} \right) = y'' \frac{\partial f}{\partial y'} + y' \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right)$, ou seja:

$$\frac{d}{dx} \left(y' \frac{\partial f}{\partial y'} \right) = \frac{df}{dx} - \frac{\partial f}{\partial x} - y' \left[\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) \right] = 0 \text{ (eq. de Euler)}$$

TM eq. (6.39)

Equação de Euler (II):

$$\frac{\partial f}{\partial x} - \frac{d}{dx} \left(f - y' \frac{\partial f}{\partial y'} \right) = 0$$

Assim, se $f = f(y, y'; x)$ não depende explicitamente de x :

TM eq. (6.40); Taylor eq. (6.43)

$$\frac{\partial f}{\partial x} = 0 \quad \rightarrow \quad f - y' \frac{\partial f}{\partial y'} = \text{constante}$$

Problemas isoperimétrico e da curva catenária revisitados

Nestes dois casos: $\frac{\partial \tilde{f}}{\partial x} = 0 \rightarrow \tilde{f} - y' \frac{\partial \tilde{f}}{\partial y'} = \text{constante}$.

Problema isoperimétrico: $\tilde{f}(y, y'; x) = y + \lambda \sqrt{1+y'^2}$

$$\tilde{f} - y' \frac{\partial \tilde{f}}{\partial y'} = \tilde{f} - \frac{\lambda y'^2}{(1+y'^2)^{1/2}} = \tilde{f} - \frac{(\tilde{f}-y)y'^2}{1+y'^2} = \frac{\tilde{f}+yy'^2}{1+y'^2} = y + \frac{\lambda}{(1+y'^2)^{1/2}} = \text{constante}$$

$$(y-y_0)(1+y'^2)^{1/2} = -\lambda \rightarrow y' = \frac{dy}{dx} = \frac{\pm[\lambda^2 - (y-y_0)^2]^{1/2}}{y-y_0} \rightarrow \pm(x-x_0) =$$

TM eq. (6.85)

$$= \int_{y=y_0+\lambda}^y \frac{(y'-y_0)dy'}{[\lambda^2 - (y'-y_0)^2]^{1/2}} = -[\lambda^2 - (y-y_0)^2]^{1/2} \rightarrow (x-x_0)^2 + (y-y_0)^2 = \lambda^2$$

Problema da curva catenária: $\tilde{f}(y, y'; x) = (y + \lambda) \sqrt{1+y'^2}$

$$\tilde{f} - y' \frac{\partial \tilde{f}}{\partial y'} = \tilde{f} - \frac{(y+\lambda)y'^2}{(1+y'^2)^{1/2}} = \tilde{f} - \frac{\tilde{f}y'^2}{1+y'^2} = \frac{\tilde{f}}{1+y'^2} = \text{constante} = \lambda c$$

Curva catenária com mínimo em $(x_0, y_0 \equiv 0)$

$$\rightarrow 1+y'^2 = \frac{1}{(\lambda c)^2} (y+\lambda)^2 \rightarrow x(y) = x_0 + \lambda c \int_{y_0/\lambda}^{y/\lambda} \frac{du}{\sqrt{(1+u)^2 - c^2}} \stackrel{c=1}{=} x_0 + \lambda \operatorname{arccosh}\left(\frac{y+\lambda}{\lambda}\right)$$