

FI 008 – Eletrodinâmica I

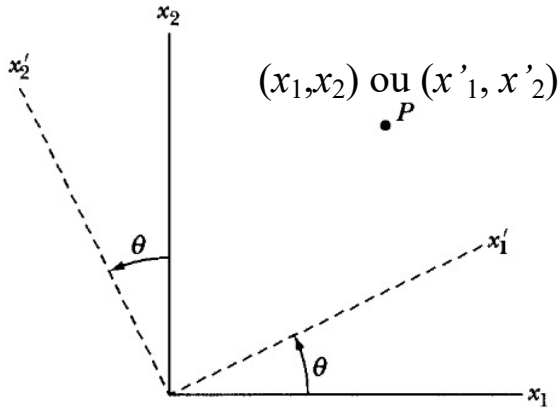
1º Semestre de 2021

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Aula 8

Propriedades de transformação dos campos eletromagnéticos

Rotações:



De maneira geral, em 3D:

$$v'_\alpha = a_{\alpha\beta} v_\beta \left[= \sum_{\beta} a_{\alpha\beta} v_\beta \right]$$
$$(a^{-1})_{\alpha\beta} = a^T_{\alpha\beta} = a_{\beta\alpha} \Rightarrow a^{-1} = a^T$$

A matriz de rotação depende de 3 ângulos (de Euler, e.g.) e é ortogonal.

Definição de vetores, escalares, tensores

Em geral, vetores se transformam como:

$$v'_\alpha = a_{\alpha\beta} v_\beta$$

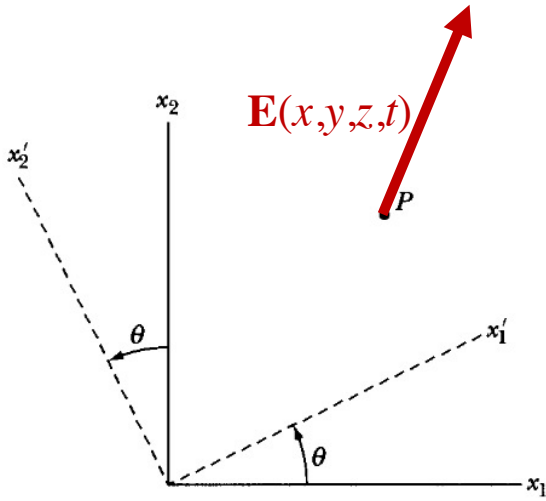
Escalares são invariantes: $e' = e$

Exemplo importante: produto escalar $v_\alpha w_\alpha = v'_\alpha w'_\alpha$

Tensores de 2a. ordem: $T'_{\alpha\beta} = a_{\alpha\gamma} a_{\beta\delta} T_{\gamma\delta}$

Tensores de ordem n : $T'_{\alpha_1\alpha_2\cdots\alpha_n} = a_{\alpha_1\beta_1} a_{\alpha_2\beta_2} \cdots a_{\alpha_n\beta_n} T_{\beta_1\beta_2\cdots\beta_n}$

Campos escalares, vetoriais, tensoriais



$$\phi'(x', y', z') = \phi(x, y, z)$$

$$v'_\alpha(x', y', z') = a_{\alpha\beta} v_\beta(x, y, z)$$

$$T'_{\alpha\beta}(x', y', z') = a_{\alpha\gamma} a_{\beta\delta} T_{\gamma\delta}(x, y, z)$$

onde: $x'_\alpha = a_{\alpha\beta} x_\beta$

Campos escalares: $\Phi(\mathbf{x}), \rho(\mathbf{x})$

Campos vetoriais: $\mathbf{E}(\mathbf{x}), \mathbf{B}(\mathbf{x}) (*), \mathbf{A}(\mathbf{x}), \mathbf{J}(\mathbf{x})$

Campos tensoriais:

$$T_{ij}(\mathbf{x}) = \epsilon_0 E_i(\mathbf{x}) E_j(\mathbf{x}) + \frac{1}{\mu_0} B_i(\mathbf{x}) B_j(\mathbf{x}) - \delta_{ij} \left(\epsilon_0 E^2(\mathbf{x}) + \frac{1}{\mu_0} B^2(\mathbf{x}) \right)$$

O operador nabla

Operador vetorial: $\nabla'_\alpha = \frac{\partial}{\partial x'_\alpha} = a_{\alpha\beta} \nabla_\beta = a_{\alpha\beta} \frac{\partial}{\partial x_\beta}$

$$\nabla\phi \rightarrow \text{campo vetorial}$$

$$\nabla \cdot \mathbf{A} \rightarrow \text{campo escalar}$$

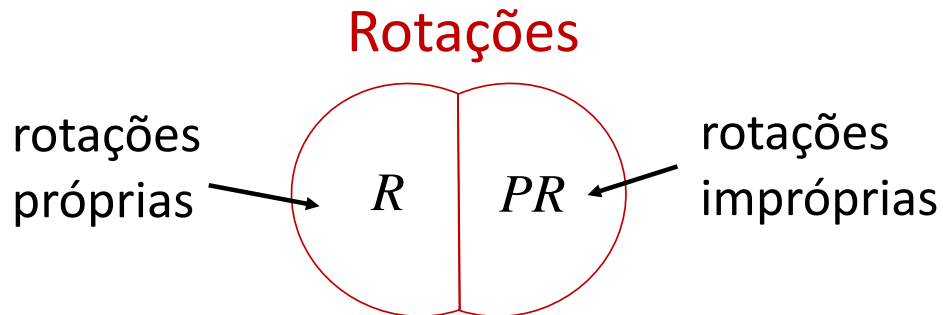
$$\nabla \times \mathbf{A} \rightarrow \text{campo vetorial } (*)$$

Rotações próprias e impróprias

$$\det(a) = \begin{cases} +1 & \text{rotação própria} \\ -1 & \text{rotação imprópria} \end{cases}$$

$$P = \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} \quad \begin{array}{l} \text{Paridade } P: \\ \text{rotação imprópria} \end{array}$$

Toda rotação imprópria pode ser escrita como uma rotação própria, seguida de P .



Vetores, escalares e tensores axiais

Vetores trocam de sinal sob P : $v'_\alpha = -v_\alpha$ (sob P)

Mas o produto vetorial de dois vetores, embora se transforme como um vetor sob rotações próprias, não muda de sinal sob P :

$$\mathbf{v} \times \mathbf{w} \rightarrow \mathbf{v} \times \mathbf{w} \text{ (sob } P\text{)}$$

Diz-se que se trata de um vetor axial (ou pseudo-vetor). Os outros vetores são vetores polares (ou vetor verdadeiro).

Vetores, escalares e tensores axiais

Sob P (ou quaisquer rotações impróprias):

Quantidades verdadeiras:

$$e \rightarrow e$$

$$v_\alpha \rightarrow -v_\alpha$$

$$T_{\alpha\beta} \rightarrow T_{\alpha\beta}$$

$$\vdots$$

$$T_{\alpha_1\alpha_2\cdots\alpha_n} \rightarrow (-1)^n T_{\alpha_1\alpha_2\cdots\alpha_n}$$

Quantidades pseudo:

$$e^{(p)} \rightarrow -e^{(p)}$$

$$v_\alpha^{(p)} \rightarrow v_\alpha^{(p)}$$

$$T_{\alpha\beta}^{(p)} \rightarrow -T_{\alpha\beta}^{(p)}$$

$$\vdots$$

$$T_{\alpha_1\alpha_2\cdots\alpha_n}^{(p)} \rightarrow (-1)^{n+1} T_{\alpha_1\alpha_2\cdots\alpha_n}^{(p)}$$

Reversão temporal

$$T\mathbf{x} = \mathbf{x}$$

$$T\mathbf{v} = -\mathbf{v}$$

SOB $T \rightarrow$ $\left\{ \begin{array}{l} \text{ÍMPAR} : \vec{y} \\ \text{PAR} : \vec{x} \end{array} \right.$

Classificação das quantidades eletromagnéticas segundo suas propriedades de transformação

Table 6.1 Transformation Properties of Various Physical Quantities under Rotations, Spatial Inversion and Time Reversal^a

Physical Quantity		Rotation (rank of tensor)	<i>P</i> Space Inversion (name)	<i>T</i> Time Reversal
<i>I. Mechanical</i>				
Coordinate	\mathbf{x}	1	Odd (vector)	Even
Velocity	\mathbf{v}	1	Odd (vector)	Odd
Momentum	\mathbf{p}	1	Odd (vector)	Odd
Angular momentum	$\mathbf{L} = \mathbf{x} \times \mathbf{p}$	1	Even (pseudovector)	Odd
Force	\mathbf{F}	1	Odd (vector)	Even
Torque	$\mathbf{N} = \mathbf{x} \times \mathbf{F}$	1	Even (pseudovector)	Even
Kinetic energy	$p^2/2m$	0	Even (scalar)	Even
Potential energy	$U(\mathbf{x})$	0	Even (scalar)	Even
<i>II. Electromagnetic</i>				
Charge density	ρ	0	Even (scalar)	Even
Current density	\mathbf{J}	1	Odd (vector)	Odd
Electric field	\mathbf{E}	1	Odd (vector)	Even
Polarization	\mathbf{P}			
Displacement	\mathbf{D}			
Magnetic induction	\mathbf{B}	1	Even (pseudovector)	Odd
Magnetization	\mathbf{M}			
Magnetic field	\mathbf{H}			
Poynting vector	$\mathbf{S} = \mathbf{E} \times \mathbf{H}$	1	Odd (vector)	Odd
Maxwell stress tensor	$T_{\alpha\beta}$	2	Even (tensor)	Even

$\nabla \times \mathbf{A} = \mathbf{B}$
 $\nabla \cdot \mathbf{M}$

PSEUDO-
VECTOR

Equações de Maxwell

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

escalar verdadeiro, par sob T

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

pseudo-vetor, par sob T

$$\nabla \cdot \mathbf{B} = 0$$

pseudo-escalar, ímpar sob T

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

vetor verdadeiro, ímpar sob T

$$x'_\alpha = a_{\alpha\beta} x_\beta \implies x_\alpha = (a^{-1})_{\alpha\beta} x'_\beta = (a^T)_{\alpha\beta} x'_\beta = a_{\beta\alpha} x'_\beta$$

$$\frac{\partial}{\partial x'_\alpha} = \frac{\partial x_\gamma}{\partial x'_\alpha} \frac{\partial}{\partial x_\gamma}$$

$$\implies \frac{\partial}{\partial x'_\alpha} = a_{\alpha\gamma} \frac{\partial}{\partial x_\gamma}$$

$$\Downarrow \\ x_\gamma = a_{\alpha\gamma} x'_\alpha$$

$$\frac{\partial x_\gamma}{\partial x'_\alpha} = a_{\alpha\gamma}$$

$$\frac{\partial x_\gamma}{\partial x'_\alpha} = (a^{-1})_{\gamma\alpha} = (a^T)_{\gamma\alpha} = a_{\alpha\gamma}$$

$$A(t) = A_0 e^{-i\omega t}$$

$$B(t) = B_0 e^{-i\omega t}$$

$$A_0, B_0 \in \mathbb{C}$$

$$A_0 = |A_0| e^{i\delta}$$

$$B_0 = |B_0| e^{i\alpha}$$

$$A_{ph}(t) = |A_0| \cos(\omega t + \delta)$$

$$B_{ph}(t) = |B_0| \cos(\omega t + \alpha)$$

$$\langle A_{ph}(t) B_{ph}(t) \rangle = |A_0| |B_0| \langle \cos(\omega t + \delta) \times \cos(\omega t + \alpha) \rangle$$

$$\langle \rangle = \frac{1}{2} \langle [\cos(2\omega t + \alpha + \delta) + \cos(\delta - \alpha)] \rangle$$

$$= \frac{1}{2} \cos(\delta - \alpha)$$

$$\langle \frac{1}{2} [A(t) + A^*(t)] \frac{1}{2} [B(t) + B^*(t)] \rangle$$

$$= \frac{1}{4} \langle [A(t)B(t) + A^*(t)B(t) + c.c.] \rangle = (*)$$

$$A(t)B(t) = A_0 B_0 e^{-2i\omega t} \quad \langle A(t)B(t) \rangle = 0$$

$$A^*(t)B^*(t) = A_0^* B_0^* e^{+2i\omega t} \quad \langle A^*(t)B^*(t) \rangle = 0$$

$$A^*(t)B(t) = A_0^* e^{+i\omega t} B_0 e^{-i\omega t} = A_0^* B_0$$

$$A(t)B^*(t) = A_0 B_0^*$$

$$\langle x \rangle = \frac{1}{4} \langle A_0^* B_0 + A_0 B_0^* \rangle = \frac{2}{4} \operatorname{Re}[A_0^* B_0] = \frac{1}{2} \operatorname{Re}[A_0^* B_0]$$

$$E(t) = E_0 e^{-i\omega t}$$

$$\langle \frac{E^2}{\rho v} (t) \rangle = \frac{1}{2} \operatorname{Re}[E_0^* E_0] = \frac{1}{2} |E_0|^2$$